Progress Report No. 1

Studies of Adsorption-Desorption and Partitioning of Dredged Material Contaminants

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## Acknowledgment

This progress report contains research performed for both the EPA Narragansett Environmental Research Laboratory and the NOAA Office of Marine Pollution Assessment. The topic of each project is similar - the particles being investigated are quite different (Sewage Sludge and Dredged Material). It is felt that it would be more instructive to combine the progress reports in order to highlight the similarities found for both types of materials and to present a more comprehensive discussion of the status of the sorption model.

### Summary

The experimental work this year has focused on the generation and analysis of consecutive desorption isotherms using atomic absorption techniques in order to analyze the desorption behavior of the native particle bound metal as it interacts with various aqueous phases. The radionuclide isotherms were used mainly to validate the desorption model, investigate the effects of adsorption and desorption times, and perfect technique.

The desorption model described below fits the experimental data quite well and generally yields consistent results for the isotherm parameters. A nonlinear least squares fitting program estimates the isotherm parameters and the desorable fraction of the metal. It also computes standard errors of the estimates so that comparisons can be made considering the uncertainty of the estimates.

It has been found that desorption of native metal from sewage sludge and Black Rock Harbor sediment follow the same framework as previously found applicable to radionuclides and clean particles (clay and quartz desorption). The difference is that a substantially smaller fraction of the total native metal is desorbable. Whereas with metal adsorbed to clean particles just prior to desorption, approximately one half or more of the metal will eventually desorb, the percentage for native metal is much less. For pure water desorption less than 1% desorbs. With seawater the percentage is higher, perhaps 5-20%, depending upon which metal is being considered, but the data are as yet not extensive enough for statistical precision. With respect to organic chemicals the primary effort has been to complete an analysis of literature data in order to investigate the applicability of the same desorption model (Appendix II). Remarkably enough the reversible component partition coefficient expression derived from the metal-clean particle experiments fits the neutral hydrophobic chemical sorption data over seven orders of magnitude in partition coefficient (Appendix II, fig. 4). In addition, the isotherm parameters can be predicted from chemical (octanol/water partition coefficient,  $K_{ow}$ ) and particle properties (fraction organic carbon,  $f_{oc}$ , and particle concentration, m). Thus for neutral hydrophobic chemicals and suspensions of particles, the desorption partitioning problem is essentially solved.

The particle interaction model (Appendix I and II), from which the expression for reversible partition coefficient follows suggests that for stationary particles as would be the case in sediments the partition coefficient would be the same as the low particle concentration limit of partition coefficient found in suspension experiments. This has been found to be the case for hexachlorobiphenyl. Syringe migration experiments produced the same partition coefficient as predicted from the particle interaction model applied to suspensions (Appendix II and III). We intend to examine whether this is a general finding by concentrating a significant effort on syringe migration experiments.

The examination of desorption of hydrophobic chemicals from dredged material and sewage sludge will be made using radiotagged chemicals with high specific activity ( $\sim$  10 Ci/m mole). Two have been purchased: benzo(a) pyrene (log K<sub>ow</sub> = 6.5) and 2-amino anthracene (log K<sub>ow</sub> = 4.13). These are suitable for both consecutive desorption experiments and syringe migration studies. Other chemicals are being acquired that fill in the range of K<sub>ow</sub>. These are of lower specific activity ( $\sim$  10m Ci/m mole) and are only suitable for syringe experiments.

The question of complete tagging of a particle with the tracer neutral organic chemical has been addressed by work reported by Karickhoff and Morris (1984, in press). They estimate the time to complete adsorption and desorption equilibrium to be on the order of  $f_{C}$  (days). Hence only reversible component sorption is experimentally accessible for B(a)P regard-less of whether native chemical or tracers are used. Even for 2-amino

anthracene and Black Rock Harbor sediment ( $f_{oc} = 5\%$ ) a tagging time of 500 days is suggested. Long term tagging experiments will be attempted but it appears that no complete equilibration is possible.

Additional experiments using model particles have been performed in order to elucidate the particle concentration effect. Two micron spherical glass beads have been used to generate nickel adsorption-desorption isotherms at varying particle concentrations. These particles also exhibit the particle concentration effect (fig. 11) in accord with the proposed model (fig. 12). We regard this as strong proof that the effect is not an artifact of particle separation. Field data for nickel and PCB sorption also conform (fig. 13-14).

Publication of results is proceeding. Appendix I has been submitted to Environmental Science and Technology and the reviews (mostly favorable) have just been received. Appendix II will be submitted to Chemosphere shortly. The syringe migration paper (Appendix III) has been submitted to Environmental Science and Technology. We expect to write a paper detailing the findings of the sewage sludge and sediment results.

# I. Introduction

The purpose of this research project is to continue the development and testing of a model of phase partitioning that describes the adsorption and desorption of heavy metals and organic chemicals to suspended and sedimented particles. The model is based upon separating sorbed chemical into two components: the reversible or labile fraction,  $r_x$ , which reversibly adsorbs and desorbs, and a resistant or non-labile fraction,  $r_o$ , which does not appreciably desorb in time scales of hours to days. The partitioning of the reversible component is described via a linear isotherm:  $r_x = \pi_x c$  with partition coefficient,  $\pi_x$ . A completely developed theory would:

- 1) specify the behavior of  $\pi_{_{\rm X}}$  as a function of chemical, particle, and aqueous phase properties
- 2) specify, or provide a measurement method for particles such as sewage sludge or dredged sediment for which the exposure history is unknown, of the reversibly bound fraction  $r_x/r_T$  of the total particle bound chemical concentration,  $r_T$ .
- 3) be applicable to both suspended particles and stationary sediments.

Progress toward this goal has been appreciable and a usable framework is now available which appears to apply to all particle types examined so far. Some preliminary evidence suggests that for the simplest situation (neutral hydrophobic chemicals and organic carbon containing particles) the above requirements for a complete theory are being approached.

II. The Model

Consider a sample of dredged sediment or sewage sludge with a total concentration  $r_T (\mu g/g)$  of sorbed chemical. A portion,  $r_X^{(0)} (\mu g/g)$  of the chemical is reversibly sorbed - the zero superscript is a convenient notation for the initial state of the particles before desorption is initiated. The remainder,  $r_o = r_T - r_X^{(0)}$ , is not desorable during the short time scales being examined.

A convenient way of thinking about desorption is to analyze the results of consecutive cycles of desorption into a fixed volume of uncontaminated aqueous phase at a fixed particle concentration, m. For the first cycle, the reversible component rapidly equilibrates to produce an aqueous concentration,  $c^{(1)}$  and a new (smaller) reversible component concentration,  $r_x^{(1)}$ . These are related via the partition coefficient:

$$r_{x}^{(1)} = \pi_{x} c^{(1)}$$
(1)

If no chemical is lost to the experimental vessel then by mass balance:

$$c_{\rm T} = mr_{\rm T} = m(r_{\rm o} + r_{\rm x}^{(0)}) = m(r_{\rm o} + r_{\rm x}^{(1)}) + c^{(1)}$$
 (2)

where  $c_T$  is the total chemical concentration on a volumetric basis (µg/L), and m is the particle concentration (g/L). Hence:

$$mr_{x}^{(0)} = mr_{x}^{(1)} + c^{(1)} = (m\pi_{x} + 1)c^{(1)}$$
(3)

where the second equality follows from the partitioning eq. (1). Thus:

$$c^{(1)} = \frac{m r_{x}^{(0)}}{\frac{1}{1 + m\pi_{x}}}$$
(4)

For the second (i=2) and subsequent cycles the vessel is centrifuged, a volume of aqueous phase is removed ( $\alpha$  = volume fraction remaining), replaced with uncontaminated aqueous phase such that m is unchanged, and the vessel is re-equilibrated. Again by mass balance:

$$mr_{o} + mr_{x}^{(1)} + \alpha c^{(1)} = mr_{o} + mr_{x}^{(2)} + c^{(2)}$$
 (5)

so that:

$$c^{(2)} = \frac{\alpha + m\pi}{1 + m\pi} c^{(1)}$$
(6)

$$c^{(1)} = \left(\frac{\alpha + m\pi}{1 + m\pi}\right) c^{(1)} c^{(1)}$$
(7)

so that the dissolved concentration in consecutive desorptions follow a decreasing geometric series and these data,  $c^{(i)}$ , would plot as a straight line versus i on semi-log paper. The slope would give  $\pi_x$ . Using eq. (4) for  $c^{(1)}$  yields

$$c^{(i)} = \frac{m}{1 + m\pi_{x}} \left(\frac{\alpha + m\pi_{x}}{1 + m\pi_{x}}\right) r_{x}^{(0)} r_{x}^{(0)}$$
(8)

so that the intercept yields  $r_x^{(0)}$ , the initial reversible component concentration. Finally the sorbed concentration at each cycle is:

$$r^{(i)} = r_0 + r_x^{(i)}$$
 (9)

$$= r_{o} + \pi_{x} c^{(i)}$$
(10)

which yields r, the resistant component concentration.

For each particle type (sludge, Black Rock Harbor sediment) and aqueous phase (nanopure water, seawater) a set of isotherms are generated by using different particle concentrations, m. Typically four particle concentrations spanning the range 100-3000 mg/L are used. Fig. 1 illustrates typical results for aqueous concentrations  $c^{(i)}$  (left hand side) and  $r^{(i)}$  (right hand side). The semi-logarithmic behavior of  $c^{(i)}$  is apparent as is the curvature of  $r^{(i)}$  reflecting the presence of a resistant component concentration.

In order to analyze each of these isotherm sets it is necessary to know how  $\pi_x$  varies with respect to  $m_j$  since each isotherm within a set is obtained at a differing particle concentration. Contrary to expectations,  $\pi_x$ can be a strong function of m. This central feature of reversible partitioning is discussed below and in Appendices I and II where a particle interaction model is proposed to establish the functional form of the relationship. That such an effect is important can be seen directly from the aqueous concentration versus desorption cycle behavior. The logarithmetic concentration behavior from eq. (7) is:

$$\ln c_{j}^{(i)} = (i-1)\ln(\frac{\alpha + m_{j}\pi_{x}}{1 + m_{j}\pi_{x}}) + \ln c_{j}^{(1)}$$
(11)

As solids concentration,  $m_j$ , increases the slope should progressively become less negative until it approaches  $\ln(m_j \pi_x/m_j \pi_x) = \ln(1) = 0$ . Note, however, in fig. 1 that at the higher solids concentrations, and in the case of the Black Rock Harbor sediment all the slopes of each isotherm, are nearly the same (the lines are parallel) indicating that  $m_j \pi_x$  is constant so that  $\pi_x$ must be decreasing with increasing  $m_j$ . Table 1 presents the  $\pi_x$  derived from the slopes of  $\ln c^{(1)}$  versus i and  $m_j$  for these isotherms. The reciprocal relationship is evident.

The vertical displacement of  $c_j^{(1)}$ , the cycle 1 concentration, as  $m_j$  increases is also related to this effect. As  $m_j$  increases the aqueous concentration should increase following:

$$c_{j}^{(1)} = \frac{m_{j}}{1 + m_{j}\pi_{x}} r_{x}^{(0)}$$
(12)

Again if  $\pi_x$  were constant than as  $m_j$  increases  $c_j^{(1)}$  should approach  $r_x^{(0)}/\pi_x$ , a constant. But these data indicate that the displacement persists even at large particle concentrations, suggesting again that  $m_j \pi_x$  is constant.

These arguments are meant only illustrate the consequences of  $\pi_x$  varying with m. The quantitative results are obtained by fitting the data to eq. (8) and eq. (10) using a nonlinear least squares program. The program chooses the model parameters such that they minimize the sum of the squared differences between observations  $\stackrel{\Lambda}{c}$ ,  $\stackrel{\Lambda}{r}$  and model calculations  $c_i^{(i)}$ ,  $r_i^{(i)}$ :

$$\sigma^{2} = \sum_{i j} \sum_{j} (\log c_{j}^{(i)} - \log(c_{j}^{(i)}))^{2} + (\log^{\Lambda} r_{j}^{(i)} - \log(r_{j}^{(i)}))^{2}$$

where  $c_j$  and  $r_j$  are given by eq. (8) and (10). The reversible component partition coefficient is given by:

$$\pi_{x} = \frac{\pi_{xc}}{1 + m\pi_{xc}/\nu_{x}}$$
(13)

which establishes the functional dependency of  $\pi_x$  to particle concentration, m. Appendix I and II present the rationale and the derivation. The parameters are  $\pi_{xc}$  (the classical, low particle concentration limiting partition coefficient  $\pi_x \rightarrow \pi_{xc}$  as  $m \rightarrow 0$ ) and  $\nu_x$ , the slope of  $\pi_x$  versus m for large particle concentration ( $\pi_x \rightarrow \nu_x/m$  for m large). These specify the behavior of  $\pi_x$ . Fig. 2 presents some early adsorption-single desorption isotherm data fit to eq. (13) using nonlinear least squares. Note the strong inverse relationship between  $\pi_x$  and m for large m as well as the approach to  $\pi_{xc}$  as m decreases. For a complete analysis of consecutive desorption isotherms the remaining parameters are  $r_x^{(0)}$  and  $r_o$ , the reversible and resistant component concentrations. Thus four parameters are fit to the data. In fact since  $r_x^{(0)} + r_o = r_T$ , which is presumably constant, only three parameters are needed to completely characterize a set of isotherms. This is a reasonably parsimonious representation of a large quantity of data.

If a complete theory were available it would then provide the parameters  $\pi_{xc}$ ,  $\nu_{x}$  and  $r_{x}^{(0)}/r_{T}$  as a function of chemical, particle, and aqueous phase properties. It is clear at this stage of the research that the three parameter formulation is descriptive of data for many chemical and particle types and in particular for sewage sludge and dredged sediment - heavy metal sorption. In the following sections this framework is applied to consecutive desorption experiments using both radionuclides and native heavy metals.

III. Consecutive Desorption - Cd<sup>115</sup>: Effect of Adsorption and Desorption Times

The original research plan was to investigate the effects of various experimental variables: adsorption time, desorption time, etc. using radionuclides. The hope was that by using sufficiently long adsorption times the resistant component sites could be radiolabeled. Fig. 3-5 present the results for Ridgewood sludge and Black Rock Harbor sediment in seawater. Table 2 lists the isotherm parameters. Fig. 3 compares adsorption times of  $t_{ads} = 1$ , 4 and 9 days from the same tagging vessel. Some increase in r is observed but the effect is not dramatic. For  $t_{ads} = 18$  and 25, Fig. 4, a dramatic change is evident. However these isotherms were obtained using a previously uncapped vessel. Progressive color changes in the first adsorption vessel (Fig. 3 results) indicated oxidation was occurring whereas the second tube maintained its color until it was sampled.

Similar results were obtained for the Black Rock Harbor sediment (Fig. 5). Small changes were observed and the somewhat erratic results are attributed to partial oxidation.

It was decided that the adsorption tagging procedure using radionuclides was not convincing since one could never be sure that the radiotag was isotopically exchanged with all sites. Hence the useful information in Table 2 is only the reversible component parameters,  $v_x$  and  $\pi_{xc}$ . These ranged from  $v_x = 0.2 - 2.0$  and  $\pi_{xc} \sim 2000 - 8000$  L/kg. Where  $\pi_{xc}$  is not listed in this and subsequent tables it is because the data are representable by the high particle concentration limit:  $\pi_x = v_x/m$  so that  $\pi_{xc}$  could not be estimated.

The effect of increasing desorption time for each cycle is examined in Fig. 6 and Table 3. No significant changes were observed from  $t_{des} = 1$  hr to 5-7 days.

IV. Consecutive Desorption - Native Metal Results

The major experimental effort for the latter half of this year was directed toward perfecting the techniques for directly measuring total and dissolved metal concentrations using atomic absorption flame and graphite techniques. Clean procedures and digestion techniques were established and checked. Duplicates are employed for each particle concentration. One hour desorptions are used for each cycle.

Fig. 7 and Table 4 illustrate the results for Ridgewood sludge and nanopure water as the aqueous phase. Copper, chromium and zinc were analyzed. Only a small percent of any of the total metal is in the reversible

Isotherm	1					
:	Lads	××	т хс	л×	ц Ч	$\sigma(\log_{10})$
#	(days)		(L/kg)	(µg/g)	(a/a)	17
186		11.2 ± 2.65	4570 ± 716	2.30 ± 0.10	0.214 ± 0.067	0.348
188	4	2.24 ± 0.724	3390 ± 1370	2.76 ± 0.20	$1.67 \pm 0.20$	0.484
193	6	2.29 ± 0.68	5500 ± 2300	2.47 ± 0.19	1.12 ± 0.18	0.546
196*	18	5.30 ± 16.7	832 ± 620	$0.147 \pm 0.02$	4.79 ± 0.42	0.495
202 <sup>*</sup>	25	0.238 :0.11	I	0.216 ± 0.02	<b>3.84 ± 0.363</b>	0.53
			Cd <sup>115</sup> - Black Rock Harbor - Seawater	Harbor - Seawater		
192	11	0.759 ± 0.41	8320 = 6980	19.1 ± 2.9	3.45 ± 2.0	1.18
197	18	0.466 ± 0.09	3550 ± 2300	15.1 ± 0.84	6.89 ± 0.49	0.401
201 <sup>*</sup>	25	0.473 ± 0.38	I	14.5 ± 4.3	1.86 ± 1.77	2.2

Table 2

Effect of Adsorption Time Cd<sup>115</sup> - Ridgewood Sludge - Seawater

	σ(log <sub>10</sub> )	1.16 0.88 2.53			
	r <sub>o</sub> (μg/g)	1.32 ± 0.45 1.03 ± 0.15 0.10 ± 0.08			
Table 3 Effect of Desorption Time 115 - Ridgewood - Seawater	(0) r <sub>x</sub> (μg/g)	5.11 ± 0.74 5.40 ± 0.15 3.58 ÷1.31			
Table 3 Effect of Deso Cd <sup>115</sup> - Ridgewood	π xc (L/kg)	I I I			
	~×	$0.562 \pm 0.24$ $0.295 \pm 0.08$ $0.229 \pm 0.16$			
	t <sub>ads</sub> (days)	0.042 (1 hr) 5 7		-	
	Isotherm #	115 121 114		<u>, , , , , , , , , , , , , , , , , , , </u>	

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Table	
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# AA Isotherm Parameters

Ridgewood Sludge - Nanopure Water (pH = 7.8-8.1)

$1.97 \pm 0.88$ $6410 \pm 8350$ $12.4 \pm 1.14$ $2900 \pm 194$ $0.546$ $3.70 \pm 0.75$ $5150 \pm 1550$ $2.11 \pm 0.12$ $807 \pm 26$ $0.260$ $2.80 \pm 1.70$ $ 9.57 \pm 3.03$ $3440 \pm 340$ $0.795$ $2.80 \pm 1.70$ $ 9.57 \pm 3.03$ $3440 \pm 340$ $0.795$ $2.80 \pm 1.71$ $2.900 \pm 829$ $7.21 \pm 0.67$ $123$ $0.280$ $4.28 \pm 1.71$ $2090 \pm 829$ $7.21 \pm 0.67$ $1306 \pm 211$ $1.06$ $2.74 \pm 7.16$ $1850 \pm 5800$ $2.24 \pm 0.66$ $1306 \pm 211$ $1.06$		Sorbate	×	π (L/ŘĜ)	τ (0) (μ <sup>X</sup> /g)	r (μĝ/g)	σ(log <sub>10</sub> )
5150 ± 1550     2.11 ± 0.12     807 ± 26       -     9.57 ± 3.03     3440 ± 340       200 ± 829     7.21 ± 0.67     2660 ± 123       2090 ± 829     7.21 ± 0.67     2660 ± 123       1850 ± 5800     2.24 ± 0.66     1306 ± 211	Cu		$1.97 \pm 0.88$	6410 ± 8350	12.4 ± 1.14	2900 ± 194	0.546
<ul> <li>9.57 ± 3.03 3440 ± 340</li> <li>240 ± 340</li> <li>240 ± 340</li> <li>240 ± 340</li> <li>240 ± 340</li> <li>25800 ± 829</li> <li>7.21 ± 0.67</li> <li>1306 ± 211</li> <li>1306 ± 211</li> </ul>	5		3./U ± U./S	$5150 \pm 1550$	$2.11 \pm 0.12$	807 ± 26	0.260
ock Harbor Sediment - Borate Buffer (pH = 9)         2090 ± 829       7.21 ± 0.67       2660 ± 123         1850 ± 5800       2.24 ± 0.66       1306 ± 211	Zn	· · ·	2.80 ± 1.70	I	9.57 ± 3.03	3440 ± 340	0.795
ock Harbor Sediment - Borate Buffer (pH = 9)         2090 ± 829       7.21 ± 0.67       2660 ± 123         1850 ± 5800       2.24 ± 0.66       1306 ± 211							
2090 ± 829 7.21 ± 0.67 2660 ± 123 1850 ± 5800 2.24 ± 0.66 1306 ± 211			Black Rock 1	Harbor Sediment - Bo	orate Buffer (pH = 9)		
	Cu		4.28 ±1.71 2.74 ± 7.16	2090 ± 829 1850 ± 5800	7.21 ± 0.67 2.24 ± 0.66	2660 ± 123 1306 ± 211	0.28 1.06

	σ(log <sub>10</sub> )	0.435 0.486		0.508 1.44 0.783	
	ro (µg/g)	613 ± 440 558 ± 3300	Seawater	1832 ± 168 937 ± 170 1350 ± 119	
le 5 Parameters <u>- Nanopure Water</u> 1 for each cycle	r <sub>x</sub> (µg/g)	24.9 ± 4.9 1350 ± 1240	25%	71.6 ± 29 3.42 ± 0.64 362 ± 67	
Table 5AA Isotherm ParametersRidgewood Sludge - Nanopure Wate1 day desorption for each cycle	<sup>π</sup> xc (L/kg)	4680 ± 3400 23,990 ± 35,000	Ridgewood Sludge - 75% Nanopure Water,	2360 ± 2700 - -	
	Sorbate v <sub>x</sub>	Cr 6.56 ± 3.12 Zn 31.6 ± 28	Ridgewo	Cu 22.8 ± 108 Cr 0.875 ± 0.50 Zn 1.91 ± 0.85	
	Isotherm #	12		5 6 6	

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# Table 7 Isotherm Parameters for N<sup>63</sup> Sorption Borate Buffer - pH=9

Sorbent	v x	<sup>π</sup> xc (L/kg)	σ(log <sub>10</sub> )
2µ glass spheres	0.550 ± 0.344	826 <del>*</del> 647	0.300
Quartz particles	$0.934 \pm 0.246$	$1965 \pm 618$	0.217
Cellulose	0.454 ± 0.087	16,070 ± 4900	0.100
Montmorillonite	$0.931 \pm 0.110$	48,380 ± 8510	0.080

# Individual Data Sets

# Joint Data Set

Sorbent	<sup>π</sup> xc
2µ glass spheres	715 ± 246
Quartz particles (Mil-U-Sil 5)	2380 ± 938
Cellulose	11,240 ± 3700
Montmorillonite	57,100 ± 21,100

component:  $r_x^{(0)}/r_T < 0.5\%$  although the partition coefficients are not large:  $\pi_{xc} \sim 5000 - 6400$  L/kg and  $v_x \sim 2-4$  is in the range of that found for the radioisotopes. The conformity of the model to the data is evident as well as the reproducibility of the data - note that the log scale for  $r^{(i)}$ is one cycle.

Fig. 8 (Table 4) present the results for Black Rock Harbor sediment in borate buffer (0.1 mM) which was necessary to stabilize the pH. Again only a small percentage of the metal is in the reversible phase  $r_x^{(0)}/r_T < 0.3\%$ . Partitioning parameters are not dissimilar to those for Ridgewood sludge although the  $\pi_{xc}$ 's do appear somewhat lower. The similarity is somewhat surprising considering the difference between the two particle types.

Fig. 9 (Table 5) illustrates the effect of 1 day desorption times between cycles. Some increase in reversible component sorption is suggested but the statistical errors of estimation are too large to draw any firm conclusions.

Fig. 10 (Table 5) illustrates the effects of desorption into dilute seawater. It appears that substantially more copper  $(r_x^{(0)}/r_T \sim 4\%)$  and zinc  $(r_x^{(0)}/r_T \sim 20\%)$  is desorbable in contrast to the nanowater case of < 0.5%. These seawater experiments are initial attempts. The matrix effects are being eliminated as well as problems with contamination. However it is clear even from these data that substantial fractions of the native heavy metal are not desorable into seawater. The use of completely reversible partitioning theory would greatly overestimate the aqueous concentration and greatly underestimate the quantity of metal that remains bound to the sludge and/or sediment as it settles to the sediment.

# V. Summary of Metal Isotherm Parameters

The data generated and analyzed to date are summarized in Table 6 the purpose of which is to provide an overview of the behavior of the isotherm parameters. The most striking feature is that there is relatively little variation in  $v_x$  and  $\pi_{xc}$  for all the metal/sorbent/aqueous phase combinations that have been investigated:  $v_x = .7 - 4.0$  and  $\pi_{xc} = 2000 - 6500$  L/kg. One would have expected  $\pi_{xc}$  to vary more dramatically especially as the aqueous phase is changed from nano-pure water to seawater reflecting the ionic strength increase. Only small variations in  $v_x$  are not too surprising anymore (although still puzzling) since only small changes are observed for organic chemicals as well (Appendix II).

The most dramatic variation is in the fraction desorable. For Cd which is adsorbed prior to the desorption, 30-60% is desorable. For desorption of native metal into low ionic strength solutions, the desorable fraction is less than 1%. The preliminary dilute seawater results suggest that this fraction may be higher for copper and zinc at least. However it is too soon to tell conclusively. More complete seawater desorption isotherms are planned to establish the validity of this trend.

Sorbate	Sorbent	Aqueous Phase	ν <sub>x</sub> /	"xc (L/kg)	r <sub>x</sub> <sup>(0)</sup> /r <sub>T</sub> (%)
Cd <sup>115</sup>	Ridgewood	Seawater	2.5	4000	60
$Cd^{115}$	Black Rock	Seawater	0.65	6000	30
Cu	Ridgewood	Nanowater	2.0	6500	0.4
Cr	Ridgewood	Nanowater	3.7	5200	0.25
Zn	Ridgewood	Nanowater	2.8	-	0.28
Cu	Black Rock	Borate	4.3	2000	0.27
Cr	Black Rock	Borate	2.7	1900	0.17
Cu	Ridgewood	1/4 Seawater	-	2400	3.9
Cr	Ridgewood	1/4 Seawater	0.88	-	0.36
Zn	Ridgewood	1/4 Seawater	1.9	-	21
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Table 6 Summary of Isotherm Parameters

# VI. Particle Concentration Effect

The most surprising and puzzling feature of reversible component partitioning is the particle concentration effect. Attempting to establish that this effect is not an artifact has prompted a number of experimental designs (Resuspension and Dilution experiments, Appendix I).

Another approach is to use a particle that is unambiguously separated by centrifugation so that no third phase arguments apply. Toward this end we have examined 2  $\mu$ m glass spheres (Duke Scientific Corp.) in the Ni<sup>63</sup>-Borate system. Fig. 11 presents the adsorption data which clearly shows the particle concentration effect. It is difficult to imagine what artifact could account for these results.

In addition, as a model particle for organic coatings we examined the behavior of cellulose particles over a range of particle concentrations. These data together with montmorillonite and quartz particle data are analyzed using nonlinear least square fits to obtain  $v_x$ ,  $\pi_{xc}$  for each particle type. They are listed in Table 7. Essentially the same  $v_x$  results for each data set. A joint fit with a common  $v_x$  is also tabulated and shown in Fig. 12 as  $\pi_x$  vs m and in normalized form. Note the excellent fit of the data regardless of particle type. Whereas  $\pi_{xc}$  ranges from  $\sim$  700 L/kg for the glass spheres to 11,200 L/kg for cellulose to 57,000 L/kg for montmorillonite  $v_y$  is essentially constant.

The glass sphere data will be refined but it is clear that as simple a particle as can be imagined exhibits the same qualitative and quantitative behavior as more complex polydisperse particle suspensions (clay, quartz, sediment, sludge).

# VII. Field Data

The sorption model employed in this project is not just descriptive of laboratory experimental data. Field data for which both total and dissolved concentrations and suspended solids measurements are available also appear to display the particle concentration effect and are in conformity with the particle interaction model prediction, eq. (13). An analysis of STORET heavy metal data is in process. Fig. 13 presents the results for nickel partition or, more precisely, distribution coefficient for all appropriate STORET-stream data. The distribution coefficients are log averages over one-half log 10 cycles of suspended solids concentration intervals: 3-10, 10-30; 30-100 mg/L etc. The top left panel compares the log mean + log standard deviation to the reversible partition coefficient expression eq. (13). Note that  $\pi_{\rm xc}$  is an order of magnitude larger than that found for montmorillonite but that  $\nu_{\rm x}$  is essentially the same. The top right panel compares the same data to an equation which includes a lower bound. It fits the data marginally better but does not change the estimates for  $\pi_{\rm xc}$  and  $\nu_{\rm x}$ . The lower panels compare the behavior of the distribution coefficients to pH and alkalinity. Note that although the trends are in the direction to be expected from conventional theory - increasing sorption with increasing pH and decreasing sorption with increasing alkalinity - the effects are small when compared to the changes that are related to particle concentration effects.

Suitable organic chemical field data are more scarce. Only for PCB's has a substantial quantity of data been accumulated. Fig. 14 illustrates the results. As with the nickel data, these are log averages over one half log cycles in particle concentrations for each data set. Both the field data distribution coefficients and the lab data for hexachlorobiphenyl reversible partition coefficients follow the prediction of the particle interaction model eq. (13). In particular the data from the Black Rock Harbor field sampling after the disposal (LI Sound) is in conformity with the model predictions. Note the significant differences between the laboratory adsorption and the reversible component partition coefficients, indicating the extent of non-reversibility exhibited by short time scale adsorption-desorption.

Further efforts to collect such data sets are planned since the description of field data phase partitioning is a goal of the research. These comparisons illustrate both the utility of the approach and the importance of both laboratory derived experimental information that is used to formulate the model, and the field data analyses that can be made within that framework in order to assess its applicability.

# VIII. Research Plan for '84-'85

The final year's effort will be devoted to investigating two principal questions. The first is to establish the isotherm parameters and, most important, the fraction of total heavy metal in sewage sludge and Black Rock Harbor sediment that is desorable in seawater. Toward this end it is planned to analyze three different sludge samples, one of which is heavily contaminated, and one other sample of dredged material. The procedures we have developed will be applied and the results of the analysis will be the isotherm parameters and the fraction desorable. Future work would then be required to establish relationships between the isotherm parameters, the desorable fraction, and the characteristics of the particles and aqueous phases.

The second question to be addressed concerns the behavior of stationary sediments. The hexachlorobiphenyl results suggests that the low particle limit partition coefficient,  $\pi_{\rm xc}$ , is the partition coefficient that applies to stationary sediments. The generality of this result is the topic for this year's investigations. Syringe migration experiments are in progress in order to examine whether in fact the low particle concentration partition coefficient applies to stationary sediments. It is planned to examine both heavy metal and organic chemical migration using this technique. Sewage sludge and Black Rock Harbor sediments will be employed as well as clean sorbents (clay and quartz particles). Seawater will be the aqueous phase and the syringes will be sterilized to prevent microbial activity.

In order to prevent microbial degradation, we will investigate autoclaving as a sterilization method. This appears to be the only feasible method that guarantees sterility. In addition it is not a violent treatment. The temperature and pressure conditions are such that they do not physically disrupt even bacterial cells (H. Pritchard EPA-ERL Gulf Breeze, personal communication). Comparisons using short term adsorption-desorption isotherms with and without autoclaving will be used to confirm that no significant changes in partitioning occurs. The purpose of these syringe experiments is to examine the universality of the particle interaction model prediction that the low particle concentration limit derived from suspension experiments applies to stationary sediments. If this is indeed the case then the problem of phase partitioning in stationary sediments and in particle suspensions will have been shown to be essentially the same problem, and the progress that has been made in the latter problem (Appendix II) can be applied directly to the stationary sediment issue.

- Fig. 1 Examples of Consecutive Desorption Isotherms Nickel-Ridgewood Sludge-Seawater m = 86,430,860,2590 mg/L Cadmium-Ridgewood Sludge-Seawater m = 114,570,1140,3400 mg/L Cadmium-Black Rock Harbor Sediment-Seawater m = 180,690,2080,3470 mg/L
- Fig. 2 Reversible Component Partition Coefficient versus Particle Concentration
- Fig. 3 Effect of Adsorption Time. Cadmium-Ridgewood Sludge-Seawater. Adsorption time = 1 day (top), 4 day (mid), 7 day (bottom). m = 114,570,1140,3400 mg/L
- Fig. 4 Effect of Adsorption Time. Cadmium-Ridgewood Sludge-Seawater. Adsorption time = 18 day (top), 25 day (bottom). m = 114,570,1140, 3400 mg/L
- Fig. 5 Effect of Adsorption Time. Cadmium-Black Rock Harbor Sediment-Seawater. Adsorption time = 11 day (top), 18 day (mid), 25 day (bottom). m = 170,690,2080,3470 mg/L
- Fig. 6 Effect of Desorption Time. Cadmium-Ridgewood Sludge-Seawater. Desorption time = 1 hr. (top), 5 day (mid), 7 day (bottom). m = 85,425,850,2550 mg/L
- Fig. 7 Native Metal Desorption. Ridgewood Sludge-Nanopure Water. Desorption time = 1 hr. Copper (top) Chromium (mid), Zinc (bottom). m = 340,680,1370,2730 mg/L
- Fig. 8 Native Metal Desorption. Black Rock Harbor Sediment-Borate Buffer. Desorption time = 1 hr. Copper (top) Chromium (bottom). m = 340,680,1370,2730 mg/L
- Fig. 9 Native Metal Desorption. Ridgewood Sludge-Nanopure Water. Desorption time = 1 day. Copper (top) Chromium (mid), Zinc (bottom). m = 340,680,1370,2730 mg/L
- Fig. 10 Native Metal Desorption. Ridgewood Sludge-1/4 Seawater, 3/4 Nanopure Water. Desorption time = 1 hr. Copper (top), Chromium (mid), Zinc (bottom). m = 340,680,1370,2730 mg/L
- Fig. 11 Adsorption Isotherms. Nickel-Borate-Glass Spheres. Particle concentrations as indicated.
- Fig. 12 Reversible Component Partitioning, π versus m (top). Normalized Partition Coefficient versus normalized particle concentration (bottom).

Fig. 13 Nickel Distribution Coefficient versus Particle Concentration. STORET stream data, log mean ± log std. deviation No. of data point indicated.

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Fig. 14 PCB Distribution Coefficient versus Particle Concentration.

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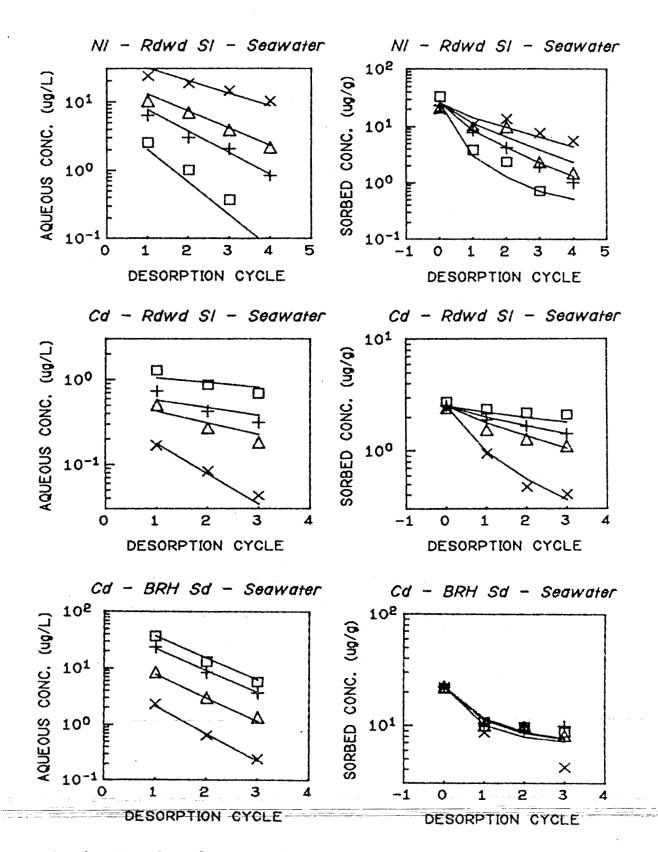
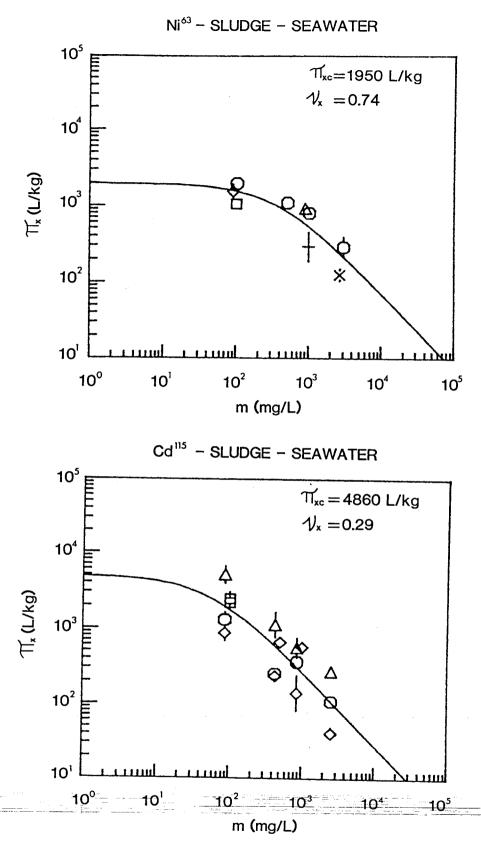
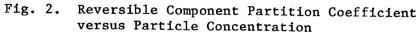


Fig. 1. Examples of Consecutive Desorption Isotherms - Nickel-Ridgewood Sludge-Seawater m = 86,430,860,2590 mg/L, Cadmium-Ridgewood Sludge-Seawater m = 114,570,1140,3400 mg/L, Cadmium-Black Rock Harbor Sediment-Seawater m = 180,690,2080,3470 mg/L





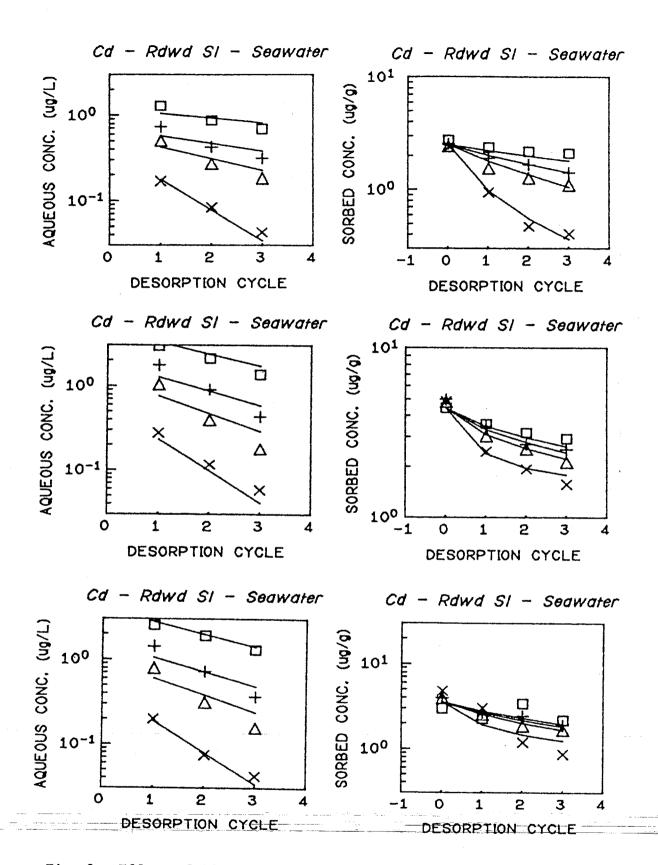


Fig. 3. Effect of Adsorption Time. Cadmium-Ridgewood Sludge-Seawater. Adsorption time - 1 day (top), 4 day (mid), 7 day (bottom). m = 114,570,1140,3400 mg/L

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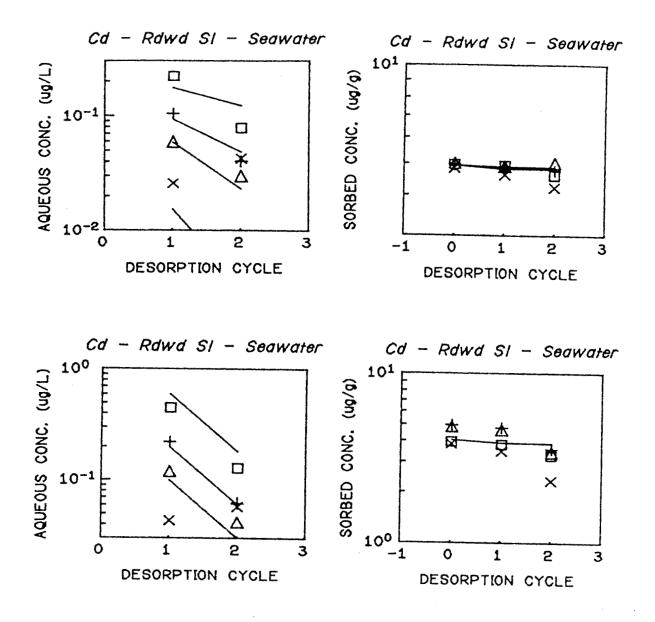


Fig. 4. Effect of Adsorption Time. Cadmium-Ridgewood Sludge-Seawater. Adsorption time = 18 day (top), 25 day (bottom), m = 114,570, 1140,3400 mg/L

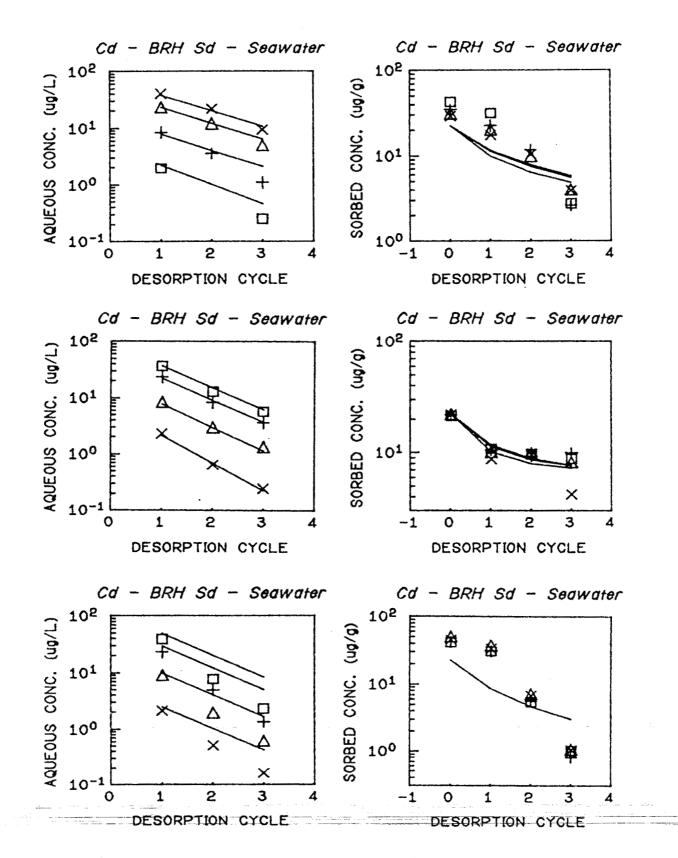


Fig. 5. Effect of Adsorption Time. Cadmium-Black Rock Harbor Sediment-Seawater. Adsorption time = 11 day (top), 18 day (mid), 25 day (bottom). m = 170,690,2080,3470 mg/L.

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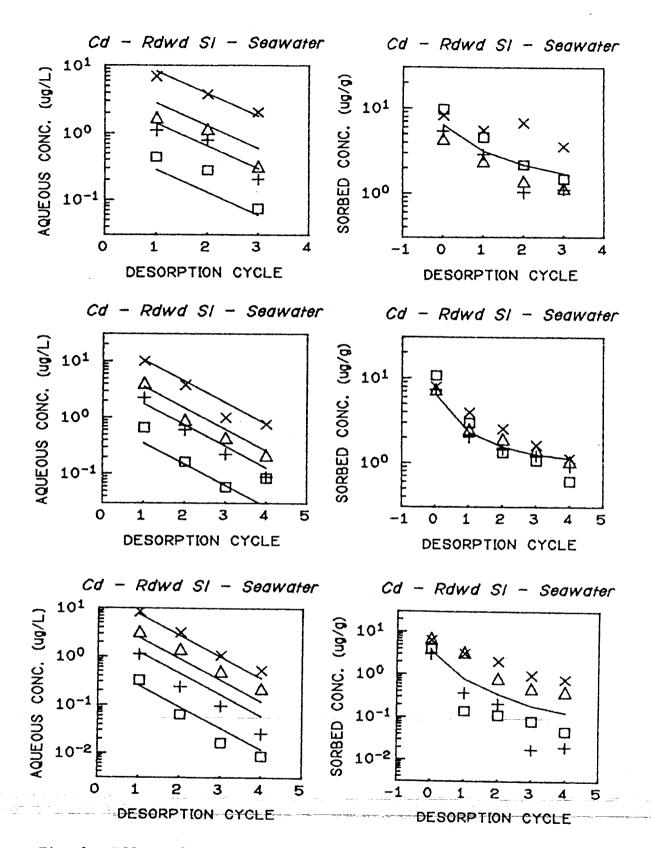


Fig. 6. Effect of Desorption Time. Cadmium-Ridgewood Sludge-Seawater. Desorption time = 1 hr. (top), 5 day (mid), 7 day (bottom). m = 85,425,850,2550 mg/L.

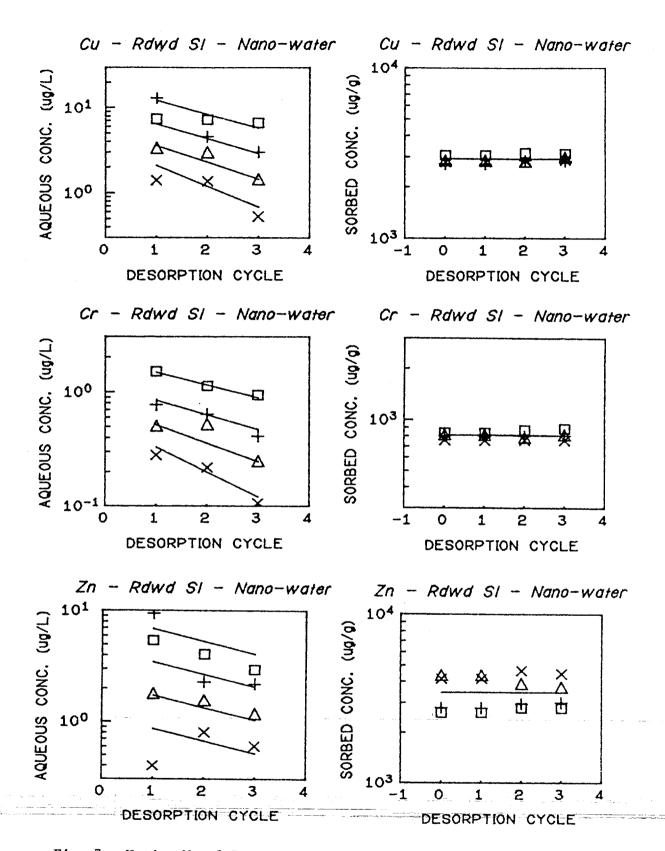


Fig. 7. Native Metal Desorption. Ridgewood Sludge-Nanopure Water. Desorption time = 1 hr. Copper (top) Chromium (mid), Zinc (bottom). m = 340,680,1370,2730 mg/L.

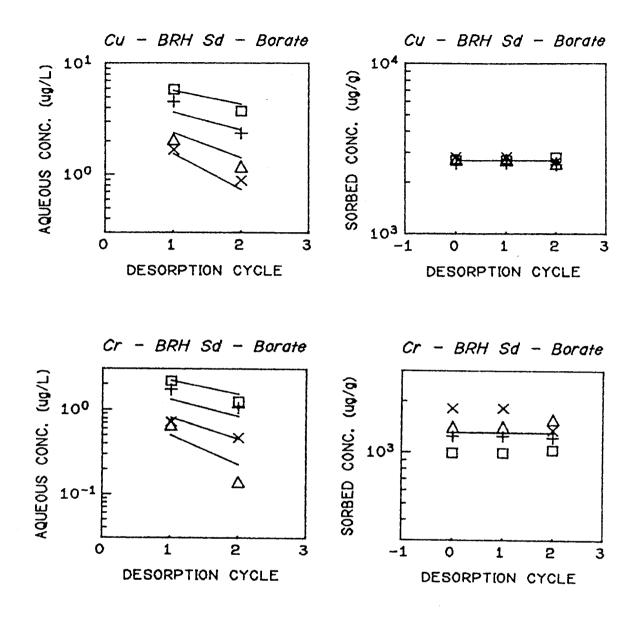


Fig. 8. Native Metal Desorption. Black Rock Harbor Sediment-Borate Buffer. Desorption time = 1 hr. Copper (top) Chromium (bottom). m = 340, 680,1370,2730 mg/L.

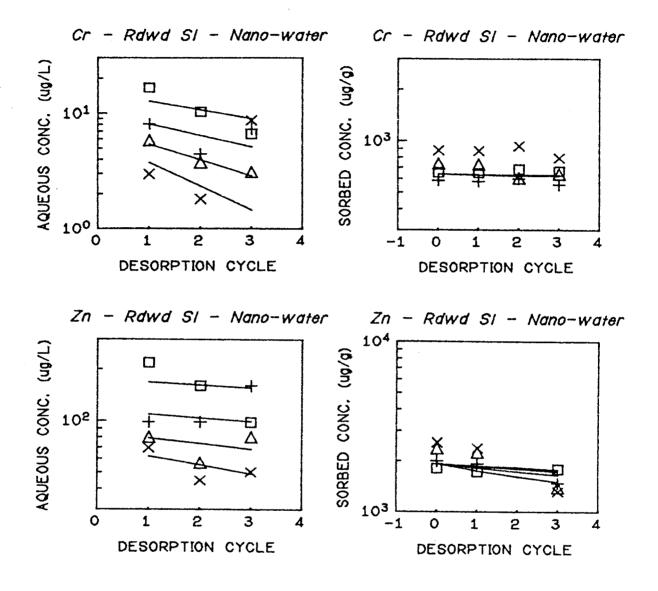


Fig. 9. Native Metal Desorption. Ridgewood Sludge-Nanopure Water. Desorption time = 1 day. Copper (top) Chromium (mid), Zinc (bottom). m = 340,680,1370,2730 mg/L.

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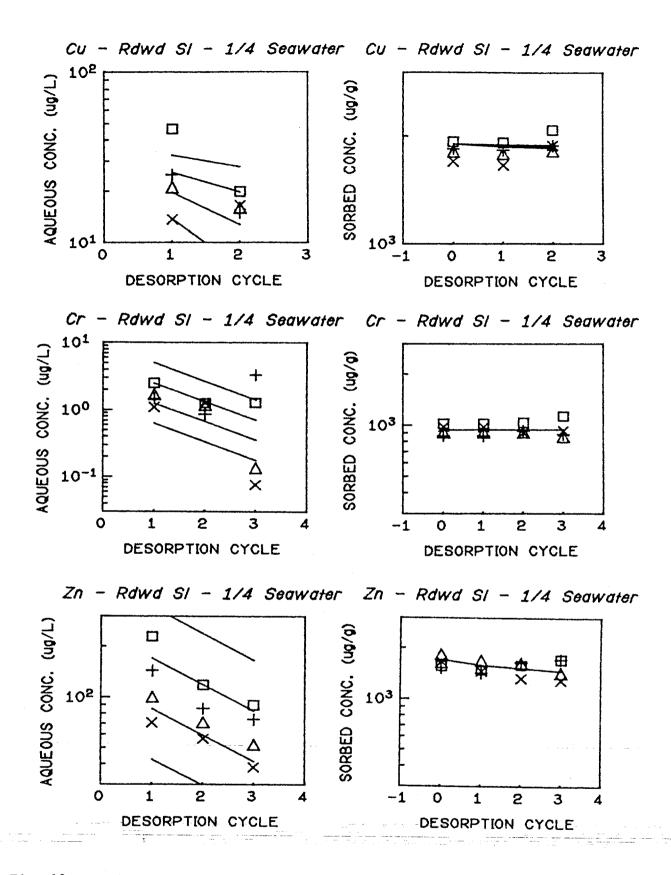


Fig. 10. Native Metal Desorption. Ridgewood Sludge-1/4 Seawater, 3/4 Nanopure Water. Desorption time = 1 hr. Copper (top), Chromium (mid), Zinc (bottom). m = 340,680,1370,2730 mg/L.

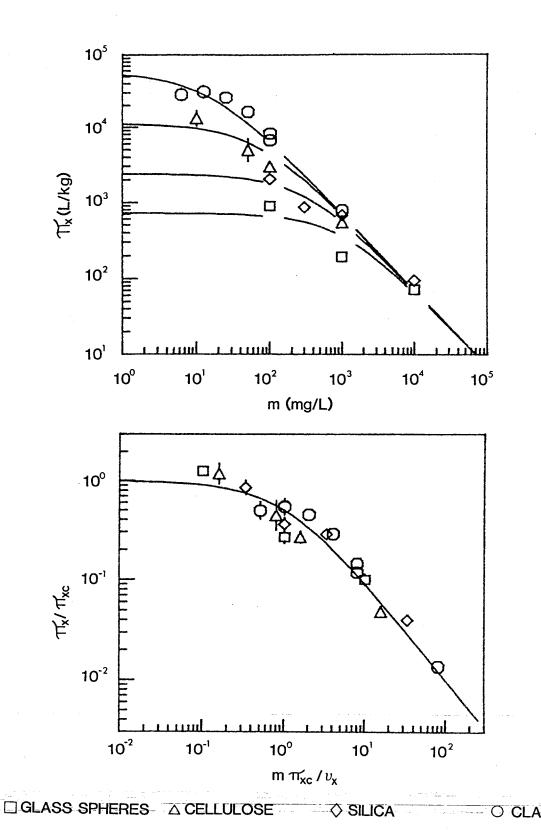
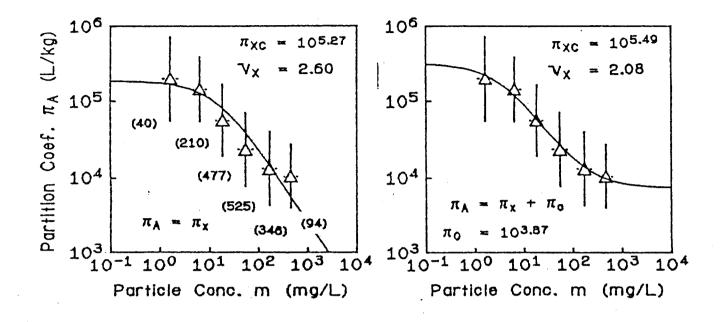


Fig. 12. Reversible Component Partitioning.  $\pi_x$  versus m (top). Normalized Partition Coefficient versus normalized particle concentration (bottom).

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Ni Partition Coefficient STORET DATA - Streams

# Effect of Particle Conc.



# Effect of Chemical Variables

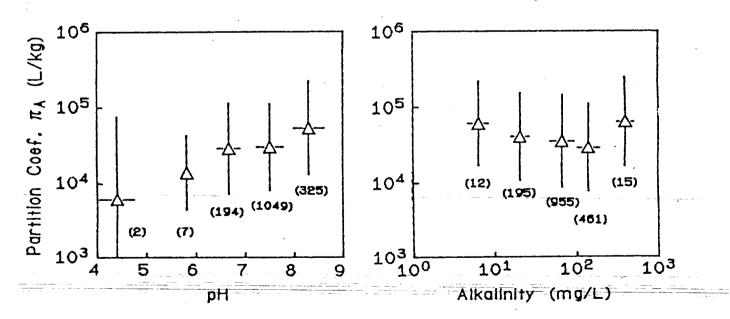
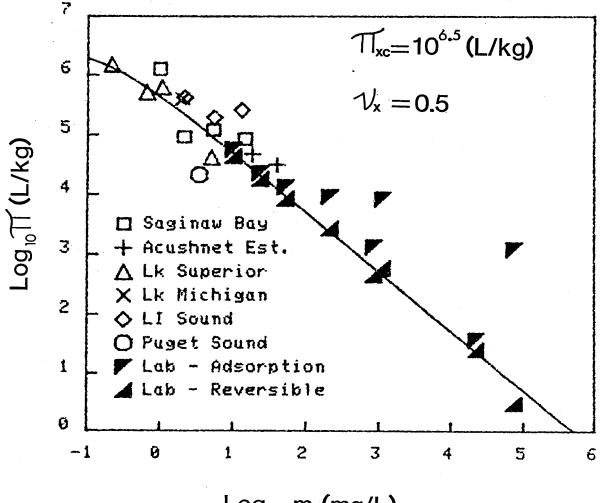


Fig. 13. Nickel Distribution Coefficient versus Particle Concentration. STORET stream data, log mean ± log std. deviation No. of data point indicated.

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 $Log_{10}$  m (mg/L)

Fig. 14. PCB Distribution Coefficient versus Particle Concentration.

### APPENDIX 1

## The Effects of Nonreversibility, Particle Concentration, and Ionic Strength on Heavy Metal Sorption

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#### Introduction

The extent of heavy metal sorption is a critical component in the evaluation of their fate in natural water systems. Adsorption at trace concentrations is usually quantified using a linear isotherm parameterized by a distribution or partition coefficient,  $\pi$ , which is the slope of the isotherm. It is well established that the magnitude of the partition coefficient is a function of particle properties such as particle size (1), the nature of the oxide coating (2), organic carbon content (3), zero point of charge (4); and of chemical variables such as pH (5), ionic strength (6), and the concentration of complexing ligands (7,8). Such effects can usually be rationalized within the frameworks of solution and surface chemistry.

Less well understood is the desorption reaction. The conventional desorption tests for heavy metals are extractions using concentrated solutions of various ionic compositions (9-10) rather than adsorbate-free aqueous phase. For example after adsorption in freshwater, mixtures of freshwater and seawater are used to simulate the desorption that might occur at freshwater-estuarine transition zones (11). This is in contrast to the usual desorption tests for organic chemicals in which the same aqueous phase is used in both the adsorption and desorption steps of the experiment and which provides a direct test of reversibility.

Heavy metal desorption into various extraction solutions has been found to be incomplete, suggesting that the reaction is not completely or even moderately reversible (e.g. 12). The usual finding is that substantial quantities of the heavy metals remain associated with the particle even at extractant concentrations ( $\sim 0.1M$ ) which should displace all the physically

adsorbed metal (see Forstner and Wittman, 13, for a review and other references). A review of the applicable desorption literature (11-22) indicates that there are almost no published heavy metal data which include adsorption and desorption isotherms using the same aqueous phase that would be suitable for testing the extent of reversibility in environmentally realistic aqueous phases.

Data for nickel and cobalt montmorillonite sorption are presented below which demonstrate that reversibility is not complete and varies as a function of chemical (pH and ionic strength) and physical (particle concentration) variables. A method for analyzing adsorption-single desorption isotherms which yields the partition coefficients of the reversible and resistant components (23) is applied in order to isolate the behavior of the reversibly sorbed heavy metal.

Perhaps the most puzzling phenomenon exhibited by heavy metal and organic chemical partitioning is the effect of the particle concentration itself. Of course the extent of sorption, that is the fraction of chemical in the sorbed state, increases with increasing particle concentration. However it has been noted that the partition coefficient tends to decrease with increasing particle concentration (24), in some cases quite dramatically. It is as if, for the analogous partitioning in gas-water systems, the Henry's constant is dependent upon the volume of the gas phase. Since such behavior has not been noted for other phase distribution coefficients a certain amount of skepticism attends these observations in particle-water systems. Although the effects of particle concentration on the partition coefficient are not infrequently noted for various heavy metal-sorbent systems in a wide variety of settings (25,26,27) they are usually explained as either experimental artifacts, or as the result of an improper chemical system descrip-

tion which ignores, for example, the presence of other important but unobserved phases such as colloidal particles or dissolved organic carbon (28).

We present below the results of simple resuspension and dilution experiments that strongly suggest the effect is due to particle interactions and not experimental artifacts or hidden variable fluctuations. An equation derived from this hypothesis describes the experimental data from a number of sorbate-sorbent systems. The model suggests that the classically conceived partition coefficient is indeed descriptive of reversible sorption at low particle concentrations but that as particle concentration is increased above a certain concentration, which is chemical and sorbent specific, the reversible partition coefficient starts to decrease and its magnitude is dependent only on particle concentration and independent of the specific sorbate-sorbent or auxiliary aqueous phase parameters such as ionic strength. It is suggested that physical rather than chemical factors are controlling in this regime.

We begin by examining the effects of the usual experimental variables, sorption time and pH, and then proceed to the particle concentration effects. Additional experiments at varying ionic strengths provide further insight into the influence of particle concentration.

#### Methods and Materials

Initial experiments indicated that buffered aqueous phases are necessary since pH cannot be controlled in the usual fashion by addition of acid or base to the reactors during equilibration. The reasons are that a-desorption step must follow the adsorption equilibration which requires that

the vessel be a centrifuge tube making pH adjustments difficult and requiring multiple vessels for an isotherm. The inconvenience of this procedure prompted the use of buffers.

Both inorganic (borate, carbonate) and organic (sulphonic acid) buffers that were found to be noncomplexing with respect to nickel were used for the experiments. Differential Pulse Polarographic analysis was employed to measure the presence of any complexation between the metal and these buffers. No change in the half wave potential for the metal ion alone compared with that for the metal ion in the presence of the buffer was observed. It is significant to note that this method is somewhat more sensitive than those previously used for this purpose (29). In fact we have measured metal binding constants for systems previously reported to be negligible, e.g. for copper and EPPS (29),  $K_f = 10^3$ . Finally the similar adsorption-desorption results obtained using quite different buffers suggest that complexing is not occurring to a significant extent. This is discussed in detail below.

Adsorption-desorption isotherms were obtained by the following procedure. The aqueous phase, containing buffered distilled water and metal of known concentration tagged with  $^{65}$ Ni or  $^{60}$ Co, was combined with montmorillonite or quartz particles to a volume of 9 ml in a 20 ml centrifuge tube. The sorbents employed were montmorillonite (Wards Scientific) and 5 µm  $\alpha$ -quartz particles ( Min-U-Sil 5, Penn. Glass Sand Corp.). The former was chosen for its high sorptive capacity and the latter for its uniform particle size distribution and well understood surface chemistry. The initial metal concentration in the systems ranged from 200 to 3.3 µg/L; the sediment concentrations studied were 5, 10, 30, 100, 300, and 1000 mg/L for montmorillonite and 200-10,000 mg/L for quartz. The buffers were kept at the mini-

mum effective level to maintain pH during the course of the experiment  $(10^{-3}$ M). The following buffers were employed: bicarbonate (pH 7), TAPS (29)(pH 8), borate (pH 8-10), and carbonate (pH 10-11). The capped tubes were agitated until adsorption equilibrium was reached. A sample of the sediment-aqueous phase mixture is removed and the total metal concentration,  $c_{Ta}$ , was determined by liquid scintillation counting. The system was then centrifuged (6000-12,000g for 15-20 min) and the metal concentration of the aqueous phase,  $c_a$ , was similarly determined. The sorbed concentration at adsorption equilibrium is calculated by difference:  $r_a = (c_{Ta} - c_a)/m$ , where m is the adsorbent concentration.

For desorption steps, the contaminated aqueous phase is carefully removed leaving the sedimented solids in the tube, and uncontaminated aqueous phase is added to achieve a total volume that produces the same particle concentration as initially present at adsorption. The capped tube is agitated until desorption equilibrium is achieved. A sample of sediment-aqueous phase mixture is removed and analyzed yielding  $c_{Td}$ . After centrifugation an aqueous phase sample is removed and analyzed for  $c_d$ . The sorbed concentration is obtained by difference,  $r_d = (c_{Td} - c_d)/m$ .

In order to check our experimental procedures we have examined the behavior of particle-free but otherwise identically configured experiments and found no evidence of precipitation. No adsorption-desorption differences were found if preceding the adsorption step, a particle agitation, centrifugation, and resuspension are performed. Thus the particle separation is efficient in the sense that only an insignificant fraction of sorbing particles are not removed by the centrifugation.

#### Analysis Framework

The two component model of adsorption-desorption is designed to interpret incompletely reversible behavior illustrated in Fig. 1 as linear and logarithmetic plots of isotherm data. Sorbed chemical at adsorption equilibrium, r, is assumed to be the sum of a reversibly sorbed component described by a linear isotherm:  $r_x = \pi_x c$ , and a resistant component that adsorbs following a linear adsorption isotherm:  $r_0 = \pi_0 c_0$  but resists desorption over the time periods characteristic of experimental desorption equilibration (hours to days). Fig. 2 illustrates the model framework. From the assumption of two components it follows that for adsorption,  $r_a = (\pi_x + \pi_x)$  $\pi_0$ )  $c_a = \pi_a c_a$  so that the conventional adsorption partition coefficient,  $\pi_a$ , is the sum of the reversible and resistant component partition coefficients. During desorption, only the reversible component interacts with the solution phase following linear desorption isotherms, all with slope  $\pi_x$ . For each pair of adsorption, r, c, and desorption, r, c, concentrations it is possible to calculate the resistant component concentration by extrapolating the consecutive desorption isotherm to zero aqueous concentration:

$$\mathbf{r}_{0} = (\mathbf{r}_{d} - \beta \mathbf{r}_{a})/(1 - \beta) \tag{1}$$

where  $\beta = c_d/c_a$  and to calculate, by difference, the reversible component concentrations  $r_{xa} = r_a - r_o$  and  $r_{xd} = r_d - r_o$  at adsorption and desorption equilibrium respectively. The lower panels of fig. 2 illustrate an example of the resulting component concentrations calculated in this way and their conformity to linear isotherms. The component partition coefficients can also be calculated from the adsorption,  $\pi_a$ , and single desorption "partition coefficient",  $\pi_d$ , using the relationships (23):

$$\pi_{x} = \frac{\pi_{a}}{1 + m(\pi_{d} - \pi_{a})}$$
(2)

$$\pi_{o} = \pi_{a} - \pi_{x} \tag{3}$$

where  $\pi_a$  and  $\pi_d$  are the slopes of the linear isotherms (fig. 1). These equations require that no adsorption or desorption occurs from the vessel during desorption. The mass balance requirement is:

$$mr_{a} = c_{d} + mr_{d} = c_{Td}$$
<sup>(4)</sup>

where m is the particle concentration. Fig. 3 presents the comparison for the data to be used subsequently. The histogram of the relative error = 100% (mr<sub>a</sub> - c<sub>Td</sub>)/c<sub>Td</sub> illustrates conformity to this assumption within ±20%.

For each experimental condition a complete adsorption and desorption isotherm is generated, 4 to 8 duplicated concentrations over the range 3.5-200 µg/L, from which the partition coefficients are determined as the log average of the ratios:  $\pi_a = r_a/c_a$ ,  $\pi_d = r_d/c_d$ . Equations (2-3) are used to compute  $\pi_x$  and  $\pi_o$ . The experimental log standard deviations are used to compute the log standard errors of  $\pi_a$ ,  $\pi_x$  and  $\pi_o$  which are presented in the subsequent figures as error bars if they exceed the symbol size.

#### Kinetics

The kinetics of metal sorption are usually found to be rapid (30). For the nickel-montmorillonite-borate system both adsorption and desorption kinetics were examined by generating isotherms at varying adsorption,  $t_{ads}$ , and desorption,  $t_{des}$ , times.

The partition coefficients for varying adsorption times are shown infig. 4. No significant change in  $\pi_a$  is observed for adsorption times of up to 5 days. Some suggestion of increasing irreversibility as  $t_{ads}$  increases can be seen from the decreasing  $\pi_x$  and the increasing ratio  $\pi_o^* = \pi_o/\pi_a$ which is the resistant component fraction:  $r_o/r_a$ , but the trend is not consistent.

A more important question is the effect of increasing desorption time,  $t_{des}$ , since incomplete reversibility may be the result of insufficient desorption time. Fig. 4 illustrates the results of increasing desorption time after a fixed  $t_{ads} = 1$  hr. adsorption. The adsorption partition coefficient,  $\pi_a$ , variations are due to experimental fluctuations of the replicate isotherms. No significant change in  $\pi_x$  or  $\pi_o^*$  is seen to occur for up to  $t_{des}$ = 5 days. Hence the component concentrations seem to be in metastable equilibrium for time scales of up to 5 days. In summary no substantial kinetic effects were observed with the possible exception of an increase in the resistant component partition coefficient at longer adsorption times. At long desorption times incomplete reversibility persisted.

### Effect of pH and Buffering

The strong effect of pH on metal adsorption is well known and the subject of numerous experiments and models (5,6,31,32). The usual focus is the shape of the curve of percent metal sorbed versus pH and, in particular, the pH at which adsorption begins. Since sorption is only important at a pH above this "pH edge", we concentrated at pH valuess above the edge and at low montmorillonite concentrations, m = 100 mg/l. Full adsorption-desorption isotherms were generated using different bufters in order to detect possible complexing effects. The results, shown in fig. 5, suggest a slightdecrease in  $\pi_a$  and  $\pi_x$ , together with an increase in resistant component frac-

tion,  $\pi_0^{\pi}$ , at pH  $\sim$  9.5 which corresponds to the pK of Ni(OH)<sup>+</sup> (33). Less extensive cobalt data exhibit the reverse trend with increasing  $\pi_a$  and  $\pi_x$  between pH 8 and 9 and no significant change in resistant fraction.

In summary, no major pH effects were noted for either nickel or cobalt. Significant (20-60%) nonreversibility persisted throughout the pH range investigated with no apparent complexing effects due to the buffers.

# Particle Concentration Effects: Isotherms, Resuspension and Dilution Experiments

In contrast to the effects of kinetics and pH, a significant relationship exists between partition coefficients and montmorillonite or quartz concentration. Adsorption isotherms at pH = 9 (borate) for Ni-clay at m = 30 - 1000 mg/L, cobalt-clay for m = 100 and 1000 mg/L, and Ni-quartz for m = 600 and 10,000 mg/L, fig. 6, clearly show the effect of varying particle concentration.

The inverse relationship found between particle concentration and partition coefficient is certainly a most unexpected result and additional experiments were performed in order to rule out any artifacts. The most popular explanation of this phenomenon is to invoke an additional complexing agent associated with the particles which experimentally appears to be dissolved – it is not removed by the particle separation technique – but actually complexes (if it is a ligand such as dissolved organic carbon (28)) or sorbs (if it is colloidal particles (27)) some dissolved chemical. If it is assumed that the concentration of the third phase increases with the concentration of added particles then in fact an apparent reduction in partition coefficient is predicted with increasing particle concentration since the experimentally determined "dissolved" concentration is actually the sum of

the truly dissolved, and the additional sorbed or complexed chemical.

The experiments described below were designed to specifically exclude an additional complexing agent as an explanation. The design was based upon the desire to change only the particle concentration while maintaining all else constant. Two experimental designs were developed which are in a sense the inverse of each other: a resuspension procedure and a dilution procedure.

For the resuspension experiment (fig. 7), replicate controls and experimental vessels are prepared. After adsorption equilibrium is achieved, the particles and supernatant are separated by centrifugation and, only in the experimental vessels, a portion of the supernatant is removed. The volume of supernatant in the control vessels is not changed. Then the particles are resuspended into the supernatant and equilibrated, after which measurements of the total and dissolved nickel are made. In both the experimental and control vessels the particles are resuspended into the same supernatant with which they were in equilibrium. If a "dissolved" third phase is present which is not separated by centrifugation then its <u>concentration</u> is not changed by the partial removal of supernatant in the experimental vessels. The only modification is that the particle concentration is increased in the experimental vessels since the supernatant volume is decreased.

The results for both the controls and the experimental vessels are shown in Table 1 for montmorillonite and quartz particles. The montmorillonite controls behaved as expected: the dissolved nickel concentration at both adsorption and resuspension equilibrium are essentially equal and their ratio is unity. For the quartz experiments a slight decrease occurred in the controls presumably due to lack of complete equilibrium. The experimen-

tal vessels for all experiments clearly exhibit the particle concentration effect. The result of increasing particle concentration is apparently to decrease the partition coefficient. This causes nickel to desorb from the particles thereby increasing the dissolved nickel concentration. The average ratio of resuspension to adsorption dissolved concentrations increases in all cases whereas the average ratio for the controls is either equal to or less than one.

The dilution experiment is designed to decrease the particle concentration without any other alteration. Three parallel sets of vessels are employed: a control, an experimental vessel, and a vessel to supply additional supernatant with which to dilute the experimental vessel. After adsorption equilibrium, the supernatant from the third vessel is obtained by centrifugation and this particle-free supernatant is added to the experimental vessel. If no experimental fluctuations were present this supernatant should have the same dissolved nickel concentration as the experimental and control vessels since they are identically configured. Note that the experimental vessel is not centrifuged. The particle dilution is achieved by adding additional (identical) supernatant. The vessels are then equilibrated and the total and aqueous concentrations are measured and a comparison is made between the controls and experimental aqueous concentrations. Again if a "dissolved" third phase existed it would have the same concentration in both the experimental and control vessels. As shown in Table 2, the aqueous concentrations in the experimental vessels decrease as predicted. The particle concentration is decreased by dilution, so that the partition coefficient increases, and more nickel is removed from solution. In all the resuspension and dilution experiments, pH fluctuations were not significant ( $\Delta pH < 0.3$ ).

Both of these experiments confirm the particle concentration dependence of the reversible partition coefficient for increasing and decreasing particle concentrations. The resuspension experiment involves centrifuging the experimental vessels but the controls for the montmorillonite experiment indicate that centrifuging and resuspension into the same supernatant volume causes no change in aqueous concentration. The only possible remaining cause is the changed particle concentration. The dilution experiment dispenses with any particle separation in the experimental vessels. The only possible experimental source of variation is that the added supernatant is at a slightly different aqueous concentration than that in the experimental vessel. A comparison of the control and supernatant aqueous concentrations do suggest that such fluctuations are present but the experimental aqueous concentrations are lower than either in all cases, thus confirming the particle concentration effect for lowered particle concentrations.

Both these experimental designs specifically exclude the possibility of a "dissolved" third phase since if it were present its <u>concentration</u> would be <u>the same</u> in both the experimental and control vessels so that its presence could not affect the difference in partitioning observed at differing particle concentrations. The hypothetical third phase could also be a competing cation. Decrease in the adsorption partitioning of lead at increasing Ca<sup>++</sup> saturated montmorillonite concentrations has been explained using a competitive Langmuir isotherm with Ca<sup>++</sup> as the competing cation (34). However, these experiments were conventional adsorption isotherms at various (large) clay concentrations (1-20 g/L) so that the additional clay could supply the additional cation. The design of the resuspension experiment precludes this explanation. In fact, any hypothesis which attempts to ex-

plain the results via an increased desorption of complexing ligands or an increased adsorption of competing cations just transfers the problem to explaining why the changes in ligand or cation concentrations occur with increasing particle concentration for an aqueous phase that is initially in equilibrium with the particles.

The third phase changes are specifically excluded as an explanation of the results of the resuspension and dilution experiments. Similar results have been observed for hexachlorobiphenyl - sediment and clay experiments (35) although those experiments did not include parallel controls.

Another explanation for the particle concentration effect is flocculation or at least some agglomeration of the sediment particles. This may occur initially or at some later point in the experimental procedure. The result of such a process would be to decrease the overall surface area of the adsorbent, and thus reduce the number of sites available for metal binding leading to a reduction in the value of the partition coefficient in higher sediment-concentration systems.

Extensive microscopy of representative clay and quartz suspensions including the resuspension experiments was carried out. In all cases, both the size and distribution of particles was essentially constant. These results clearly rule out particle size alteration as a means of accounting for the particle concentration effect in our experiments.

## Particle Concentration Effects - A Particle Interaction Model

The resuspension and dilution experiments strongly suggest that it is the presence of particles themselves, and not a surrogate variable, that is influencing the extent of reversible partitioning. Thus it appears reasonable to expect that particle interactions somehow affect sorption. We tentatively suggest the following model for the sorption of the reversible component:

$$\frac{K_{ads}}{Adsorption: c + m \rightarrow c=m}$$
(5)

Spontaneous Desorption: 
$$c=m \rightarrow c + m$$
 (6)

$$\frac{K_{p-d}}{Particle Interaction} \qquad c=m+m \rightarrow c+2m \qquad (7)$$

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where c is the dissolved chemical concentration, c=m is the reversibly sorbed chemical concentration and m is the particle concentration. The adsorption and spontaneous desorption mechanisms are the conventional representations for reversible partitioning. Particle interaction induced desorption assumes that a binary interaction between particles causes desorption to occur. The mechanism that prompts this desorption is uncertain at present.

For these three reactions at equilibrium, the reversible partition coefficient is:

(8)

$$\pi_{x} = \frac{[c=m]}{[c][m]} = \frac{K_{ads}}{K_{s-d} + mK_{p-d}}$$

A useful normalization of this expression is:

$$\pi_{x} = \frac{\pi_{xc}}{1 + (m \pi_{xc} / \nu_{x})}$$
(9)

where  $\pi_{xc} = K_{ads}/K_{s-d}$ , the classical partition coefficient and  $v = K_{ads}/K_{p-d}$ , which is the slope of the  $\pi_{x}$  versus m relationship at large particle concentration:

$$\pi_{x} \approx \frac{\nu_{x}}{m} \qquad m > \nu_{x}/\pi_{xc} \qquad (10)$$

A convenient parameter that quantifies the particle concentration at which particle interactions become important is the ratio  $K_m = v_x/\pi_{xc}$ , which can be thought of as the half saturation constant for the  $\pi_x - m$  relationship. That is at  $m = K_m$ ,  $\pi_x = \pi_{xc}/2$  so that  $K_m$  is the particle concentration at which the particle interaction desorption and spontaneous desorption are of equal magnitude:  $K_{s-d} = m K_{p-d}$ . Fig. 8 presents Ni - montmorillonite data and the  $\pi_x$  relationship, eq. (9). For a tentative assignment of  $\pi_o^* = 0.4$ , which assumes that the resistant fraction is independent of m - the data are too scattered for a more refined hypothesis - the resulting  $\pi_a = \pi_x + \pi_o$ relationship is compared to observation in fig. 8.

These  $\pi_x$  data and those for the Co-montmorillonite and Ni - quartz are fit to eq. (9) using a nonlinear least squares fit of log  $\pi_x$  versus log m to estimate  $\pi_{xc}$  and  $\nu_x$ . The results are given in Table 3 (Individual Data Sets). The resulting  $\nu_x$  are essentially equal for the three systems. A four parameter least square fit: constant  $\nu_x$  and an individual  $\pi_{xc}$  for each system yields the same result within the standard error estimates (Table 3 -Joint Data Set). The results are shown in fig. 9 for the three systems. A

$$\frac{\pi x}{\pi xc} = \frac{1}{1 + m/K_{m}}$$
(11)

illustrates that these data are well represented by eq. (11) and that each system has data in both the range of m for which  $\pi_x$  is essentially constant:  $\pi_x = \pi_{xc}$  (m << K<sub>m</sub>) and for which  $\pi_x$  varies inversely with particle concentration:  $\pi_x = v_x/m$  (m >> K<sub>m</sub>).

Two observations are worthy of note. For  $m \ll K_m$ ,  $\pi_x \approx \pi_{xc} = K_{ads}/K_{s-d}$ , the classical reversible partition coefficient. It is a function of both the nature of the sorbate (Ni-clay:  $\pi_{xc} = 45$  L/g; Co-clay:  $\pi_{xc} = 5.8$  L/g) and the sorbent (compare Ni-clay to Ni-quartz:  $\pi_{xc} = 1.1$  L/g) as would be expected. The surprising result is that  $\nu_x$  is essentially independent of both sorbate type and sorbent properties (Table 3 - Individual Data Sets). Since:

$$v_{\rm x} = \frac{K_{\rm ads}}{K_{\rm p-d}}$$
(12)

is the ratio of adsorption to particle induced desorption rate, the fact that it is constant suggests that its magnitude is physically rather than chemically determined. This would be the case if, for example, both the adsorption rate and particle interaction induced desorption rate were functions of hydrodynamic factors.

Following this idea and assuming that K<sub>ads</sub> is constant, the extent of reversible partitioning is determined by the magnitude of the two desorption mechanisms. The data in fig. 9 can be interpreted as follows. For small particle concentrations, the rate of spontaneous desorption exceeds the rate of particle interaction induced desorption (which is small due to the low

particle concentrations) and spontaneous desorption dominates the total desorption rate. Spontaneous desorption is expected to be a function of the strength of the chemical and electrostatic bonding and so it is chemical and surface specific. The spontaneous desorption rate is smallest for Ni-clay, yielding the highest reversible partition coefficient at small m, in decreasing order followed by Co-clay and Ni-quartz (left hand side of fig. 9). However as m increases, the rate of particle interaction induced desorption increases for all systems. Since the spontaneous desorption rate is smallest tor Ni-clay, and as the particle interaction desorption rate, m  $K_{d-n}$ , becomes comparable to  $K_{s-d}$  for this system first, the Ni-clay reversible partition coefficient starts to decrease at relatively small particle concentrations  $K_m = 18.7 \text{ mg/L}$ . However, the partitioning of Co-clay and Ni-quartz systems, with their larger spontaneous desorption rates, are unaffected at these particle concentrations. As m increases further, however, m  $K_{n-d}$  approaches  $K_{s-d}$  for Co-clay ( $K_m = 146 \text{ mg/L}$ ) and  $\pi_x$  for this system starts to decrease. Not until m approaches  $K_m = 786 \text{ mg/L}$  for the Ni-quartz system is m  $K_{p-d}$ comparable to  $K_{s-d}$  for this system.

Hence  $\pi_x$  is a specific chemically determined constant,  $\pi_{xc}$ , for each system until m reaches  $K_m = K_{s-d}/K_{p-d}$  for that system at which point m  $K_{p-d}$ =  $K_{s-d}$  and the additional particle interaction induced desorption causes  $\pi_x$ to start decreasing. For large m >>  $K_m$ , the spontaneous desorption rate is overwhelmed by the particle interaction induced desorption and

$$\pi_{x} \cong \frac{K_{ads}}{mK_{p-d}} \stackrel{\nu}{=} \frac{x}{m} \qquad m \gg K_{m} \qquad (13)$$

which is constant for the all sorbate-sorbent pairs. Hence all  $\pi_x$  curves eventually meet at high particle concentrations as shown in fig. 9. Note

that at m = 1000 mg/L the three systems have virtually the same  $\frac{\pi}{x}$  and the pattern continues for Ni-quartz at m = 10,000 mg/L. This is a consequence of the virtually constant  $\frac{\nu}{x}$  = 0.68 - 0.91 found experimentally.

#### Ionic Strength Effects

Metal-clay sorption is known to decrease significantly as ionic strength increases (36,37). We examine this effect at two particle concentrations (m = 100 and 1000 mg/L) using adsorption and desorption isotherms as before (pH = 9, borate buffer). The component partition coefficients are shown in fig. 10. At the low ionic strength of the buffer, the effect of particle concentration is clearly in evidence. However, as ionic strength is increased by adding NaNO<sub>3</sub>, the  $\pi_a$  values at both particle concentrations tend to become equal. The reversible component  $\pi_x$ 's are approaching each other, whereas the resistant component is behaving quite differently. For m = 100 mg/L the system is essentially approaching complete reversibility as ionic strength increases whereas for m = 1000 mg/L the resistant component persists as  $^{-1}$ 50% of the total sorbed nickel with some decrease apparent at high ionic strength.

We offer no explanation for the resistant component behavior. However, the reversible component can be analyzed using the particle interaction desorption model. Using the observed  $\pi_x$  at m = 100 and 1000 mg/L, and eq. (8) in the form:

$$\frac{1}{\pi} = \frac{K_{s-d}}{K_{ads}} + \frac{K_{p-d}}{K_{ads}} m$$
(14)

the ratios:  $\pi_{xc} = K_{ads}/K_{s-d}$  and  $v_{x} = K_{ads}/K_{p-d}$  can be computed at each ionic strength. They are shown in fig. 11 together with the Co-clay data, and Ni-

quartz data at the ionic strength of the buffer. The results of additional Ni-quartz experiments at I = 0.01M, pH = 9, borate buffer m = 100 - 10,000 mg/L, fit using eq. (9) to obtain  $\pi_{xc}$  and  $\nu_x$  are also included. Once again the influence of the chemical factors: ionic strength and the specific sorbate-sorbent pair, are clearly seen for the estimated classical partition coefficient,  $\pi_{xc}$ , =  $K_{ads}/K_{s-d}$ . If we assume that  $K_{ads}$  is constant and not chemically controlled, then the decrease in  $\pi_{xc}$  is due to an increase in the spontaneous desorption rate,  $K_{s-d}$ . The Ni-clay data suggest that  $K_{s-d} \sim 1^{\frac{1}{2}}$  which is the same ionic strength power law dependency for the double layer thickness (38) suggesting that the reduction in double layer thickness as I increases lowers the energy barrier and enhances spontaneous desorption. The use of more refined chemical models (31,32) to explain these relation-ships as well as the observed pH variations (fig. 5) is clearly possible.

The point is that the experimental data should be collected using adsorption-desorption isotherms at two particle concentrations so that  $\pi_{xc}$  can be estimated. The formulation of equilibrium sorption models assumes both complete reversibility and no particle interaction effects. Thus the models cannot be expected to apply to simply the adsorption partition coefficient,  $\pi_a$ , which is confounded by both non-reversibility:  $\pi_a = \pi_x + \pi_o$ , and by  $\pi_x$ itself varying with particle concentration (fig. 10). It is only  $\pi_{xc}$ , the classical low particle concentration limiting partition coefficient, that satisfies the implicit model assumptions of reversibility and no particle interaction effects.

On the other hand, for large particle concentrations (m >>  $K_m$ ),  $\pi_x = v_x/m$  and the estimates of  $v_x$  are essentially constant and independent of x ionic strength and the identity of the sorbate-sorbent pair. This is

= 0.3 (39), suggests that this finding is not just a laboratory curiosity but has wider applicability.

If in fact  $v_x$  is essentially constant and of order one, then the half saturation constant for the particle concentration effect becomes:

$$K_{m} = \frac{v_{x}}{\pi} \sim \frac{1}{\pi}$$

which suggests that highly sorbed chemicals (large  $\pi_{xc}$  say  $10^6$  L/kg) will exhibit particle concentration effects at low particle concentrations (m = 1.0 mg/L) whereas less highly sorbed chemicals ( $\pi_{xc}$  =  $10^3$  L/kg) are less likely to exhibit the effect until larger particle concentrations (m = 1000 mg/L) are encountered.

#### Acknowledgements

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## Table 1

## Resuspension Experiment Borate Buffer, pH = 9 Aqueous Ni Concentrations (µg/l)

# Montmorillonite

Cor	ntrol Vessels		Experi	mental Vessels	
Adsorption	Resuspension	Ratio	Adsorption	Resuspension	Ratio
$m = 100 mg/\ell$	$m = 100 mg/\ell$	· · · · · ·	$m = 100 mg/\ell$	m = 243 mg/l	
c <sub>a</sub>	c <sub>rs</sub>	c <sub>rs</sub> /c <sub>a</sub>	с <sub>а</sub>	<sup>c</sup> rs	c <sub>rs</sub> /c <sub>a</sub>
72.6	72.2	1.00	74.7	97.3	1.30
38.9	37.8	0.97	36.8	40.4	1.10
15.9	17.0	1.07	19.1	23.4	1.23
6.18	6.33	1.02	5.32	5.87	1.10
	Average ratio	= 1.017		Average ratio	= 1.182

# Quartz

m = 1000 mg/	L m = 1000 mg	/L	m = 1000 mg/	L m = 2430 n	mg/L
c a	c <sub>rs</sub>	c <sub>rs</sub> /c <sub>a</sub>	с а	°rs	c <sub>rs</sub> /c <sub>a</sub>
102.	91.9	0.901	87.3	96.0	1.10
46.8	41.6	0.889	38.3	43.7	1.41
20.1	19.0	0.945	19.7	20.4	1.04
10.1	9.32	0.923	9.80	11.9	1.21

Average ratio = 0.914

Average ratio = 1.123

Quartz

$m = 1000 m_{\odot}$	g/L m = 1000 mg	/L	m = 1000 mg/	L m = 3750 n	ng/L
c a	c <sub>rs</sub>	c <sub>rs</sub> /c <sub>a</sub>	c	c rs	c_/c_ rs a
72.3	61.2	0.846	74.9	87.5	1.17
36.0	32.6	0.906	30.3	36.2	1.195
16.1	13.9	0.863	17.5	24.5	1.40
9.10	8.16	0.897	8.29	12.0	- <u>1.45</u> -
	Average rat	io = 0.878		Average rat	tio = 1.303

TABLE 2 Dilution Experiment Ni-Montmorillonite (pH = 9 Borate Buffer)

Total N	i Concentrations	s (µg/l)	
Control Vessels	Supernatant Vessels	Experimental Vessels	
m = 500  mg/l	m = 500  mg/l	$m = 188 mg/\ell$	
249 266	261 270	263 264	
269	275	266	
Dissolved Ni	Concentrations	(µg/l)	Experimental to Control Vessel Ratio
115 118 98.1	105 98.0 98.7	103 89.3 83.4	0.898 0.756 0.850
		Average	e Ratio = 0.835

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	Least		* ster Estimates of	$f \pi_{xc}$ , $v_x$ and	Square Parameter Estimates of $\pi_{xc}$ , $\nu_{x}$ and $K_{m}$ (pH = 9, Borate Buffer)	Suffer)
		Individual	Individual Data Sets		Joint Data Set: $v_{x} = 0.85(0.077)$	x = 0.85(0.077)
	No. of	#xc	x	Κ m (mo/1.)	<sup>π</sup> xc (L/g)	K (mg/L)
System	Points	(L/g) 42.5(7.3)	0.907(0.078)	21.3(4.9)	45.5(9.9)	18.7(5.0)
NI-CLAY	5	6.44	0.773	120	5.83(2.2)	146(59)
N1-Quartz	Ĩ	1.22(0.36)	0.677(0.20)	555 (280)	1.08(0.18)	786(161)

Estimated value (standard error of estimate)

\*

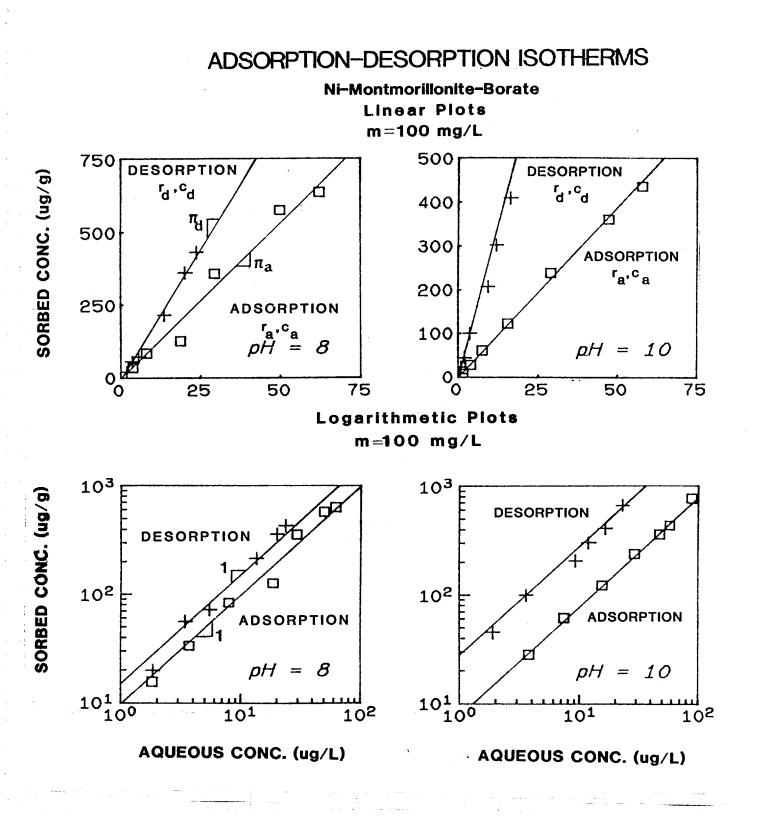
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TABLE 3

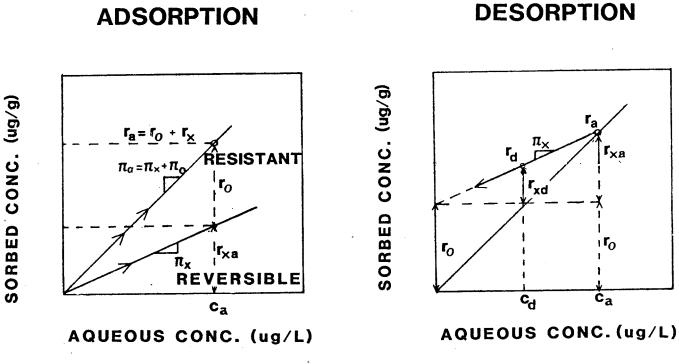
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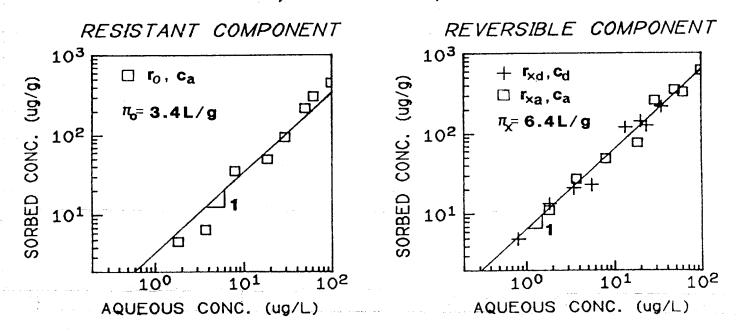


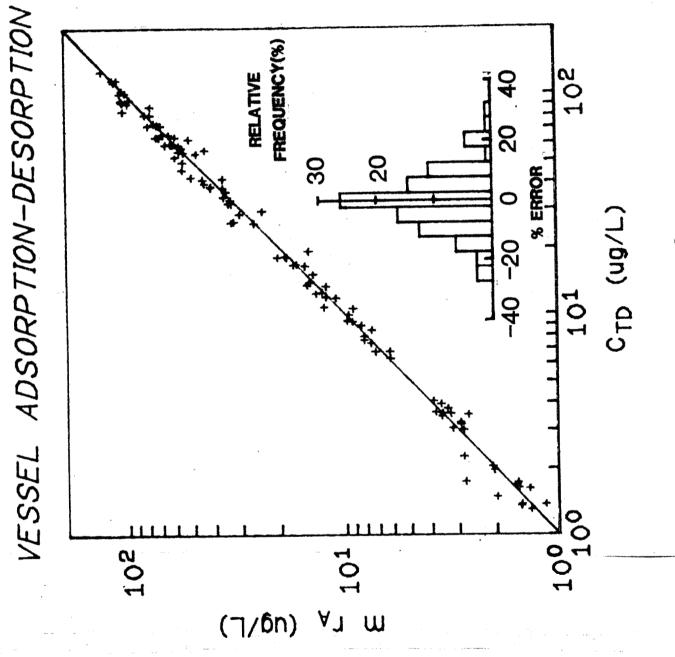




**Ni-Montmorillonite** 

pH=8, Borate Buffer, m=100mg/L



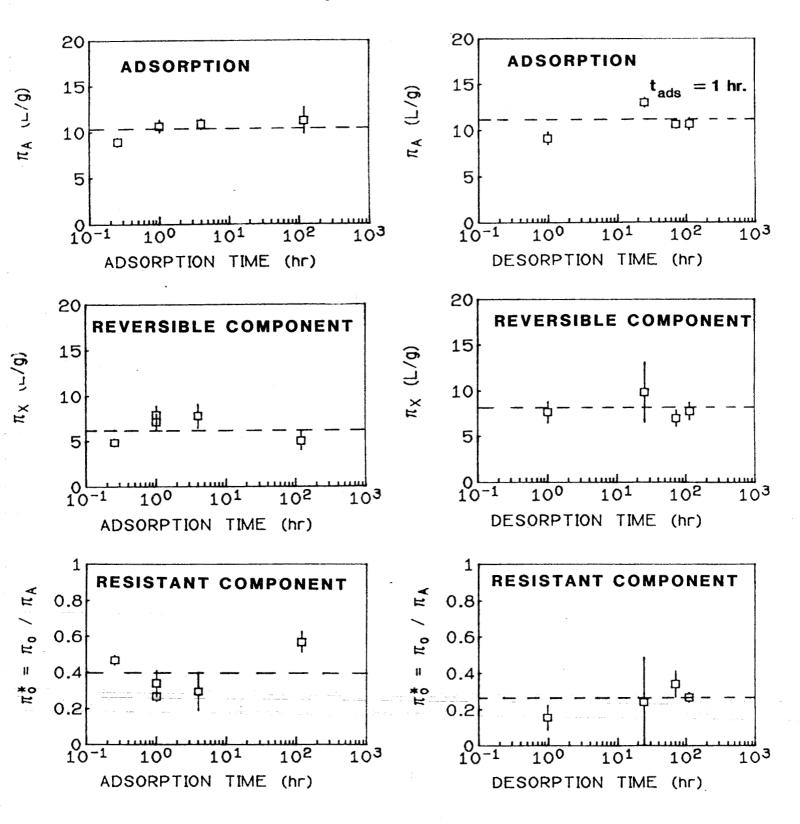


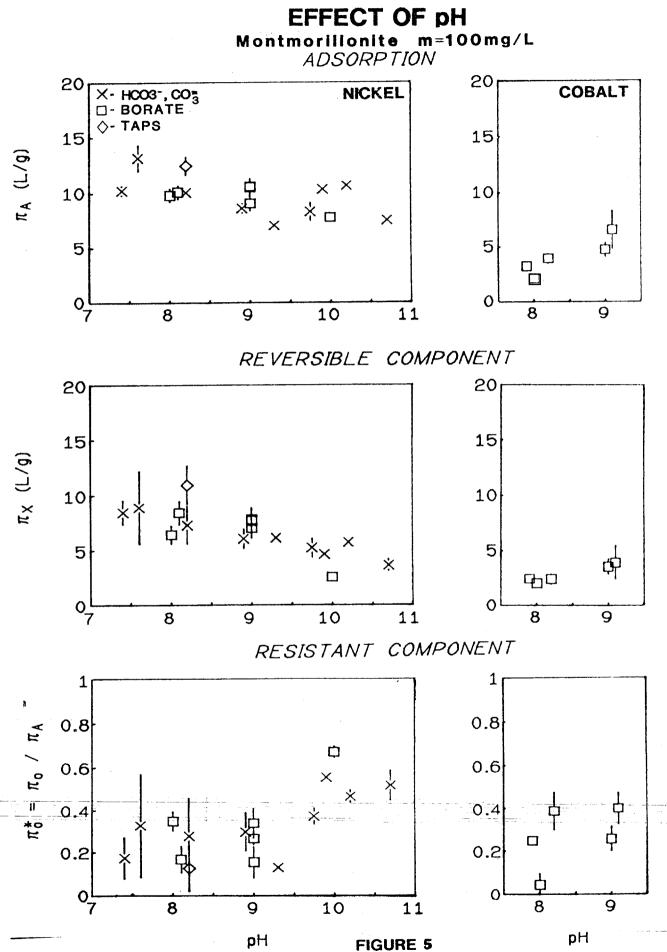
**ADSORPTION KINETICS** 

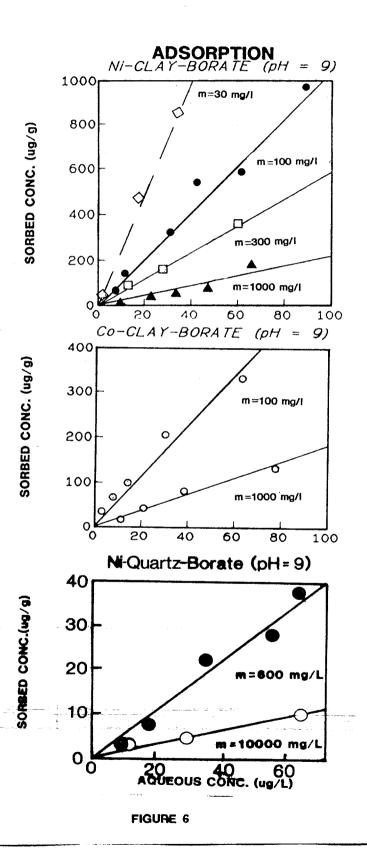
**DESORPTION KINETICS** 

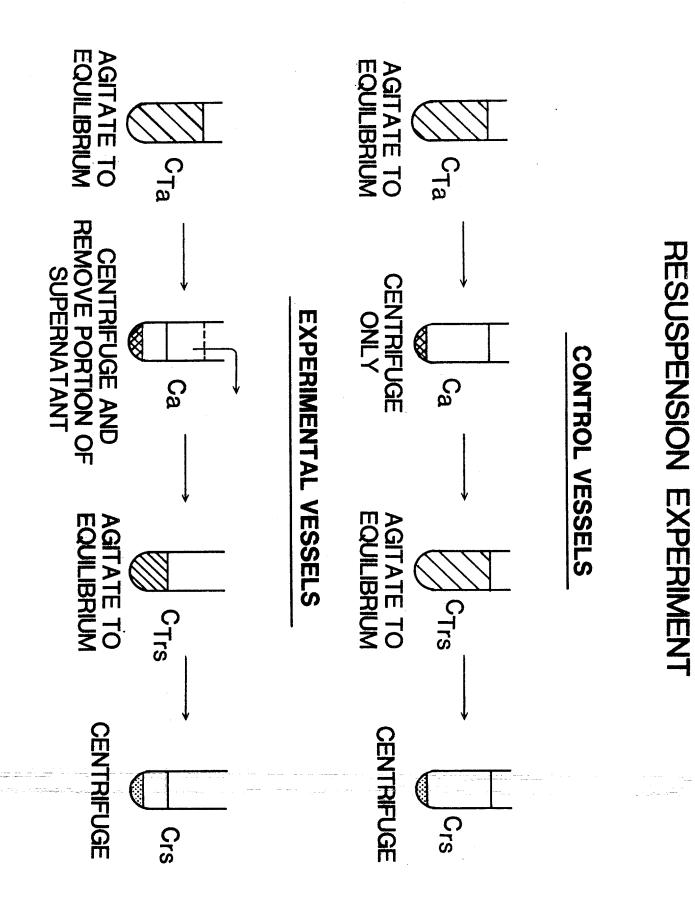
### Ni-Monmorillonite-Borate

pH=9 m=100 mg/L

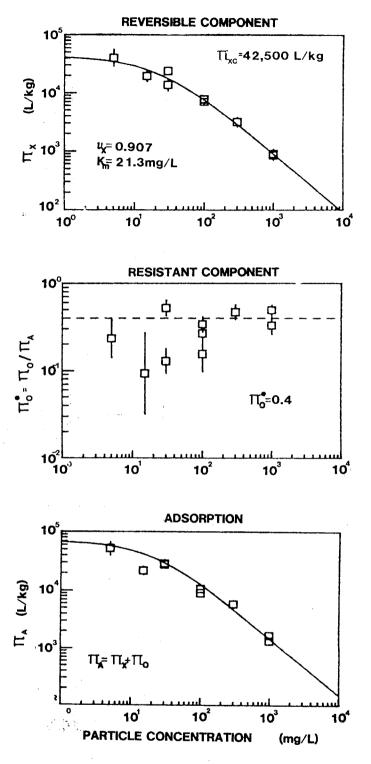




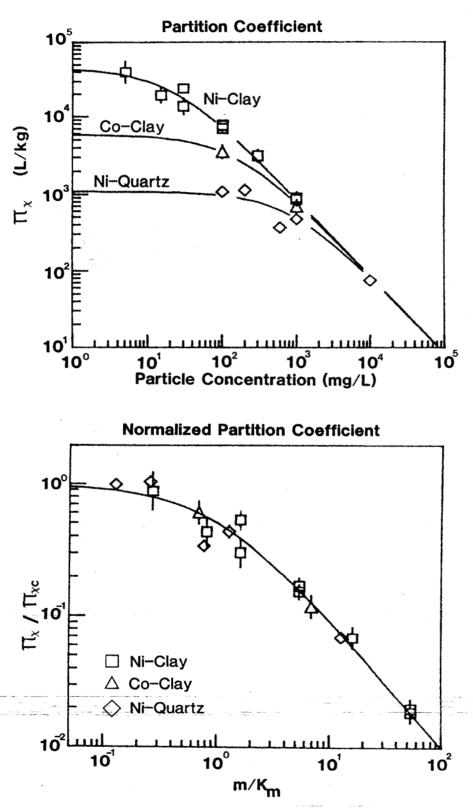




# PARTICLE CONCENTRATION EFFECT









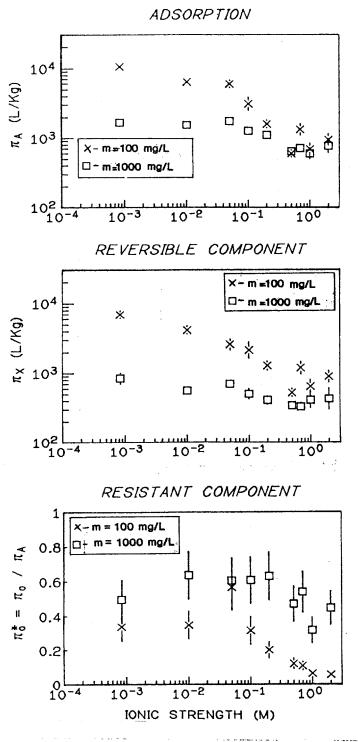


FIGURE 10

## REVERSIBLE COMPUNENT

## CLASSICAL PARTITION COEFFICIENT

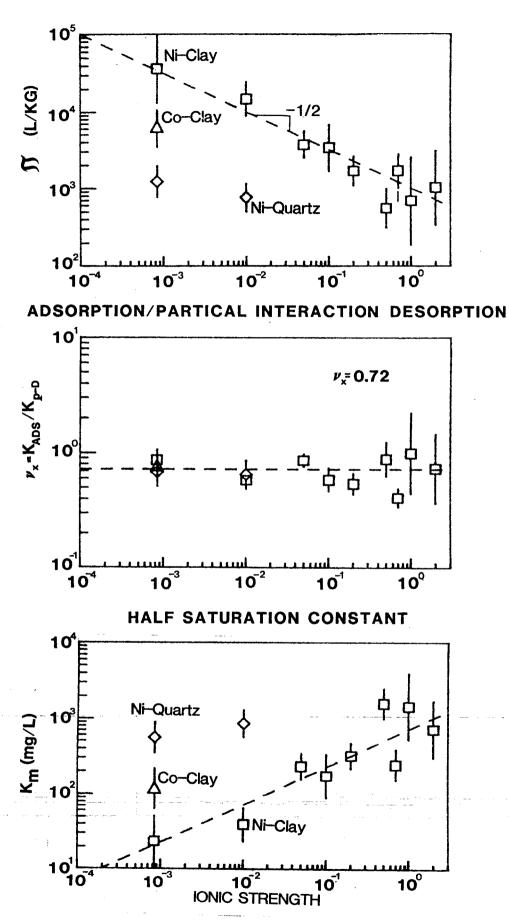


FIGURE 11

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### APPENDIX II

# A PARTICLE INTERACTION MODEL OF REVERSIBLE ORGANIC CHEMICAL SORPTION

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August 1984

#### INTRODUCTION

The sorption of hydrophobic organic chemicals onto soils, suspended, and sediment particles is the reaction which distinguishes the environmental fate of these chemicals from more conventional pollutants. The usual parameterization of this reaction involves the use of a partition or distribution coefficient that determines the fraction of total chemical concentration that is in the particulate and dissolved phases. The typical assumption is that the reaction is completely reversible.

Unfortunately there is a large body of experimental adsorptiondesorption data (1-54) which, in almost all cases, illustrates that complete reversibility is not commonly observed. Rather it is found that, for the usual time scales of these experiments (hours-days) an apparent desorption steady state occurs that is not in conformity with the particulate and dissolved concentration distribution predicted from the adsorption steady state partition coefficient. We employ an empirical model to quantify the partition coefficient of the reversibly sorbed component and attempt to relate it to chemical and particle properties.

The most surprising and controversial of these is the particle concentration itself, whose effect on the adsorption partition coefficient has been recently stressed by O'Connor and Connolly (55). We attempt to reconcile this effect for reversible component partitioning by proposing a particle interaction model that posits an additional desorption reaction. This model is shown below to be in conformity with observations for a large set of adsorption-desorption data (1-54). However the mechanism responsible for this desorption reaction is still uncertain. This is discussed below in more detail.

#### METHODS

The reversible and resistant component model of sorption (38) assumes that adsorbed chemical concentration,  $r_a$ , is composed of two components; a resistant component,  $r_o$ , that does not appreciably desorb under the conditions of the experiment, and a reversible component,  $r_x$ . The key assumption is that reversible component adsorption and desorption are both governed by the same linear isotherm:

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$$r_{\rm X} = \pi_{\rm X} c \tag{1}$$

where c is the aqueous phase concentration at either adsorption or desorption steady state.

Consider an adsorption experiment which yields an adsorbed concentration r<sub>a</sub> at aqueous concentration c<sub>a</sub>. Then since  $r_a = r_0 + r_x$ :

$$\mathbf{r}_{\mathbf{a}} = \mathbf{r}_{\mathbf{o}} + \pi_{\mathbf{x}} \mathbf{c}_{\mathbf{a}} \tag{2}$$

If, for a desorption into uncontaminated aqueous phase, the resulting sorbed concentration is  $r_d$  at aqueous concentration,  $c_d$ , then

$$\mathbf{r}_{d} = \mathbf{r}_{o} + \pi \mathbf{c}_{d} \tag{3}$$

since only the reversible component is affected. Thus the reversible component partition coefficient can be estimated by subtracting eq. (3) from eq. (2) yielding:

$$\pi_{\rm x} = \frac{r_{\rm a} - r_{\rm d}}{c_{\rm a} - c_{\rm d}} \tag{4}$$

Consider, now an adsorption isotherm comprised of  $i=1,\ldots,N_a$  concentrations r(i,o), c(i,o), the zero denoting the adsorption cycle. If, for each adsorption vessel a sequence of desorption cycles,  $j=1,\ldots,N_d$  are performed, yielding r(i,j), c(i,j), then the component model predicts that the slopes of the consecutive desorption isotherms:

$$\pi_{x}(i,j) = \frac{r(i,j) - r(i,j+1)}{c(i,j) - c(i,j+1)} \qquad \begin{array}{c} i=1,\dots,N\\ j=1,\dots,N, \end{array}$$
(5)

are all constant, independent of the isotherm index, i, or the desorption cycle index, j. Measurement errors and the idealized model assumptions will yield variations in each estimate. The geometric mean ignoring any negative values of  $\pi_{\chi}(i,j)$ , is used as the best estimate of  $\pi_{\chi}$  for each isotherm analyzed. Resistant component concentrations,  $r_{0}(i)$ , are estimated by subtraction:

$$r_{o}(i,j) = r(i,j) - \pi_{x}c(i,j)$$
 (6)

and for each isotherm the geometric mean is used as the best estimate of  $r_{o}(i)$ . The component model equation for the complete isotherm is:

$$r(i,j) = r_{0}(i) + \pi_{x}c(i,j)$$
 (7)

with  $\pi_{v}$  and  $r_{o}(i)$  as the fitted constants.

Two examples are presented in fig. 1. The top panels illustrate the isotherm data and the conformity to the model eq. (7) using the indicated value of  $\pi_x$ . The middle panels compare the calculated resistant component data  $r_o(i,j)$  (symbols) from eq. (6) together with the estimates of the presumed constant resistant component concentration,  $r_o(i)$  (horizontal lines). The bottom panel illustrates, in these cases at least, that the adsorption of the resistant component follows a linear (unity log-log slope) isotherm

$$\mathbf{r}_{0}(\mathbf{i}) = \pi_{0}\mathbf{c}_{0} \tag{8}$$

For parathion, the model provides an excellent fit to the data. For atrazine some deviations are present - the resistant component estimates are not exactly constant - but the fit to the experimental data (top panel) is reasonable and parsimonious since only two parameters are necessary:  $\pi_x$  and  $\pi_o$ . Fig. 2 presents other adsorption-consecutive desorption isotherms and the reversible-resistant component model isotherms (eq. 7), fit to the data as described above. These examples are chosen to illustrate that the model is applicable to isotherms from a wide range of organic and inorganic sorbates and sorbents: soils, lake sediments, and inorganic particles.

The division of sorbed chemical into two components - reversible and resistant - can be viewed as a convenient operational method that distinguishes between the rapidly sorbed, labile component, which also rapidly desorbs; and a nonlabile slowly desorbing component. In fact a kinetic description of the sorption process is more fundamentally correct. The gas purge experiments of Karickhoff (56) are a direct examination of these phenomena. He finds both a labile (rapidly adsorbing

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and desorbing) component and a nonlabile component with slower desorption rates. The separation we propose essentially isolates the labile component partitioning behavior within the context of available consecutive desorption experimental data.

A substantial portion of the reversible component partition coefficients analyzed below are computed from reported adsorption and consecutive desorption data using the average slope method (Table 1). Unfortunately not all reported desorption experiments include the actual data. For these cases estimates of  $\pi_x$  are made in various ways, depending on the reported results. The estimation equations are based upon the definition of  $\pi_x$ , eq. (4), and a mass balance requirement:

$$\operatorname{mr}_{a} + \alpha c_{a} = \operatorname{mr}_{d} + c_{d}$$
(9)

where m is the particle concentration and  $\alpha$  is the volume fraction of aqueous phase not removed before uncontaminated aqueous phase is added to initiate desorption. This equation requires that the chemical remaining in the vessel both on the particles,  $mr_a$ , and in the remaining aqueous phase,  $\alpha c_a$ , be accounted for after desorption as either on the particles,  $mr_d$ , or in the aqueous phase  $c_d$ . This is equivalent to assuming that no significant vessel adsorption or desorption occurs during desorption equilibration. Only experiments for which essentially the same particle concentration was maintained throughout are considered. Where possible the standard deviation of the estimated  $\pi_x$  is also computed. The result is over 200 estimates of  $\pi_x$  from these reported experiments for a wide variety of sorbates and sorbents.

#### REVERSIBLE PARTITION COEFFICIENTS: NEUTRAL ORGANIC CHEMICALS

The adsorption partitioning of neutral organic chemicals to soils and sediment particles has been shown to be a function of the weight fraction of sediment organic carbon,  $f_{oc}$ , and the octanol-water partition coefficient,  $K_{oW}$ , of the chemical (see Ref. 57 for an excellent review). The basic relationship is that the organic carbon normalized adsorption partition coefficient  $K_{oc} = \pi_a/f_{oc}$  is sediment independent so long as the swelling clay weight fraction of that sediment is not large and that a relationship exists between K and K of the form:

$$\log K_{oc} = a_{o} + a_{1} \log K_{ow}$$
(10)

Fig. 3 examines this relationship for the reversible component partition coefficients,  $\pi_{\chi}$ , sequentially plotted using various symbols with standard deviations indicated if they exceed the symbol size. The conventional organic carbon normalization applied to the reversible partition coefficient  $\pi_{\chi}/f_{oc}$  is plotted directly above the symbol as a plus. Directly below the symbol the quantity  $m\pi_{\chi}$  is displayed as a cross if m < 1 kg/L. For m = 1 kg/L,  $m\pi_{\chi} = \pi_{\chi}$  numerically and the cross is omitted for clarity of presentation. For each chemical-reference set, the data are ordered by increasing m and, for each m, by increasing f<sub>oc</sub> of the sorbent. The horizontal dotted lines are the K<sub>ow</sub> for each chemical.

It is apparent that for low  $K_{ow}$  chemicals  $\pi_x/f_{oc} \sim K_{ow}$  but that as log  $K_{ow} > 3$  this equality breaks down and in some cases  $\pi_x/f_{oc}$  is significantly less than  $K_{ow}$ . A careful examination of fig. 3 indicates that the deviation from  $\pi_x/f_{oc} \sim K_{ow}$  is correlated to the extent that  $m\pi_x$  approaches 1 and in particular that  $m\pi_x$  appears to be bounded in the region of  $m\pi_x \sim 1$ . Note especially the lower panel of fig. 3 where  $\pi_x/f_{oc}$  is markedly lower than  $K_{ow}$  and the crosses are all of order one. This suggests that  $m\pi_x \sim 1$  sets an upper limit to the extent of reversible sorption and in particular that the particle concentration itself influences the extent of sorption.

The observation that particle concentration affects the adsorption partition coefficient has been suggested previously by O'Connor and Connolly (55) who noted that certain chemicals strongly exhibited the effect whereas other chemicals were less affected. Significant particle concentration effects have been observed for hexachlorobiphenyl (HCBP) sorption (39). In particular it has been observed that  $\pi_x$  is inversely related to m via  $\pi_x = v_x/m$  and that the values of  $v_x$  observed for various particle types, to within a factor of four (0.3-1.2), are independent of sediment particle properties such as  $f_{oc}$  and particle identity (clay, silica etc). The experimental range of particle concentration was m = 10-22,000 mg/L.

This puzzling finding prompted the present analysis of available adsorption-desorption data for other chemicals in order to assess the generality of the finding.

#### PARTICLE INTERACTION MODEL

One immediate problem is that the relationship  $\pi_x = v_x/m$  cannot persist indefinitely as  $m \rightarrow 0$  since one is left with the absurd prediction that, in the limit, a single particle can reversibly sorb the same quantity of HCBP as can collectively be sorbed by many particles.

We have recently suggested a model for reversible component sorption (58) which relies upon the hypothesis that, in addition to the usual adsorption and desorption reactions:

$$k_{ads} c + m \rightarrow c=m$$
 (11a)

where c is dissolved chemical, m is the particle concentration, and c=m is sorbed chemical concentration, there exists an additional desorption reaction for the reversible component:

$$k_{p-d} c=m + m \rightarrow c + 2m$$
 (12)

which occurs as particles interact, possibly via close encounters or collisions. At steady state these reactions imply that

$$\pi_{x} = \frac{[c=m]}{[c][m]} = \frac{k_{ads}}{k_{s-d} + mk_{p-d}} = \frac{\pi_{xc}}{1 + m\pi_{xc}/\nu_{x}}$$
(13)

where  $\pi_{xc} = k_{ads}/k_{s-d}$ , the classical reversible partition coefficient, and  $\nu_x = k_{ads}/k_{p-d}$ , the ratio of adsorption to particle interaction induced desorption rate. Note that at low particle concentrations  $\pi_x \rightarrow \pi_{xc}$  which predicts that at sufficiently low particle concentration, there should be no effect on the extent of reversible sorption but that at large particle concentrations,  $\pi_x \rightarrow \nu_x/m$  so that  $m\pi_x \sim \nu_x$ . This suggests that the reason  $\pi_x/f_{oc}$  is systematically less than K for certain chemicals (fig. 3) is that additional particle induced desorption is lowering  $\pi_y$ .

This model provides a framework within which to analyze the  $\pi_{\chi}$  data presented in fig. 3. Following conventional adsorption theory we expect that  $\pi_{\chi c} = f_{oc} K_{oc}^{\chi}$  so that  $K_{oc}^{\chi}$  is the proper organic carbon normalization for the reversible partition coefficient at the low particle concentration limit; and that the octanol-water partition coefficient parameterizes the extent of chemical sorption to particle organic carbon. The result is that  $\pi_{\chi}$  is a function of m,  $f_{oc}$ , and  $K_{ow}$  via:

$$\pi_{x} = \frac{f_{oc} K_{oc}^{x}}{1 + m f_{oc} K_{oc}^{x} / \nu_{x}}$$
(14)

and

$$\log_{10} K_{oc}^{x} = a_{o} + a_{1} \log_{10} K_{ow}$$
(15)

The superscript x is used to emphasize that  $K_{oc}^{x}$  applies to reversible component partitioning. The idea is that  $\pi_{x} \rightarrow \pi_{xc}$  at low particle concentrations and in this limit the classical reversible partition coefficient,  $\pi_{xc}$ , is given by the usual organic carbon, octanol-water relationships:  $\pi_{xc} = f_{oc} K_{oc}^{x}$  with  $K_{oc}^{x}$  given by eq. (15). However at larger particle concentrations, m f<sub>oc</sub>  $K_{oc}^{x} >> v_{x}$ , only the particle concentration is important so that  $\pi_{x} \rightarrow v_{x}/m$ . In this region particle induced desorption is overwhelming spontaneous desorption and only the particle concentration, m, and the ratio  $v_{x} = k_{ads}/k_{p-d}$  controls the extent of reversible partitioning.

A nonlinear least squares fit (59) of eqs. (14, 15) to the neutral organic chemical data using the K<sub>ow</sub> given in Table 2, and for which f<sub>oc</sub> > 0 (N = 129 points) yields the parameter values ( $v_x$ ,  $a_o$ ,  $a_1$ ) shown in Table 3. The observed versus calculated  $\pi_x$ 's for each data point and a normalized plot of the entire data set as log averages over ranges of the x-axis are shown in fig. 4. The fit to eqs. (14, 15) is remarkable considering the range of chemicals, particle types, and the number of independent experiments and experimentalists. Figs. 5 and 6 present the

individual chemical data in normalized form. It is interesting to note that even for moderate log  $K_{oc}^{X}$  chemicals ( $\sim 2.5 - 3.0$ , fig. 5) the particle concentration effect is observed at higher mf values. Conversely high  $K_{oc}^{X}$  chemicals (e.g. Kepone, pp' DDT, fig. 6) can be examined at low enough mf so that  $\pi_{x} \cong f_{oc} K_{oc}^{X}$  as expected from hydrophobic partitioning theory. At large mf values these and other high  $K_{oc}^{X}$  chemicals exhibit a 3-4 order of magnitude supression of  $\pi_{x}$  relative to the low mf limit.

For certain chemicals enough data are available to fit eq. (14), individually. The results of the nonlinear least squares fit ( $\nu_x$ ,  $K_{oc}^x$ ) are listed in Table 3 and displayed in fig. 7 versus the chemical's  $K_{ow}$ . The log linear relationship of  $K_{oc}^x$  to  $K_{ow}$  is apparent. What is curious is the tendenty of  $v_x$  to be larger,  $v_x \sim 2-4$ , for the lower  $K_{ow}$  chemicals and to drop abruptly with Terbufos and remain at  $v_x \sim 0.5-1.0$  for the remaining chemicals. Whether this is an artifact of the data sets analyzed or a real indication that, in fact,  $v_x$  is not universally a constant is not clear. Certainly a variable  $v_x = k_{ads}/k_{p-d}$  is closer to one's intuitive expectations than is a constant  $v_x$ . More experiments in the transition region are clearly warranted. From the above data analysis it appears that  $v_x$  is of order one plus or minus roughly one half an order of magnitude.

#### EFFECTS OF CHEMICAL AND PARTICLE PROPERTIES AND CONCENTRATION

The proposed reversible component partitioning model has two distinct regions in which particle concentration either does or does not have an effect on  $\pi_x$ . For mf<sub>oc</sub> $K_{oc}^x < \nu_x$ ,  $\pi_x \approx f_{oc}K_{oc}^x$  which is a linear function of f<sub>oc</sub> and  $K_{oc}^x$ . Fig. 8 presents the data for mf<sub>oc</sub> $K_{oc}^x < 1$ . The expected relationships:  $\pi_x/f_{oc}$  versus K<sub>ow</sub> with the line corresponding to eq. (15), and  $\pi_x/K_{oc}^x$  versus f<sub>oc</sub> with the line corresponding to the expected linear relationship are shown. There are no apparent systematic deviations.

Conversely for mf  $_{oc}K_{oc}^{X} >> v_{x}$  only particle concentrations determine the magnitude of  $\pi_{x}$ . Fig. 9 presents the data for mf  $_{oc}K_{oc}^{X} > 3$ . Again no systematic deviations are apparent. It is interesting to note that even for chemicals with fairly low K<sub>ow</sub>'s - the chemicals in the upper

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panel range from  $\log_{10} K_{ow} = 2.81$  to 3.81 - the particle concentration can dominate  $\pi_x$  if it and f are large enough. In addition chemicals with large  $K_{ow}$ 's can exhibit extremely low reversible partitioning at high particle concentrations (e.g. HCBP, DDT, and Aroclor 1254).

We illustrate these effects on reversible partitioning using actual adsorption-desorption data in fig. 10 which are plots of adsorption - and single or consecutive desorption isotherms. Only one of the isotherms (there are typically 4 or 5) are plotted for each sorbent for clarity. For Diuron  $(\log_{10} K_{ow} = 2.81)$  we expect  $K_{oc}^{X} = 595$  and for m = 0.2 kg/L the breakpoint in  $f_{oc}$  above which  $f_{oc}$  should not influence  $\pi_{x}$  is  $f_{oc} = v_{x}/(mK_{oc}^{X})$  (Table 3) i.e.  $f_{oc} = 5.5\%$ . Note that the slope of the consecutive desorption isotherms,  $\pi_{x}$ , is increasing from  $f_{oc} = 1.9$  to 2.3% but that it slows its increase at  $f_{oc} = 6.2\%$  and is essentially constant as  $f_{oc}$  increases to 9.3%. For Terbufos at the same particle concentration, since  $K_{ow}$  is almost an order of magnitude larger than for Diuron, the breakpoint in  $f_{oc}$  is correspondingly less:  $f_{oc} = 0.24\%$  (Table 3). And, in fact, as shown in fig. 10 the slopes for  $f_{oc} = 0.51$  - 18.4% are all essentially equal as predicted.

Conversely for large K<sub>ow</sub> chemicals, varying the particle concentration for the same particle type strongly affects the slope. For both Kepone and HCBP at  $f_{oc} \approx 5\%$ , the particle concentration breakpoint m =  $v_x/(f_{oc}K_{oc}^x)$  is m = 84 mg/L and m = 5 mg/L respectively (Table 3) and as m increases beyond these values indeed the slope of the desorption isotherms decrease.

The predicted relationships: both the absence of  $\pi_x$  increases as  $f_{oc}$  increases at constant m (Terbufos), as well as the particle concentration effect (Kepone and HCBP) can be directly seen in these examples. Thus they are not artifacts of the data reduction techniques used to generate  $\pi_x$  but are real features of reversible component partitioning.

#### REVERSIBLE PARTITIONING - ORGANIC ACIDS, BASES AND INORGANIC SORBENTS

The particle interaction model can only be directly tested if, for each chemical of interest, the data span a large enough range in mf<sub>oc</sub> so that  $\pi_{xc}$  and  $\nu_{x}$  can both be estimated. Alternately an independent method of estimating  $\pi_{xc}$  can be utilized such as that employed in the previous section for neutral organic chemicals and organic carbon containing particles where  $\pi_{xc} = f_{oc} K_{oc}^{x}$  and  $K_{oc}^{x}$  is estimated from K. For ionizable organic chemicals or for sorption onto inorganic particles no analogous methods are available. Thus no direct test is possible.

However the model does make one prediction that is not dependent upon the magnitude of  $\pi_{\mathbf{vc}}$ :

$$\pi_{x} = \frac{\pi_{xc}}{1 + m\pi_{xc}/\nu_{x}} = \frac{\nu_{x}/m}{\nu_{x}} < \frac{\nu_{x}}{m}$$
(16)

namely that  $\pi_x$  should be less than  $\nu_x/m$  regardless of the magnitude of  $\pi_{xc}$ . Hence plots of  $\pi_x$  versus m should either exhibit values below  $\nu_x/m$  corresponding to insignificant particle interaction induced desorption, or along the line  $\pi_x = \nu_x/m$  corresponding to desorption being dominated by the particle interaction desorption mechanism.

Fig. 11 presents the observations for neutral organic chemicals and inorganic sorbents (nearly all clays). With rare exceptions the points are either below or on the line corresponding to  $v_x \cong 1$ . Fig. 12 presents similar results for organic acids and bases and the few available inorganic chemical data. Without an independent or at least generally applicable method of estimating  $\pi_{xc}$ , no further test is possible. However the substantial number of observed  $\pi_x$  that are near  $v_x/m$  suggests that the particle concentration effect on reversible partitioning is not limited to just neutral organic chemicals and organic carbon containing particles, but is a ubiquitous feature of reversible component partitioning.

A large set of metal sorption data: Ni, Co as sorbates with montmorillonite and quartz as sorbents that has the requisate range in particle concentration are also in conformity with the same reversible component partitioning model (58) employed above, eq. (14), thus supporting this view.

#### ALTERNATE POSSIBILITIES - A THIRD PHASE

The most popular explanation (60-62) for the particle concentration effect is that in addition to the aqueous and particle phases there exists a sorbing (third) phase which is not removed from the supernate by the particle separation technique (usually centrifugation) employed to operationally measure the "dissolved" concentration. This third phase is identified as being either dissolved organic matter or colloidal particles. To apply this hypothesis to reversible component partitioning we assume that the "dissolved" concentrations at adsorption and desorption steady state,  $c_a$  and  $c_d$ , are the sum of the truly dissolved concentrations,  $c'_a$  and  $c'_d$ , and that sorbed (or complexed) to the nonseparated third phase,  $m'_a r'_a$  and  $m'_d r'_d$ , where  $m'_a$  and  $m'_d$  are the nonseparated third phase aqueous concentrations (kg/L) and  $r'_a$  and  $r'_d$  are the sorbed chemical concentrations (mg/kg-third phase) at adsorption (a) and desorption (d) steady states. Thus

$$c_{a} = c_{a}^{\dagger} + m_{a}^{\dagger}r_{a}^{\dagger}$$
(17)

$$c_{d} = c'_{d} + m'_{d}r'_{d}$$
(18)

and  $\pi_{x}$  (eq. 3) becomes:

$$\pi_{x} = \frac{r_{a} - r_{d}}{(c_{a}^{\dagger} + m_{a}^{\dagger}r_{a}^{\dagger}) - (c_{d}^{\dagger} + m_{d}^{\dagger}r_{d}^{\dagger})}$$
(19)

or

$$\pi_{x} = \frac{r_{a} - r_{d}}{c_{a}'(1 + m_{a}'\pi_{a}') - c_{d}'(1 + m_{d}'\pi_{d}')}$$
(20)

where  $\pi_a^{\prime}$  and  $\pi_d^{\prime}$  are the adsorption and single desorption partition coefficients for the nonseparated phase. To convert this expression to that found to be descriptive of the data, eq. (14), a number of assumptions are required: (1) that the concentration of nonseparated third phase at adsorption and desorption steady states (for each consecutive desorption step) are equal:  $\pi_a^{\prime} = \pi_d^{\prime} = \pi^{\prime}$  and (2) that the partitioning to this phase is reversible:  $\pi_a^{\prime} = \pi_d^{\prime} = \pi^{\prime}$  in which case eq. (20) becomes:

$$\pi_{x} = \frac{r_{a} - r_{d}}{(c_{a}' - c_{d}')(1 + m'\pi')}$$
(21)

Defining:

$$\pi_{xc} = \frac{r_a - r_d}{c_a' - c_d'}$$
(22)

as the "true" partition coefficient the result is:

$$\pi_{\mathbf{x}} = \frac{\pi_{\mathbf{x}\mathbf{c}}}{1 + \mathbf{m}^{\mathsf{T}}\pi^{\mathsf{T}}}$$
(23)

In order to convert this equation to eq. (14), which describes the data, multiply  $m'\pi'$  by  $m\pi_{xc}/m\pi_{xc} = 1$  so that:

$$\pi_{x} = \frac{\pi_{xc}}{1 + m\pi_{xc}} \left(\frac{m'\pi'}{m\pi_{xc}}\right)$$
(24)

from which  $v_x$  in eq. (14) is:

$$v_{\mathbf{x}} = \frac{\mathbf{m}\pi_{\mathbf{x}\mathbf{c}}}{\mathbf{m}^{\dagger}\pi^{\dagger}}$$
(25)

The empirical finding that  $v_x = 1$  requires that

$$\mathfrak{m}'\mathfrak{n}' \cong \mathfrak{m}\mathfrak{n}_{\mathbf{X}}$$
(26)

For neutral organic chemical and organic carbon containing particles:  $\pi_{xc} = f_{oc} K^{x}_{oc}$  and if we assume (3) that the same normalization applies to the nonseparated phase:  $\pi' = f'_{oc} K^{x}_{oc}$ , then eq. (26) requires that

$$\mathbf{m}'\mathbf{f}'_{\text{oc}} \stackrel{\text{and}}{=} \mathbf{m}\mathbf{f}_{\text{oc}} \tag{27}$$

that is, that the quantity of organic carbon associated with the nonseparated phase is approximately equal to that contained in the separated particles. Alternately if specific surface area,  $\sigma(m^2/kg)$ , is a proper normalization for inorganic particles then  $\pi_{xc} = \sigma K_{\sigma}$  and  $\pi' = \sigma' K_{\sigma}$  so

that  $v_x \cong 1$  requires that

i.e. that the surface area/unit aqueous phase volume  $(m^2/L)$  for the nonseparated and separated phase, which presumably are particles in the case of inorganic sorbents, be equal.

There are a number of difficulties with this explanation. If we assume that the third phase are colloidal particles, it requires that nonseparable particles be in the same concentrations at adsorption, m', and desorption, m'. If they are not separated at the adsorption step they remain in the aqueous phase after the first centrigugation - they are then discarded when the supernatant is removed and uncontaminated aqueous phase is added. Thus during the desorption agitation they must be regenerated to the same concentration as they were present during adsorption. Further the requirement that follows from  $v_{\mathbf{y}} \approx 1$ , eqs. (27-28) demands that they be quantitatively of equal sorptive capacity as the separated particles. That these specialized requirements obtain for all experiments, particles, etc. analyzed above appears to be highly unlikely. In fact for one HCBP experiment (Ref. 39, Sta. #69) the particles were wet seived and only the greater than 75 µm particles were employed. This  $\pi_{v}$  is in conformity with the rest of the HCBP data (fig. 6).

If, on the other hand, the third phase is assumed to be dissolved organic carbon that desorbs from the particles, then the requirement (eq. 27) that  $m'f'_{oc} \cong mf_{oc}$  forces the conclusion that there is as much dissolved organic carbon in the aqueous phase as there is particulate organic carbon in the particulate phase, per unit aqueous phase volume, at both adsorption and at each desorption. In fact  $m'f'_{oc} = mf_{oc}$  requires that desorable particle organic carbon has a partition coefficient of  $\sim 1 \text{ L/kg-organic carbon which is extremely small. Even liquid octanol which has an aqueous solubility of <math>\sim 10^{-3}$ M and therefore a partition coefficient of  $\sim 10^3 \text{ L/mole-octanol} = 10.4 \text{ L/kg-organic carbon is ten-fold larger than that required for desorable organic carbon.$ 

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(28)

Another difficulty with such a low partition coefficient is that substantial organic carbon would be removed at each desorption cycle. Hence at adsorption equilibrium, one-half of the original particulate organic carbon mass must desorb and be in solution. Since this is discarded before desorption is initiated, it is difficult to see how  $m'_a = m'_d$ . Rather  $m'_d = m'_a/2^j$  where j is the desorption cycle index. It would further suggest that simply washing particles in water should remove substantial quantities of particle-bound organic carbon. Finally it is difficult to see how dissolved organic carbon could be implicated with inorganic sorbents.

Perhaps the most convincing experiments which appear to preclude nonseparable third phase as an explanation are resuspension experiments (58,63) in which, instead of desorption into new uncontaminated aqueous phase, the particles are resuspended into a reduced volume of the aqueous phase remaining after adsorption and centrifugation. For this experimental design the concentration of the nonseparated third phase in the supernatant must be the same at both adsorption and resuspension Only the concentration of the separated particles is equilibration. increased by resuspension into a reduced aqueous phase volume. For these experiments the partition coefficient is observed to decrease as particle concentration increases from the adsorption to the resuspension equilibrations. If a desorbing third phase (dissolved organic carbon) is involved then the problem of explaining the decreased partitioning is just transferred to explaining why there is increased desorption of the third phase into the original supernatant when only particle concentration is increased at the resuspension equilibration.

Dilution experiments (58,63) dispense with centrifugation of the experimental vessel completely and decreases particle concentration by adding either uncontaminated aqueous phase (63) or equivalently contaminated aqueous phase from a parallel vessel (58). These experiments yield increasing partition coefficient with decreasing particle concentration.

Thus nonseparable third phase models require that the nonseparable phase has properties that are both ubiquitous, very specialized and, for desorbing organic carbon, not chemically realistic. Further they appear to be precluded by resuspension and dilution experiments designed expressly to strongly discriminate against such models.

#### Desorption Kinetics and Particle Aggregation

An examination of desorption kinetics using an air stripping technique to remove dissolved chemical from the sorbent suspension (56) precludes the need for consecutive particle separations by centrifugation that characterizes the desorption data analyzed above. The results of a series of these experiments, presented by Karickhoff and Morris (64), indicates that there exists a labile component that desorbs rapidly (< 1 hr) and a nonlabile or resistant component whose desorption rate they correlate to the equilibrium partition coefficient:  $K_p = f_{oc}K_{oc}$ . This result and square root of time dependence of desorption of the nonlabile fraction strongly suggest intraparticle diffusion as the kinetic mechanism controlling nonlabile component desorption.

For the labile component, Karickhoff and Morris (64) found that the labile fraction of total sorbed chemical,  $X_{l}$ , decreased as mK increased, where m is the particle concentration. Their data analysis suggested:

$$X_{\ell} \sim \frac{1}{2.5 + mK_{p}/10}$$
 (29)

If we identify  $X_{\ell} = r_x/r_a$ , consistent with the component model, eq. (1-2) then  $X_{\ell} = \pi_x/\pi_a$  and if at sorption equilibrium  $\pi_a = K_p = f_{oc}K_{oc}$ , then:

$$X_{\ell} = \frac{\pi_{x}}{f_{oc}K_{oc}}$$
(30)

Using eq. (14) for  $\pi_x$  and assuming  $K_{oc} = K_{oc}^x$  yields:

$$X_{\ell} = \frac{\pi_{x}}{f_{oc}\kappa_{oc}^{x}} = \frac{1}{1 + mf_{oc}\kappa_{oc}^{x}/\nu_{x}}$$
(31)

which is the same functional form as that suggested by Karickhoff and Morris, eq. (29). The results of a fit of their data to eq. (16), is

shown in Table 4. A significantly larger  $v_x = 11.6(\pm 7.22)$  is found as well as the expected log linearity between  $K_{oc}^{x}$  and  $K_{ow}$ . The implication to be drawn from the larger  $v_x$  is unclear at present - in particular is it related in some way to the completely different experimental design but the remarkable result is that the same functional form, eq. (31), also describes their data. Since third phase models clearly do not apply to their experiment we take this to be further evidence that third phase explanations are inappropriate.

Karickhoff and Morris (64) suggest that particle aggregation may be the cause of the particle concentration effect they observed. Although this may be possible we have microscopically examined inorganic particle suspensions at various concentrations that exhibit the effect (58) and find no evidence of larger aggregations at higher particle concentrations. Also, if particle aggregation were the sole cause of the decreasing labile component sorption one would expect that the functional dependency would be solely in terms of m, rather than  $mf_{oc}K_{oc}^{x}$ . The latter expression includes the chemical property  $K_{oc}^{x}$  and the organic carbon fraction,  $f_{oc}$ , as well as the particle concentration. It is difficult to imagine a particle aggregation model that yields this functional dependence.

#### Discussion

It should also be stressed that the particle interaction model which is based upon the hypothesis that an additional desorption reaction exists, eq. (12), is no more successful in explaining why a constant  $v_x = k_{ads}/k_{p-d}$  appears to be characteristic of reversible component sorption, although fig. 7 suggests that it might vary systematically and the gas purge experiments suggest it may not be constant but rather a function of experiment design. One can speculate that both the adsorption and particle interaction induced desorption rate are physically rather than chemically determined by, for example, the level of turbulence induced by the agitation. While the turbulence level may be quite different among the experiments analyzed in this paper, it is only the ratio,  $k_{ads}/k_{p-d}$ , that needs to be constant and the absolute magnitude of the turbulence may cancel out of the ratio.

One piece of evidence that supports the particle interaction model is that, in the absence of agitation, the prediction  $\pi_x = f_{oc} K_{oc}^x$  is verified for HCBP sorption. In a series of dual tag diffusion experiments using syringes (65) with a layer of uncontaminated sediment overlying a layer of tagged (H<sup>3</sup>-HCBP and C<sup>14</sup>-graphite particles) sediment the K<sub>oc</sub> for HCBP is estimated to be  $10^{6.61 \pm 0.13}$  whereas from conventional agitated adsorption-desorption experiments (Table 3, fig. 6) K\_{oc}^x is estimated to be  $10^{6.50 \pm 0.31}$ . This concurrence can be explained by observing that in the syringe experiments the particles are stationary and, therefore, particle interaction induced desorption which is associated with particle motion is not occurring. Hence the syringe partition coefficients should correspond to the low particle concentration limit inferred from the agitated experiments. The fact that this is the case suggests that it is particle motion and, therefore, particle interactions that enhance desorption if particle concentrations are large enough and, as a consequence, particle interactions are frequent enough for this reaction to be significant.

Finally we observe that partition coefficients (actually distribution coefficients) inferred from field measurements of dissolved and particulate chemical concentrations also display an inverse relationship to particle concentration (39,66) so that the effect is an important factor in determining the partitioning of chemicals between particulate and dissolved phases in natural waters. Preliminary analysis of collections of these data sets appear to conform to eq. (14). A detailed analysis is forthcoming.

The most significant shortcoming of the model presented above is the lack of a mechanistic explanation for the particle interaction induced desorption reaction, eq. (12). One would surely like to know what, exactly, is the physics and chemistry of this reaction. Or, more critically, if the reaction itself is the correct explanation for what appears to be the correct result, eq. (14). The success of eq. (14) in correlating large amounts of adsorption-desorption data is undeniable, fig. 4-6, but its derivation via the reaction scheme, eq. (11-12) must be regarded as speculative.

From a practical point of view, the correlative power of the model for neutral organic chemicals and organic carbon containing particles is quite useful since a knowledge of the chemical's  $K_{ow}$ , the particle's  $f_{oc}$  and its concentration, m, suffices to determine  $\pi_x$ , the reversible component partition coefficient. For inorganic sorbents and ionizable or inorganic chemicals an upper bound to the expected partition coefficient,  $\pi_v < v_v/m$ , is available.

From a theoretical point of view, the most surprising and intuitively disturbing result is the observation that  $v_x$  is of order one ( $\pm$  one half an order of magnitude) for all chemicals and particle types examined. This generality suggests either that there exists a universal feature of reversible sorption which has heretofore been unsuspected or that the entire analysis itself is flawed in some subtle way. The author prefers to believe that the former is the case and that its explanation is via a particle interaction induced desorption reaction.

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TABLE	: 1. METHODS FOR CALCULATING $\pi_{\mathbf{x}}$	
Reported Data	Formula Employed <sup>(a)</sup>	References
Adsorption and Consecutive Desorption Isotherm data	Geometric Average of $\pi = \frac{r(i,j) - r(i,j+1)}{c(i,j) - c(i,j+1)}$	1,2,4,5,6,7,9,10,11,12 14 <sup>*</sup> ,15 <sup>*</sup> ,17,33,49,50,53
Adsorption and Single Desorption Isotherm data	8	3,13,16 <sup>*</sup> ,21,41 <sup>*</sup>
Single Point Adsorption, Single or Consecutive Desorption data	Ŧ	18,26,28,37,51,54
Adsorption and Single Desorption Partition Coefficients	$\pi_{x} = \frac{\pi_{a} - \alpha \pi_{d}}{1 - \alpha + m(\pi_{d} - \pi_{a})}$	20,23,27,28,30,35 38,39,40,43,44
Single Point Adsorption and Desorption: Sorbed data only	$\pi_{\rm X} = \frac{\Delta r}{(1-\alpha)(c_{\rm T}-mr) - m\Delta r}$	22,24,32,34
Adsorption Isotherm Partition Coefficient Single Point Desorption: Data as fraction desorbed	$\pi = \frac{\delta \pi}{1 - \alpha - m\delta \pi}$	19,25,29,31
Single Point Adsorption and Desorption: Aqueous concentrations	$\pi = \frac{c_d - \alpha c_a}{m(c_a - c_d)}$	28
Single Point Consecutive Desorptions: Data reported as sorbed fraction remain- ing at cycle j	Least Square fit of: $\frac{r(1,j)}{r_{T}} = \frac{r_{0}}{r_{T}} + \frac{1}{1 + m\pi} \left( \frac{\alpha + m\pi}{1 + m\pi} \right) (1 - \frac{r_{0}}{r_{T}})$	45,46,48
(a) See Notation list for definitions		

Data reconstructed from reported Freundlich constants and mass balance equation (9)

•															
Ref.	[69]	[69]	[14]	[89]	[77]	[89]	[43]	[67]	[74]	[73]	[67]	[67]	[72]	[74]	[72]
$\log_{10} \mathrm{K}_{\mathrm{ow}}$	3.80	3.81	3.87	4.99	5.18	5.34	5.50	5.57	5.91	6.03	6.19	6.31	6.50	6.89	7.75
-i   '				-								-			
1			thion	S					Iracene	54			ene		
Chemical	Beta-HCH	Alpha-HCH	Methyl Parathion	Chlorpyrifos	Pyrene	Endrin	Kepone	Γ.	Benz(a)Anthracene	Aroclor 1254	pp' DDT	Leptophos	Benz(a)Pyrene	ex	e.
	Bet	Alp	Met	Chl	Pyr	End	Kep	DDT	Ben	Aro	, dd	Lep	Ben	Mirex	HCBP
Ref.	[67]	[68]	[67]	[68]	[71]	[67]	[67]	[67]	[67]	[67]	[20]	[67]	[ 3]	[67]	[69]
	<u> </u>	Ē	<u> </u>	ĩ	, Land	Ē	Ξ	Ξ	Ξ	ũ	<b>ن</b> .	<u> </u>	<b>لىسىر</b>	ĩ	Ŭ,
$\log_{10}^{\mathrm{K}}$ ow	0		2	6	0		,i	88	2	2	0	0	80	6	8
 log_	0.70	1.60	2.12	2.19	2.20	2.81	2.81	2.888	2.92	3.02	3.10	3.40	3,68	3.69	3.72
I											a		·		•
Chemica1	Aldicarb	Carbofuran	Monuron	Linuron	Fluometron	Diuron	<b>Carbary1</b>	Permethrin	Phorate	Diazinon	Napropamide	Parathion	Terbufos	Dieldrin	Gamma-HCH
	Alc	Cal	Мог	Lin	Flu	Diu	Car	Per	Phc	Dia	Nap	Par	Ter	Die	Gam

TABLE 2. OCTANOL-WATER PARTITION COEFFICIENTS

# TABLE 3. NONLINEAR LEAST SQUARES FIT OF $\log_{10} \pi_x$

$$\pi_{\mathbf{x}} = \frac{f_{oc} K_{oc}^{\mathbf{x}}}{1 + m f_{oc} K_{oc}^{\mathbf{x}} / v_{\mathbf{x}}}$$

Entire data set (N = 129)

$$\log_{10} K_{oc}^{X} = a_{o} + a_{1} \log_{10} K_{ow}$$

 $v_x = 1.4(0.18); a_0 = 0.178(0.16); a_1 = 0.924(0.053)$ 

 $\sigma(\log_{10}\pi_{x}) = 0.408$ 

# Individual Chemicals

Chemical	No. of Data Points	$\frac{\log_{10} K_{oc}^{x}}{(L/kg)}$	ν <u>x</u>	σ(log <sub>10</sub> <sup>π</sup> x)
Aldicarb	5	1.01(0.16)	-	
Diuron	18	2.60(0.10)	4.39 (2.7)	0.312
Phorate	5	2.58(0.05)	3.02 (0.46)	0.057
Parathion	17	3.45(0.12)	4.73 (1.4)	0.246
Terbufos	5	3.29(0.71)	0.952(0.31)	0.200
Lindane	7	4.04(0.42)	0.619(0.32)	0.470
Kepone	10	5.28(0.32)	0.865(0.27)	0.328
pp' DDT	6	7.16(2.4)	1.60 (0.73)	0.343
НСВР	12	6.50(0.31)	0.694(0.06)	0.109

Parameter Estimate (Standard Error of Estimate)

\*

TABLE 4. Least	TABLE 4. Least Square Fit of $\pi_x = X_g K_p^{**}$ to Equation 14	= X <sub>&amp;</sub> K <sup>**</sup> to Equat	ion 14
	$v_{\rm x} = 11.6 \ (7.22)$	2)	
Chemical	$\log_{10}^{\mathrm{K}_{\mathrm{oc}}^{\mathrm{X}}}$	$\log_{10}\mathrm{K_{ow}}$	Ref. for K <sub>ow</sub>
Trifluralin	3.94(0.18)	3.06,5.34	[67],[68]
Pyrene	4.64(0.22)	4.88,5.18	[77],[68]
Pentachlorobenzene	4.97(0.37)	4.94,5.19	[76],[68]
Hexachlorobenzene	5.50(0.39)	5.23,5.50	[68],[75]

Parameter (Standard error of estimate)  $\sigma(\log_{10} \pi_x) = 0.302$ 

\*

\*\* Karickhoff and Morris (64)

# NOTATION

r(i,j) = sorbed concentration (mg/kg); i<sup>th</sup> isotherm, j<sup>th</sup> desorption cyclec(i,j) = aqueous concentration (mg/L); i<sup>th</sup> isotherm, j<sup>th</sup> desorption cycle= volume fraction of aqueous phase not removed before desorption  $= \pi/f_{oc}$ , organic carbon normalized adsorption partition coeffi $r_a(r_d)$  = adsorption (single desorption) sorbed concentration (mg/kg)  $c_a(c_d)$  = adsorption (single desorption) aqueous concentration (mg/L)  $\pi_a(\pi_d)$  = adsorption (single desorption) partition coefficient (L/kg) = resistant component sorbed chemical concentration (mg/kg) = reversible component partition coefficient (L/kg), eq. 3 =  $mr_a + c_a$  = initial total chemical concentration (mg/L) =  $r_a - r_d$ , desorbed chemical concentration (mg/kg) = initial sorbed chemical concentration (mg/kg) octanol-water partition coefficient (L/L) =  $\Delta r/r_a$ , fraction chemical desorbed = particle concentration (kg/L) cient (L/kg-organic carbon) ي د Mow × H

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#### FIGURE CAPTIONS

- Fig. 1. Components model fit to consecutive desorption data for Parathion [2] and Atrazine [6].
- Fig. 2. Components model fit to consecutive desorption data for Picloran [7], Diuron - Lake Sediment [17] SR-90 [11], Diuron -Montmorillonite [13], Fluometron [9], and Phenanthroline [4].
- Fig. 3. Sequential plot of  $\pi_x$  (various symbols);  $\pi_x/f_{oc}(+)$ ,  $m\pi_x(x)$ , ordered by increasing m and increasing  $f_{oc}$ . Vertical dotted lines separate chemicals, horizontal dotted lines are the chemical's  $K_{ow}$ .
- Fig. 4. Comparison of eq. (14-15) to data: observed versus calculated  $\pi_{\nu}$  (top) and in normalized form (bottom).
- Fig. 5. Normalized plots of individual chemical data.
- Fig. 6. Normalized plots of individual chemical data.
- Fig. 7. Parameters obtained from individual fits of eq. (14) to chemical data versus octanol/water partition coefficient.
- Fig. 8. Low particle concentration limit relationships: Log linearity of K to K (top) and linearity of normalized  $\pi$  to f oc (middle and bottom). The latter two lines have slope one.
- Fig. 9. The high particle concentration limit.  $\pi_x$  as a function of m only.
- Fig. 10. Consecutive desorption data for varying f (top) and particle concentration (bottom).

- Fig. 11. Comparison of data to the prediction that  $\pi_x < \nu_x/m$  for sorption to inorganic sorbents.
- Fig. 12. Comparison of data to the prediction that  $\pi_x < \nu_x/m$  for sorption of inorganic chemicals (top) and organic acids, and bases (middle and bottom).

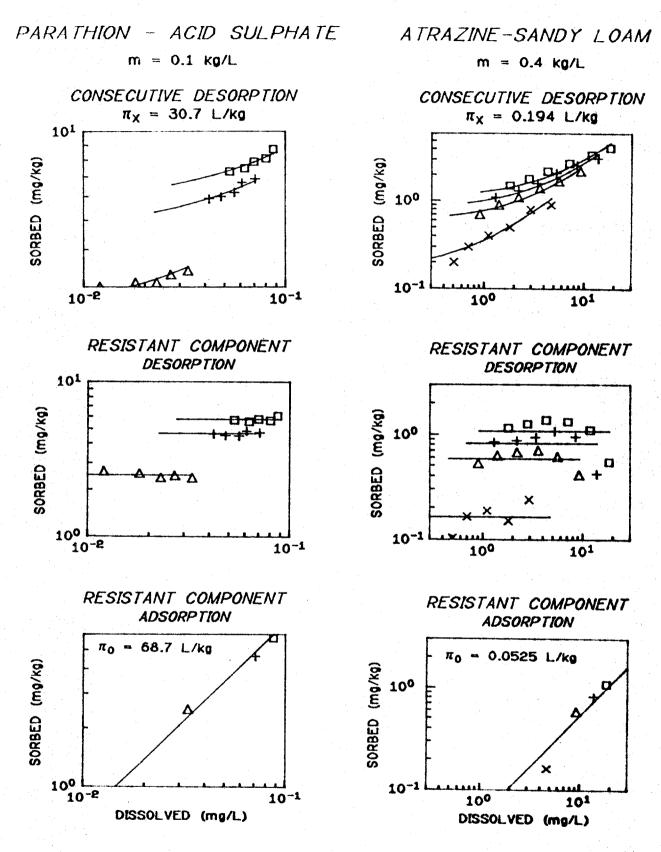
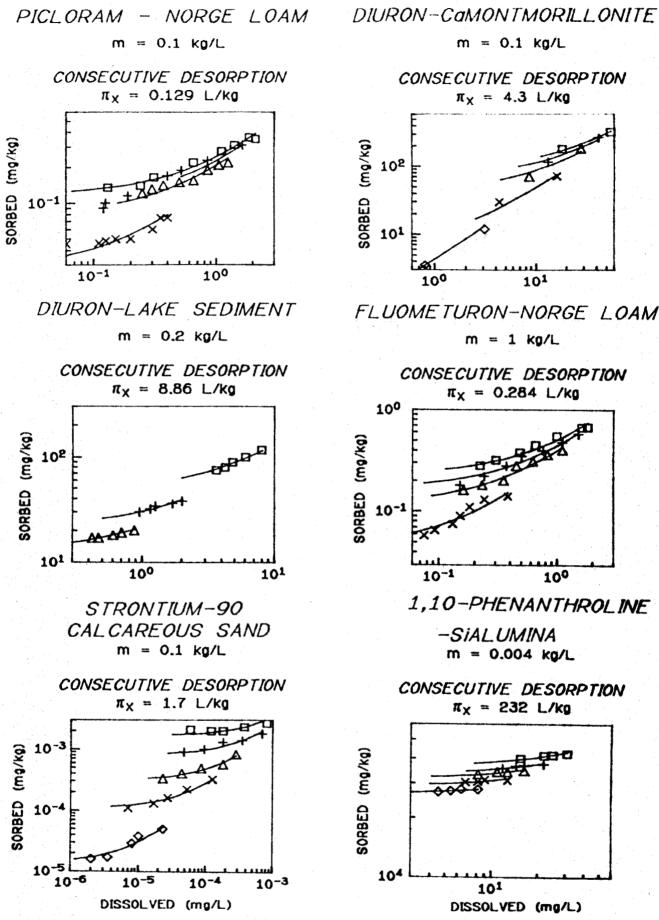
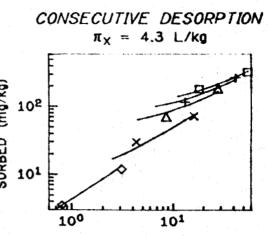


FIGURE 1

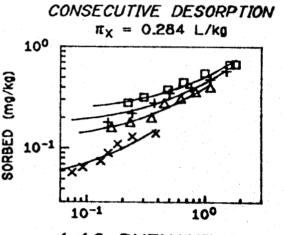


#### m = 0.1 kg/L



FLUOMETURON-NORGE LOAM

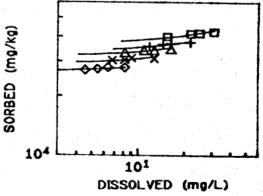
m = 1 kg/L

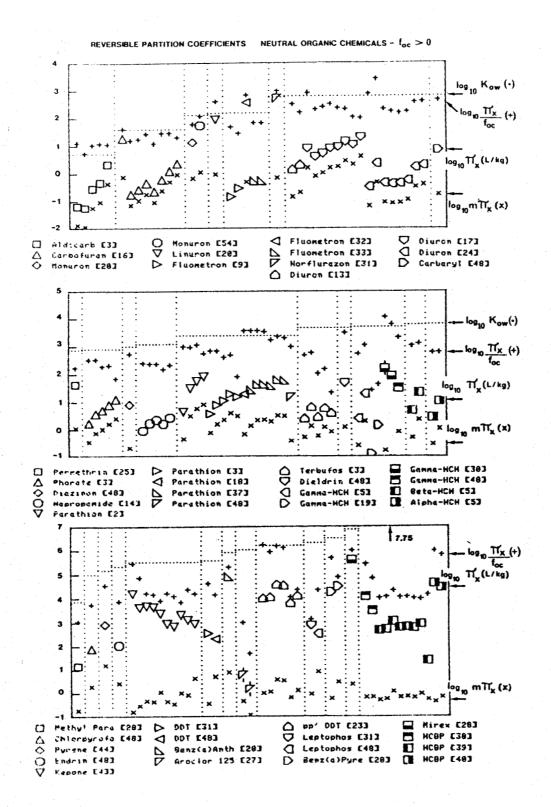


1,10-PHENANTHROLINE

-SIALUMINA m = 0.004 kg/L

CONSECUTIVE DESORPTION  $\pi_{\rm X} = 232 \, {\rm L/kg}$ 

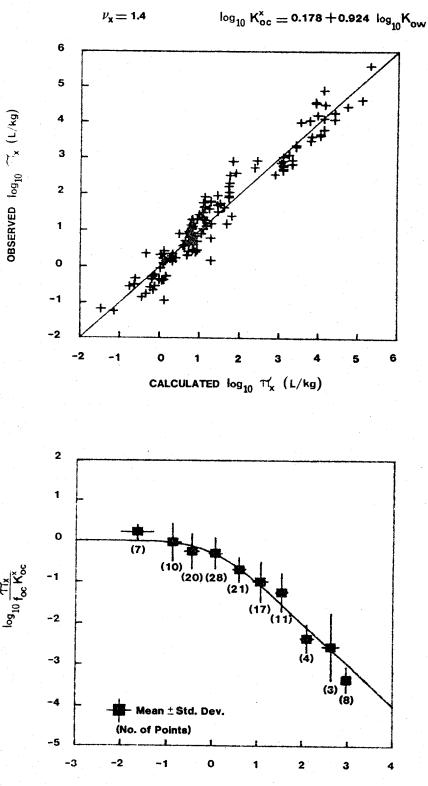




REVERSIBLE COMPONENT PARTITION COEFFICIENTS

COMPARISON TO PARTICLE INTERACTION MODEL

NEUTRAL ORGANIC CHEMICALS



 $\log_{10} m \ f_{oc} \ K_{oc}^{x} \, / \, \nu_{x}$ 

FIGURE 4

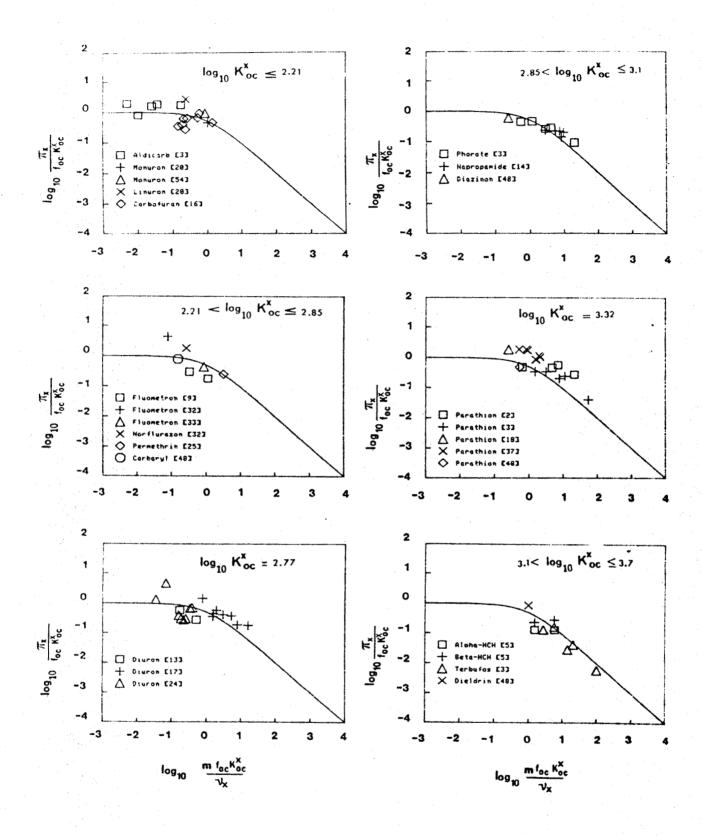


FIGURE 5

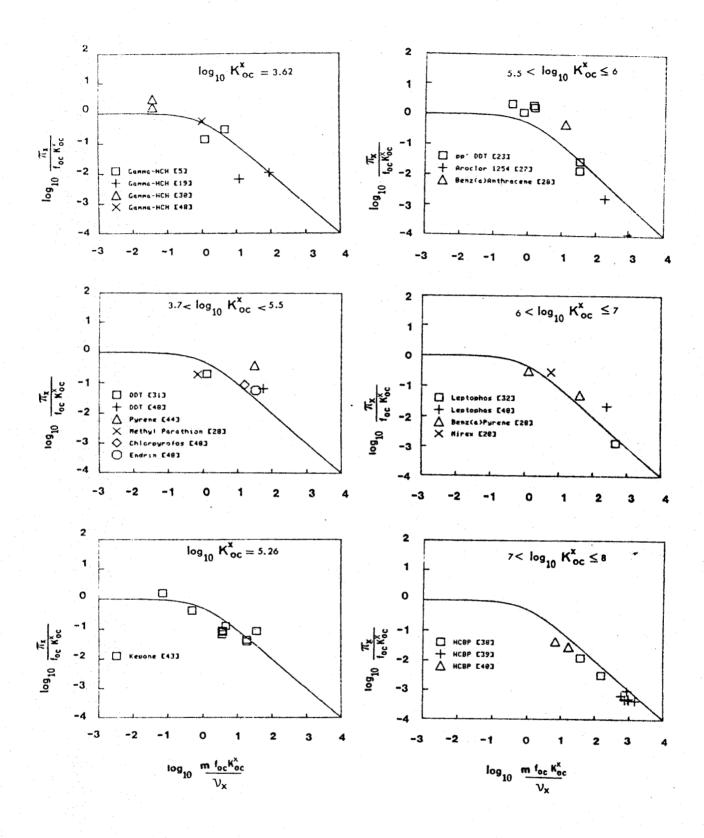


FIGURE 6

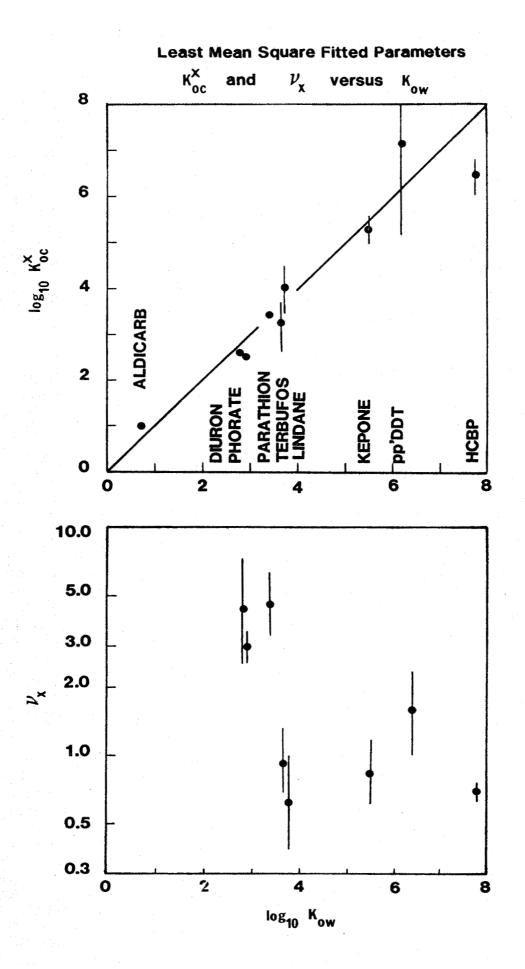
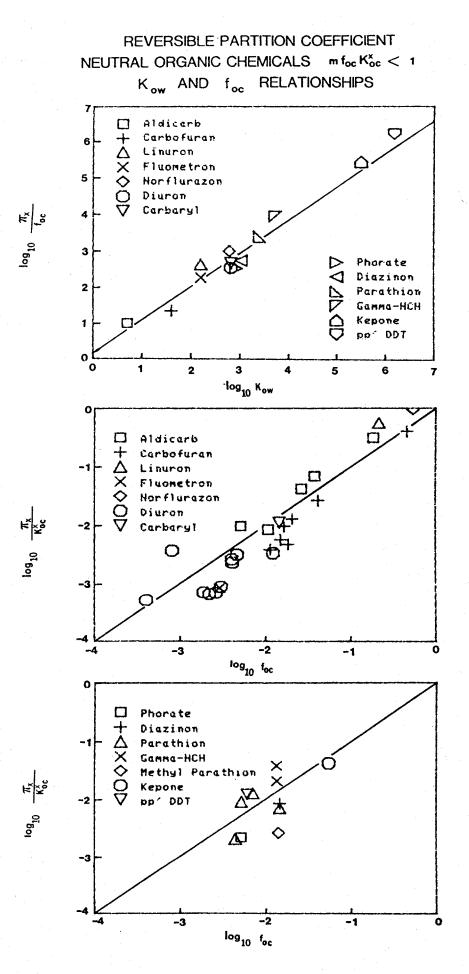


FIGURE 7

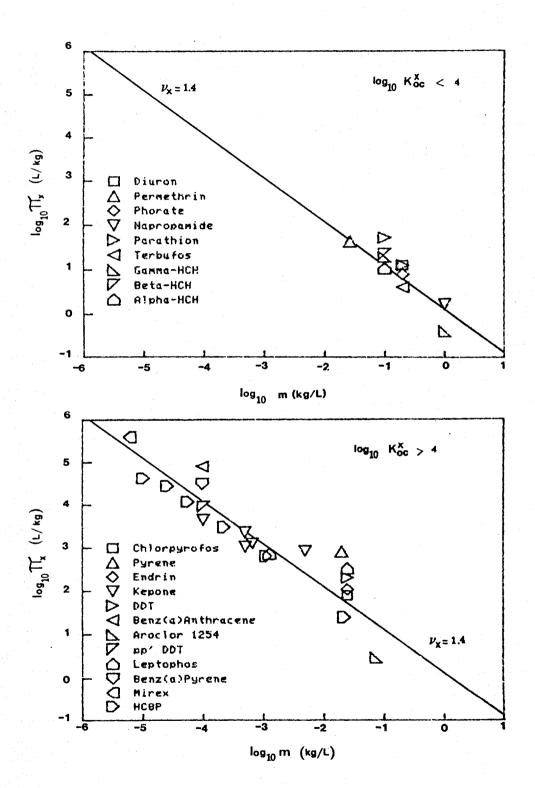
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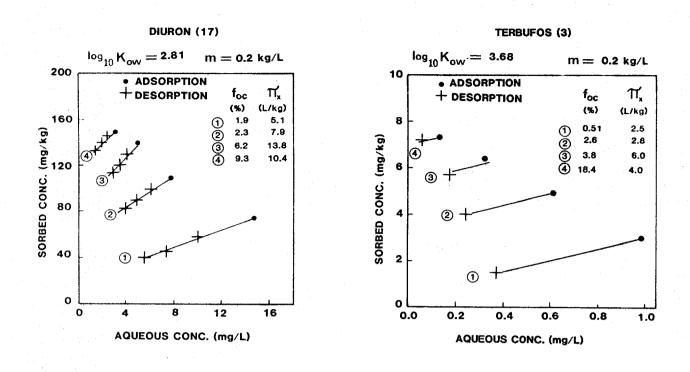


# REVERSIBLE PARTITION COEFFICENT

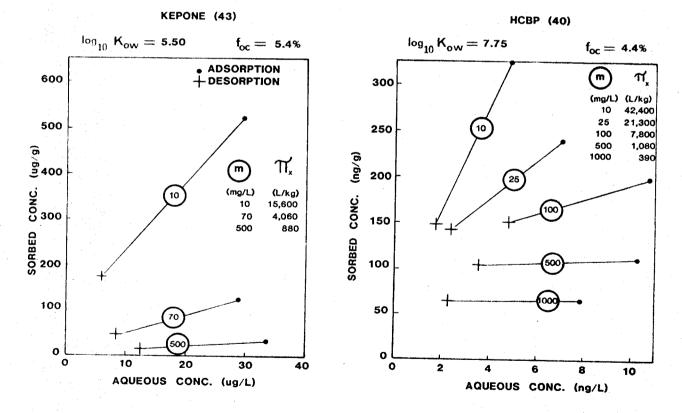
# NEUTRAL ORGANIC CHEMICALS

 $m f_{oc} K_{oc}^{\chi} > 3$ 





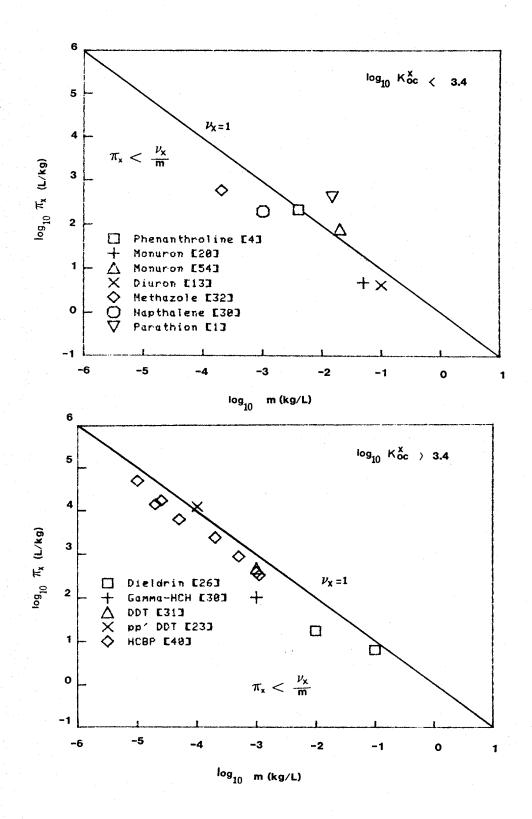
#### EFFECT OF PARTICLE CONCENTRATION

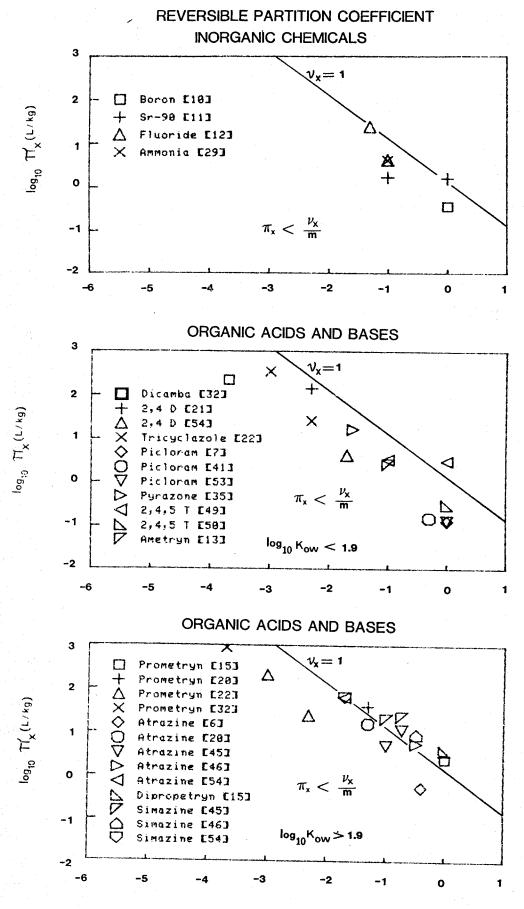


## **REVERSIBLE PARTITION COEFFICENT**

6 L .

NEUTRAL ORGANIC CHEMICALS-INORGANIC SORBENTS





بر م م

log<sub>10</sub> m (kg/L)

FIGURE 12

GURE 12

## APPENDIX III

### DIFFUSION AND PARTITIONING OF HEXACHLOROBIPHENYL IN SEDIMENTS

by

Dominic M. Di Toro John S. Jerís Daniel Ciarcía

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#### I. Introduction

The fate of persistent hydrophobic chemicals in natural bodies of water is greatly influenced by the transport mechanisms between sediments and the overlying water: particle transport (settling and resuspension) and aqueous diffusion. An analysis of each of these mechanisms is required if rational mass balance calculations of chemical fate are to be made. The purpose of this paper is to present an experimental investigation that is designed to evaluate the diffusive mass transport and partitioning of a highly sorbed chemical, hexachlorobiphenyl (HCBP), in sediments.

Conventional reversible sorption theory predicts that the apparent diffusion coefficient of total chemical,  $D_{a}^{*}$ , is (Berner, 1980):

$$D_{s}^{*} = \frac{D_{s}}{1 + m\pi/\phi}$$
(1)

where  $D_s$  is the aqueous diffusion coefficient in the interstitial water for nonadsorbing chemicals, m is the sediment solids concentration  $m = \rho_s (1 - \phi)$ ,  $\pi$  is the adsorption-desorption partition coefficient and  $\phi$  is the porosity. For the classical theory to apply, adsorption and desorption are assumed to be described by a linear reversible isotherm.

The initial concern was: what would be the effect of non-reversible HCBP adsorption-desorption, which has been observed in agitated suspended sediment experiments (Di Toro and Horzempa, 1982), on the rate of diffusion in sediments. The importance of the question is related to the burial of chemical by sedimentation. If contaminated sediment is buried by uncontaminated sediment faster than interstitial water diffusion can contaminate the newly deposited layer, then sedimentation in the absence of bioturbation provides an ultimate sink for sorbed chemical. For soil column leaching experiments, it has been found that significantly greater mass transport occurs than is implied by eq. (1) (Gersty, et al., 1979). The concern is that this effect may also be exhibited by highly sorbed chemicals such as HCBP. On the other hand, non-reversible behavior may inhibit migration as has been observed for soil column leaching experiments (van Genuchten, 1977).

The experiment is not without its practical difficulties. If l is the distance from the initially contaminated sediment layer, then the time, t, to reach this distance is on the order of:

$$\sqrt[4]{2D_{s}t}$$

l

For example, if  $m/\phi = 1 \text{ kg/L}$ ,  $\pi = 10^5 \text{ L/kg}$ , and  $D_s = 1 \text{ cm}^2/\text{day}$ , then after one day, the concentration is predicted to migrate a distance of  $\ell = 0.045$ millimeters and after 100 days, the distance is increased only ten-fold. Thus, it is necessary to be able to measure changes of concentrations in distances on the order of 0.1 mm and to be able to wait for time periods on the order of  $10^2$  days (Duursma, 1966; Walker and Crawford, 1970).

The experimental design is illustrated in fig. 1. The vessel is a hypodermic syringe from which top assembly is removed. A layer of uncontaminated sediment is deposited, allowed to settle, and the supernatant is removed. A layer of HCBP contaminated sediment is added, allowed to settle, and the supernatant is removed. A final layer of uncontaminated sediment is added, settled, the supernatant removed, and the syringe is capped.

After a specified period of time has elapsed for migration to take place, the syringe is mounted on a device which is equipped with a screwdriven mechanism that advances the syringe plunger in small (0.005-0.01 inch) increments. The exuded sediment slices are carefully removed and analyzed for radioactivity. Thus, it is possible to examine the profile of HCBP in 0.005-0.01 inch increments throughout the sediment column.

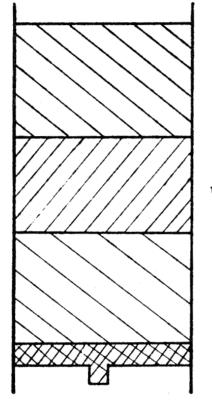
The initial design involved only HCBP tagged sediment. However, it was soon discovered that ambiguous results could be obtained since it is not possible to exclude the possibility that, initially, the interface regions between tagged and untagged sediment were not exactly perpendicular (to within  $\sim$  0.01 inch) to the syringe axis. As a consequence, the horizontal slices, through a tilted interface, would give the false impression of migration.

The solution to the problem was a double tag experiment. A small quantity of an absolute particle tracer,  $C^{14}$ -graphite particles, was added to the tritiated HCBP tagged sediment. The liquid scintillation counter can distinguish the tritium and carbon-14. The migration of the HCBP is then unambiguously monitored since if it appears in a sediment slice in excess of the  $C^{14}$ -graphite concentration, the only possibility is that aqueous phase diffusive transport has occurred.

2

(2)

# SEDIMENT MIGRATION EXPERIMENTAL DESIGN



Non-Radioactive Sediment

Sediment Containing <sup>3</sup>H-Hexachlorobiphenyl <sup>14</sup>C-Graphite

Non-Radioactive Sediment

Syringe Plunger

Figure 1.

Schematic Diagram of Experimental Vessel.

#### Materials and Methods

The sediment employed in these experiments was obtained from Saginaw Bay (Station #50)(see Di Toro et al., 1982 for station location and sediment properties). The sediment was wet sieved through a #200 sieve (75  $\mu$ m) to remove the coarse material which would interfere with the slicing. Table 1 lists the sieved sediment characteristics. The preparation of the dual tag sediment began with the transfer of the contents of an ampule containing 40  $\mu$ Ci of <sup>3</sup>H-HCBP into a 250 ml erlenmeyer flask. This <sup>3</sup>H-HCBP was a concentrated stock solution of custom synthesized 2,4,5,2',4',5' hexachlorobiphenyl in acetone with a high specific concentration (38.2 Ci/m mole New England Nuclear Corp.). The ampule was rinsed twice with acetone, and the rinse placed into the flask. The acetone was evaporated by purging the flask with a moderate flow of nitrogen which first passes through a distilled water trap. To prevent atmospheric contamination, the exit gas was passed through octanol and amyl alcohol traps.

After the acetone was evaporated, 25 grams of wet sediment (approx. 20 ml) and 10 ml distilled water were added to the flask to make the thick sediment less viscous. The flask was placed on a wrist action shaker and allowed to shake for four days. At this point, samples were removed from the flask, and analyzed for  $H^3$  radioactivity. The final step was the addition of 2  $\mu$ Ci C<sup>14</sup>-graphite particles (10 mCi/m mole) to the flask which was then shaken overnight. The weight fraction of added graphite was insignificant. The flask contents were allowed to settle for 24 hours, and 9 ml of supernatant were withdrawn with a syringe. The remaining sediment mixture containing both C<sup>14</sup>-graphite and sorbed H<sup>3</sup>-HCBP was used as the stock sediment for the dual isotope experiments. The concentration of sorbed HCBP was 23 ng HCBP/g dry sediment.

Disposable plastic syringes (3 ml, 0.8 cm I.D. Benton Dickenson) with their tips removed using a lathe to assure a parallel cut, were the experimental vessels. Approximately 0.4 ml of untagged sediment was added to the syringe and the sediment was allowed to settle for approximately 1 month. Any overlying water was carefully removed. Approximately 0.3 ml of tagged sediment was added in the same manner, and allowed to settle for four days. Overlying water was removed and 0.3 ml of untagged sediment added. Stoppers were placed in the top of the reaction vessel syringes, and the vessels placed in the humidifier.

Reaction vessels were removed from the humidifier at the appropriate time, and placed in the extraction device. The plunger of the reaction vessel syringe was moved slowly and accurately by using an unconfined compression machine (Soiltest). The proving ring normally used in this device was removed and replaced with a stationary mounting block. Displacement was measured by a strain gage connected to the mounting block. A single edged safety razor was used to slice the sediment. Glass tracks, mounted on the top of the block, were used to guide the cutting razor across the top of the reaction vessel. The upper end of the syringe is positioned flush with the glass tracks and clamped securely. A slice is taken by raising the platform the desired slice thickness and slowly moving the razor along the glass tracks on the mounting block. The sediment slice, which adheres to the razor, is carefully jetted with water into a 20 ml counting vial using a l ml syringe with 1 ml of distilled water, 10 ml of Aquasol is added, and mixed using a vortex mixer for 15 sec. The vial is counted for 2 minutes on a liquid scintillation counter using the carbon 14 window. The vial is then centrifuged at 6000 g for 15 min. in order to remove both the C<sup>14</sup>-graphite and sediment particles. The toluene base Aquasol extracts the HCBP into the liquid phase, which is decanted into another vial, and then counted for 2 minutes using the tritium window to measure the concentration of  ${}^{3}$ H-HCBP in the slice.

Identically prepared syringes were used to measure the diffusion coefficient of tritiated water. To a 3 ml syringe containing 1.4 ml of consolidated sediment, 300  $\mu$  of distilled water was added and after 1 hr, 30  $\mu$  of <sup>3</sup>H-H<sub>2</sub>O is added. Samples of this overlying water (5  $\mu$ ) are taken at intervals and counted.

#### Analysis of Tritiated Water Diffusion

An independent measurement of the interstitial water diffusion coefficient can be made using a non-adsorbing substance. Tritiated water is an ideal choice since its presence does not significantly alter the ionic composition of the interstitial water as would be the case if, for example, high concentrations of a major cation or anion were introduced.

The conventional method (e.g. Hesslein, 1980) is to introduce a quantity of tracer into the water overlying the sediment, and after the passage of a suitable length of time, to measure the depth distribution of the tracer.

The problem with this approach when applied to the syringe experimental design is that a considerable time is required to slice the sediment and, during this time, the tracer continues to diffuse within the progressively shortening sediment column. Although it is possible to numerically compute the spatial distribution of tracer to be expected for this situration, the method lacks a directness which can be achieved by simple modification.

Instead of relying on the depth distribution of the tracer at a fixed time, it is possible to measure the decline of tracer concentration in the overlying water with time as the tracer diffuses into the sediment (e.g. Officer, 1982). If the volume ratio of overlying water to interstitial water is suitably small so that the decline is appreciable, and the tracer measurement is sufficiently sensitive so that only small volumes need to be removed from the overlying water, then the analysis of this experiment is direct.

Let  $\ell_1$  be the depth of the overlying water, and  $\ell_2$  be the depth of sediment. Let  $C_{T1}(t)$  and  $C_{T2}(z,t)$  be the tracer concentrations in the overlying and interstitial water, respectively. The mass balance equation for the overlying water is:

$$\ell_1 A \frac{dC_{T1}}{dt} = D_s A \frac{\partial (\phi C_{T2})}{\partial z} \Big|_{z=0}$$
(3)

where A is the interfacial area. This equation states that the time rate of change of tracer in the overlying water is due to the diffusive flux of tracer into the sediment. The mass balance equation in the sediment is:

$$\frac{\partial(\phi C_{T2})}{\partial t} - D_s \frac{\partial^2(\phi C_{T2})}{\partial z^2} = o$$
 (4)

The boundary conditions correspond to continuity of fluid tracer concentration at z=o:

$$C_{T1}(t) = C_{T2}(o,t)$$
 (5)

and a zero flux condition at the sediment bottom:

$$\frac{\partial C_{T2}(z,t)}{\partial z} \Big|_{z=\ell_2} = 0$$
(6)

The solutions of these equations is known to be (Carslaw and Jaeger, 1959):

$$C_{T1}(t) = \frac{C_{T1}(o)}{1+k} \left[1 + 2k(k+1) \sum_{i=0}^{\infty} \frac{\exp(-\alpha_i^2 T)}{P_i}\right]$$
(7)

$$C_{T2}(z,t) = \frac{C_{T1}(o)}{1+k} \left[1 + 2k(k+1)\sum_{i=0}^{\infty} \frac{\exp(-\alpha_i^2 T) \cos(\alpha_i^2(1-z/\ell_2))}{P_i \cos(\alpha_i)}\right]$$
(8)

where:

 $\mathbf{k} = \ell_2 \phi / \ell_1 \tag{9}$ 

$$P_{i} = \alpha_{i}^{2} + k(k+1)$$
 (10)

$$T = D_{s}t/\ell_{2}^{2}$$
(11)

and  $\alpha_i$  are the positive roots of:

$$k \tan \alpha_i + \alpha_i = 0 \tag{12}$$

For small times a fairly large number of roots are required for accurate solutions but no computational problems arise. The solution is easily checked at t=o to determine the maximum number of terms required in the sums. Note that at t  $\rightarrow \infty$ , the solutions conserve mass:

$$\ell_{1} C_{T1}(o) = \ell_{1} C_{T2}(\infty) + \ell_{2} \phi C_{T2}(z,\infty)$$
(13)

as they should.

One refinement is required is to account for the effect of sample withdrawal on the volume, and therefore the depth,  $l_1$ , of the overlying fluid. Initially  $v_1 = l_1 A = 330 \ \mu l$  whereas at the end of the experiment,  $v_1 = 255 \ \mu l$ . This is enough of a change to influence the computations. However, instead of solving this problem directly, the following approximation is employed. At each time, t, the average depth:

$$\bar{\ell}_{1}(t) = \frac{1}{t} \int_{0}^{t} \ell_{1}(t) dt \qquad (14)$$

is computed and this is used in equation (7) for  $C_{T1}$ . While this is not a rigorously correct solution it appears to be a reasonable approximation. One check is to compare the actual tracer mass lost via sampling to that removed from the computation by using  $\bar{k}$ ,(t):  $(k_1 - \bar{k}_1(t))C_{T1}(o)$  at each sampling time. The average absolute error for the volume ratio employed in this experiment is 17%. Errors are positive during the early periods (more mass lost than computed) and negative in the latter stages. However the approximation appears reasonable and simple to apply.

The results of the analysis are shown in fig. 2a, for two separate experiments. The coefficients and relevant parameters are listed in Table 2. The apparent sediment diffusion coefficient is found to be  $D_{s, H_20} = 1.5 \text{ cm}^2/\text{day}$ .

As a confirmation of this result, a spatial profile that resulted after a exposure time of 2.5 hrs is shown in fig. 2b. Beyond the first few slices, the flattening of the data profile due to the increased dispersion during the slicing is evident. Nevertheless the data are in reasonable agreement with the dispersion coefficient estimated from the penetration experiment.

The apparent sediment diffusion coefficient found from these experiments can be compared to empirical correlations that apply to natural sediments. The relationship between the molecular diffusion coefficient, D, and the apparent sediment diffusion coefficient,  $D_e$ , is (Berner, 1980):

$$D_{s} = D/\theta^{2}$$
(15)

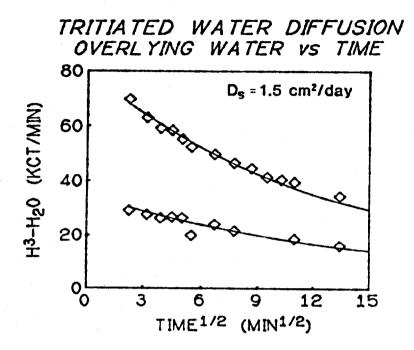
where  $\theta$  is the tortuosity: the ratio of the length of the actual diffusion path to the linear length of the sediment. It has been shown that:

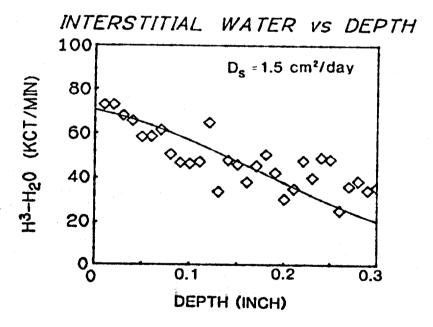
$$\theta^2 = \phi F \tag{16}$$

where F is the formation factor:  $F = R/R_0$ , the ratio of the electrical resistivity of the bulk sediment to the pore fluid. Manheim (1970) has shown that:

$$\mathbf{F} = \boldsymbol{\phi}^{-\mathbf{n}} \tag{17}$$

where n = 2.8. If n = 2, this relationship is known as Archie's law. Thus:







Tritiated Water Diffusion. Overlying water concentration versus time (top). Total  $H^3-H_2^0$  concentration versus depth after t = 2.5 hr. (bottom).

$$\mathbf{D}_{\mathbf{a}} = \mathbf{D}_{\mathbf{b}} \boldsymbol{\phi}^{\mathbf{1},\mathbf{8}} \tag{18}$$

For the syringe experiments:  $\phi = 0.75$  and  $D_{H^3-H_20} = 1.81 \text{ cm}^2/\text{day}$  (Wang et al., 1953) so that the expected value is  $D_{s,H_20} = 1.08 \text{ cm}^2/\text{day}$  which can be compared to  $D_{s,H_20} = 1.5 \text{ cm}^2/\text{day}$  found experimentally.

#### Analysis of HCBP Diffusion

The distribution of an adsorbing chemical in a sediment is determined by the diffusive migration in the interstitial water and its adsorptiondesorption behavior. If reversible adsorption-desorption is assumed, with a linear isotherm applicable to both adsorption and desorption, and if adsorption-desorption equilibrium is attained on a time scale which is short relative to the diffusive time constant, then it has been shown (e.g. Berner 1980) that the governing equation is:

$$\frac{\partial(\phi C_{T})}{\partial t} - \frac{\partial}{\partial z}(\phi D_{s}^{*} \quad \frac{\partial C_{T}}{\partial z}) = o \qquad (19)$$

For the syringe experiments, the initial condition is:

$$C_{T}(z,o) = C_{To} \qquad z > o \qquad (20)$$

$$C_{T}(z, o) = o$$
  $z < o$  (21)

where z=o is the interface boundary and the syringe boundaries are far enough from the interface so that an infinite spatial domain is a reasonable approximation. For these conditions the solution is known to be (Crank, 1975):

$$C_{T1}(z,t) = \frac{C_{T0}}{2} \operatorname{erfc}(\frac{-z}{2\sqrt{D_{t}^{*}t}}) \qquad z < o \qquad (22)$$

$$C_{T2}(z,t) = \frac{C_{T0}}{2} [1 + erf(\frac{z}{2\sqrt{p_s^{*}t}})] z > o$$
 (23)

where:

$$f_{d} = \frac{1}{1 + m\pi/\phi}$$
(24)

$$D_{s}^{*} = f_{d} D_{s,HCBP}$$
(25)

m =  $\rho_s(1-\phi)$ , the solids concentration and  $\pi$  is the partition coefficient.

The experimental measurements do not produce the continuous profiles that actually exist but rather the slice average concentrations. The theoretical slice average concentrations can be computed from eqs. (22-23) by spatially averaging the solutions:

$$\overline{C}_{T}(z,t) = \frac{1}{d} \int_{z}^{z+d} C_{T}(z,t)dt \qquad (26)$$

to yield:

$$\overline{C}_{T1}(z,t) = \frac{C_{To}}{2} \quad \frac{\sqrt{4 p_s^{*}t}}{d} \left[ \operatorname{ierfc}(\frac{-z - d}{2\sqrt{p_s^{*}t}}) - \operatorname{ierfc}(\frac{-z}{2\sqrt{p_s^{*}t}}) \right]$$
(27)

$$\overline{C}_{T2}(z,t) = \frac{C_{T0}}{2} \left\{ 2 - \frac{\sqrt{4\tilde{D}_{s}t}}{d} \left[ \operatorname{ierfc}(\frac{z}{2\sqrt{D_{s}t}}) - \operatorname{ierfc}(\frac{z+d}{2\sqrt{D_{s}t}}) \right] \right\}$$
(28)

where:

$$\operatorname{ierfc}(z) = \frac{1}{\sqrt{\pi}} e^{-z^2} - z \operatorname{erfc}(z)$$
 (29)

is the integral of the complimentary error function and d is the slice thickness.

In order to estimate the partition coefficient via eqs. (24-25) the molecular diffusion coefficient of HCBP,  $D_{HCBP}$  is required. This is estimated using the Wilke-Chang correlation to molal volume (Reid, et al., 1977). The latter is computed from the LeBas additive volume increments of the molecular constituents (C,H,Cl, corrected for the two 6-member rings). The result is  $D_{HCBP} = 0.461 \text{ cm}^2/\text{d}$  (20°C). Using the known molecular diffusivity of water,  $D_{H_2O} = 1.81 \text{ cm}^2/\text{d}$  (20°C)(Wang et al., 1953) and the measured tri-

tiated water diffusion coefficient in the syringe:  $D_{s,H_20}^* = 1.5 \text{ cm}^2/\text{d}$  it follows that the molecular diffusivity of HCBP in the syringe in the absence of adsorption is:  $D_{s,HCBP} = (D_{s,H_20}/D_{H_20})D_{HCBP} = 0.382 \text{ cm}^2/\text{d}$ . The apparent diffusivity of HCBP follows from:

$$D_{s,HCBP}^{\star} = \frac{D_{s,HCBP}}{1 + m\pi/\phi}$$
(30)

for partition coefficient,  $\pi$ .

A correction is required for the tilted interface, since slicing through it produces a profile which indicates an apparent diffusive motion. The method adopted is to simultaneously fit the  $C^{14}$ -graphite and  ${}^{3}$ H-HCBP data profiles to the solutions (eqs. 27-28 with  $C_{To} = 1$ ) normalized as:

$$N_{C}(z) = N_{CO} + (N_{C1} - N_{CO})\overline{C}_{T}(z, t_{g}, D_{s, HCBP}^{*})$$
 (34)

$$N_{H}(z,t) = N_{HO} + (N_{H1} - N_{HO})\overline{C}_{T}(z,t_{g}+t,D_{s,HCBP}^{*})$$
(35)

where  $N_{CO}$ ,  $N_{HO}$ :  $N_{C1}$ ,  $N_{H1}$  are the background (untagged region) and the tagged region counting rates for  $C^{14}$  and  $H^3$  respectively far from the interface. The location of the interface, z=o, is chosen by examining the  $C^{14}$  profiles. The idea is that at t = o both the graphite and HCBP profiles are identical (this is the initial condition) and as t increases, HCBP diffuses relative to the graphite. This strategem eliminates the need to solve eq. (19) for an arbitrary initial condition. Rather a convenient initial condition:  $C_T(z,t_g,D_{s,HCBP}^*)$  is used which is itself a solution of eqs. (27-28) with  $t_g$ chosen to fit the graphite profile. As time, t, increases the HCBP profile evolves from that initial condition.

The experimental data: the logs of the  ${}^{3}$ H,  ${}^{14}$ C, and  ${}^{3}$ H/ ${}^{14}$ C ratio are fit to eqs. (34-35) using a nonlinear least squares method (Marquardt, 1963) to estimate the six unknown parameters:  $\pi$ , t<sub>g</sub>, N<sub>CO</sub>, N<sub>C1</sub>, N<sub>HO</sub>, N<sub>H1</sub>.

Figure 3 illustrates the fits obtained for various elapsed times. Slice average concentrations are connected with a smooth curve for ease of presentation. The diffusive motion of HCBP relative to the graphite particle profile is evident. The observed HCBP profiles are remarkably regular and agree



**GRAPHITE vs DEPTH** 

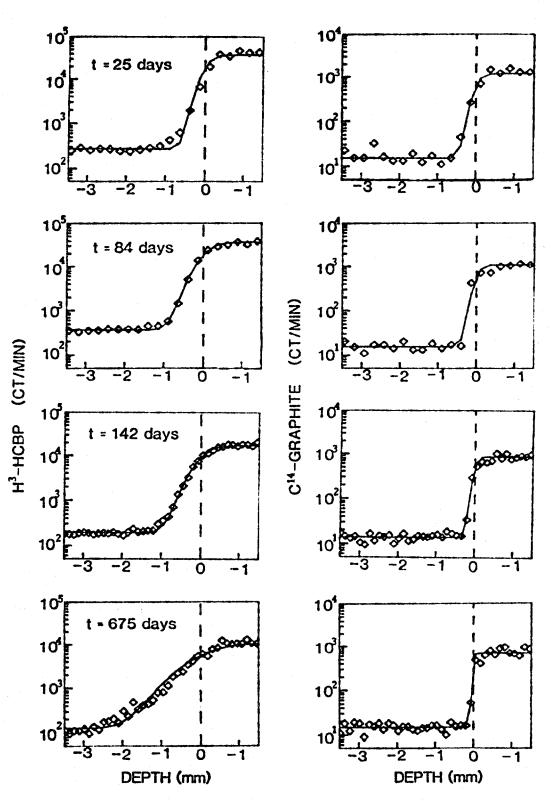


Figure 3. HCBP and graphite concentration versus depth at various elapsed times. Curves are computed using parameters in Table 3.

with the theoretical calculations over two orders of magnitude of HCBP concentration. Initially the slice thickness used was 0.01 inch. However it was found that a more refined profile could be obtained using a slice thickness of 0.005 inches. This was the procedure used for the syringes with t >140 days.

The initial apparent diffusive motion of HCBP at t = 25 day is actually due to a tilted interface. This can more clearly be seen in fig. 4 which presents the  ${}^{3}$ H-HCBP/ ${}^{14}$ C-graphite count ratio. This is the most sensitive measure of the extent of HCBP diffusion. Excess  ${}^{3}$ H/ ${}^{14}$ C counting rate over the background ratio to the left of the interface can only be the result of diffusive migration of the  ${}^{3}$ H-HCBP in the interstitial water since the graphite particles are immobile. Depletion of HCBP to the right of the interface is also detectable at longer elapsed times. The computed slice average  ${}^{3}$ H/ ${}^{14}$ C concentration ratios, N<sub>H</sub>(z,t)/N<sub>C</sub>(z) are displayed as a staircase function over the slice thickness so that a more detailed comparison can be made to the data.

The lack of significant excess HCBP for t = 25 days indicates that no substantial interfacial mixing occurred during the preparation of the syringe or during the slicing procedure. The extent of migration increases regularly as elapsed time increases. However even at the termination of the experiment (almost two years) the extent of migration is no more than 2-3 mm. It is this fact that necessitates the dual tag experimental design. Table 3 tabulates the parameter estimates.

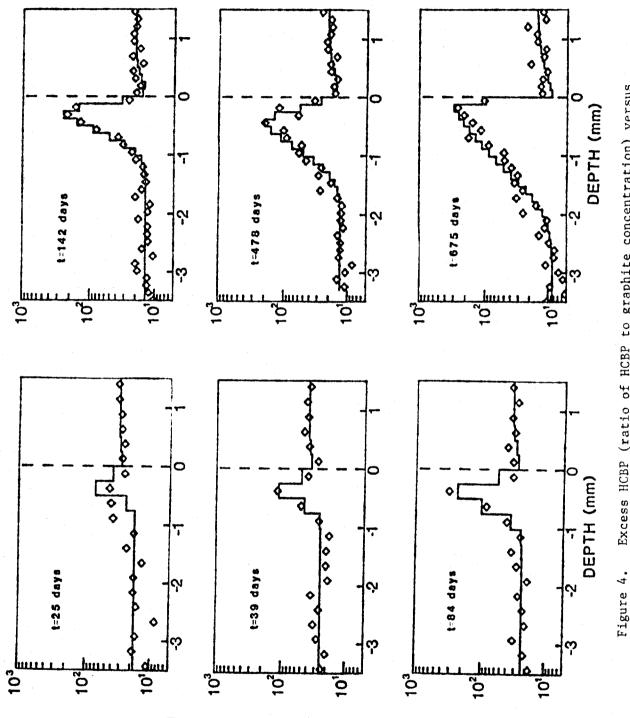
#### Discussion and Conclusions

As illustrated in figs. 3-4 the diffusion-linear reversible equilibrium model of chemical transport in sediments satisfactorily reproduces the individual syringe observations. The estimates of apparent sediment diffusion and partition coefficients versus elapsed time are shown in fig. 5. No temporal trends are evident so that the assumption of constant linear partition coefficient,  $\pi = 10^{5.06 \pm 0.13}$ , and therefore apparent sediment diffusivity,  $D_s^* = 3.95 \pm 1.12 \times 10^{-6} \text{ cm}^2/\text{d}$ , is reasonable.

The partition coefficient for HCBP obtained from these experiments is compared in fig. 6 and Table 4 to predictions from empirical hydrophobic

**EXCESS HCBP vs DEPTH** 

.



H3-HCBP/C14-GRAPHITE

• Excess HCBP (ratio of HCBP to graphite concentration) versus depth at various elapsed times. Staircase curves are slice

average ratios, computed using parameters in Table 3.

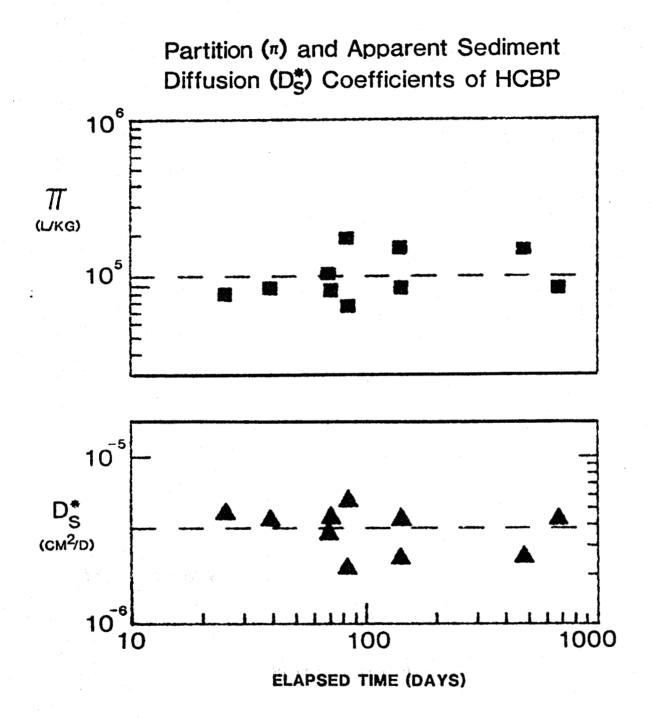


Figure 5. Partition coefficient (upper and left axis) and apparent sediment diffusion coefficient (lower and right axis) versus elapsed time for all syringes analyzed.



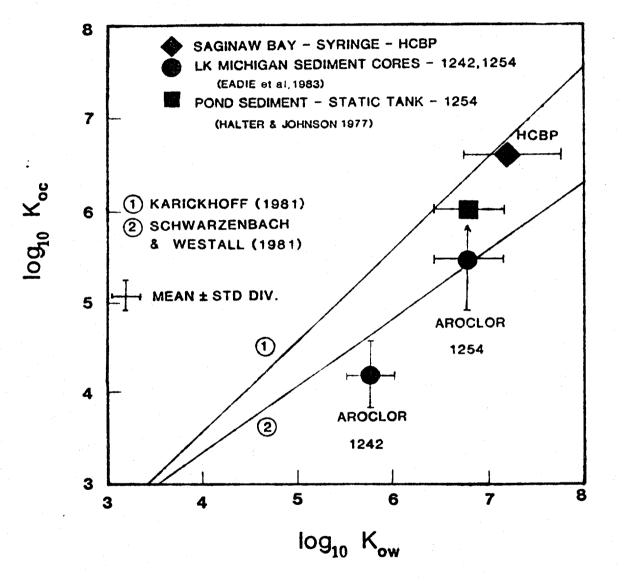


Figure 6. Sediment Organic Carbon Partition Coefficient (K<sub>oc</sub>) versus Octanol-Water Partition Coefficient (K<sub>ow</sub>) for HCBP and PCB Aroclors.

partitioning equations based on octanol-water partition coefficient and with the few available direct (Eadie et al., 1983) and indirect (Halter and Johnson, 1977) observations of PCB partitioning in sediments. The octanolwater partition coefficients,  $K_{ow}$ , are somewhat uncertain for PCBs. Table 5 summarizes the available values. The value for HCBP is reasonable accord with the predictions of the Karickhoff (1981) correlations whereas the Lake Michigan results are somewhat lower although they are in reasonable accord with the Schwarzenback & Westall (1981) correlation.

The situation with respect to inferred apparent diffusivities is less satisfactory. Static leaching experiments using contaminated sediments have been reported in two cases. An analysis based upon observed fluxes measured by recirculating the overlying water through a foam plug (Fisher et al., 1983) indicated sediment diffusivities of  $D_s^* = 0.0012 - 0.85 \times 10^{-6}$  cm/d for four PCB isomers. An analysis of the reported overlying water PCB concentration versus time in a static quiescent leaching experiment (Halter & Johnson, 1977) which parallels the analysis used for the tritium penetration experiment yields  $D_s^* = 0.05 \times 10^{-6} \text{ cm}^2/\text{day}$ . These values are one to three orders of magnitude lower than the syringe results.

For both these experiments the organic carbon fractions were similar to the syringe sediment so that these apparent diffusivities would suggest PCB partition coefficient of  $\log_{10} K_{oc} \sim 7.5 - 10.5$  by comparison to  $\log_{10} K_{oc} =$ 6.6 found for the syringe experiments (fig. 6). Such large partition coefficients appear unrealistic. It may be that a large fraction, 90-99.9%, of the total PCB in the sediments used in the leaching experiments are nonreversibly bound. Since the framework used to estimate  $D_s^*$  assumes completely reversible behavior the inferred partition coefficient would be artificially large in order to account for the large non-reversible fraction bound to the particles.

It should be pointed out that although the syringe experiments exhibited fully reversible behavior, this may not apply to native PCBs in sediment. For the Saginaw Bay sediment employed, 23 ng <sup>3</sup>H-HCBP/g were added to sieved Saginaw Bay Sta. #50 sediment. Non-sieved Saginaw Bay sediments from similar locations have native Arochlor 1254 concentrations of  $\sim$  400 ng/g (Richardson, 1983; Thomann and Mueller, 1983). Since 2,4,5,2',4',5' HCBP is 3.67% of Arochlor 1254 (Rapaport & Eisenreich, 1984) it is likely that > 15 ng HCBP/g of native chemical already was present in the sieved experimental sediment. In terms of the reversible-resistant model of PCB sorption (Di Toro & Horzempa, 1982) the additional radiolabeled chemical may not have sorbed to any available resistant sites (which would happen only if the tagged concentration were in excess of the native concentration) but, instead, only isotopically exchanged with the native reversible component already present.

Since only the reversible component was radiolabeled, and since only very small concentrations migrated into the untagged (but HCBP containing) sediment, it is likely that the syringe experiment actually measured the reversible component partition coefficient, x. Whether native PCBs are fully reversible remains to be determined.

It is also possible that unlike quiescent diffusion within a sediment the transport of PCBs between the sediment and overlying water in a leaching experiment is not well modeled using diffusion equation formulations. Very thin layers ( 1 mm) are involved in short term leaching experiments and the mass transport mechanisms may not be entirely diffusive in character. A one to three order of magnitude discrepancy invites a more careful experimental investigation of the mechanisms controlling static leaching fluxes.

The dual tag syringe experiments clearly indicate that molecular diffusion and linear reversible partitioning are the mechanisms that determine the rate of mobility of the (reversible) component of HCBP. The resulting partition coefficient is in conformity with empirical correlations. This methodology may be the most direct route for measuring partition coefficients of highly sorbed chemicals.

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# TABLE 1

Sediment Properties Saginaw Bay Sta. #50 (< 75 µm fraction)

Organic	Matter		Ŧ	7.6%
Organic	Carbon	Fraction		-
Surface	Area		=	12.8 $m^2/g$

Size Fractions	% Total Weight
(µm)	
1-2	4
2-10	15
10-30	20
30-75	61

TABLE 2

Tritiated Water Diffusion Experiment

Porosity in Syringe:	$\phi = 0.75$
Solids density:	$p_{s} = 2.65 \text{ g/cm}^{3}$
Solid-liquid phase ratio: m	$\sqrt{\phi} = \rho_{\rm s}(1-\phi)/\phi = 0.883 \text{ kg/l}$
Sediment Volume:	$v_2 = 1400 \ \mu L$
Diffusion Coefficient: D <sub>s,H</sub>	$_{20}^{2} = 1.5 \text{ cm}^{2}/\text{day}$

Overlying Water Volume

t(min)	$\frac{v_1(\mu L)^*}{2}$
0	330
30	290
120	255
180	255

\* Volume is linearly interpolated between the tabulated values

Elapsed time t(days)	Initial Condition t <sub>g</sub> (days)	$2^{\text{D}_{\text{s,HCBP}}^{\text{h}}}$ (cm <sup>2</sup> /d x 10 <sup>-6</sup> )	log <sub>10</sub> π (ℓ/kg)	log <sub>10</sub> S.E. of π (ℓ/kg)
25	23.3	4.82	4.95	0.12
39	28.3	4.43	4.99	0.058
69	37.3	3.66	5.07	0.16
71	12.4	4.56	4.98	0.042
83	0.21	2.26	5.28	0.042
84	16.1	5.70	4.88	0.036
140	28.4	2.60	5.22	0.046
142	9.7	4.44	4.99	0.019
478	45.1	2.66	5.21	0.027
675	0.35	4.43	4.99	0.022
		3.95 ± 1.12	5.06 ±0.13	

TABLE 3Estimates of Diffusion and Partition Coefficients

		PCB Sediment Part	Partition Coefficients	
Chemica1	Sediment	Organic Carbon fraction (%)	log <sub>10</sub> Koc = π/foc ± log <sub>10</sub> Std. Dev.	Reference
Aroclor 1242	Lake Michigan	0.7-3.8	4.17 ± 0.38	Eadie et al. (1983)
Aroclor 1254	Lake Michigan	0.7-3.8	5.44 ± 0.55	Eadie et al. (1983)
Aroclor 1254	Pond	2.0	6.0	Halter & Johnson (1977)
2,4,5,2',4',5'	Saginaw Bay	2.8	<b>6.61</b> ± 0.13	This study
		TAB PCB Octanol-Water P	TABLE 5 PCB Octanol-Water Partition Coefficient	
Aroclor 1242	Aroclor 1254		HCBP (2,2',4,4',5,5')	Reference
5.58 - 5.90 <sup>(a)</sup>	6.72 6.47 7.17(2	a)	- 6.72 7.44 7.44	Veith & Morris (1978) Veith et al. (1979) Chiou et al. (1977) Rapaport & Eisenreich (1984) Rapaport & Eisenreich (1984)
5.74 ± 0.23	6.79	± 0.35	7.30 ± 0.53	Mean ± Std. Deviation.
(a) Computed as Σf <sub>j</sub> míxture.	Computed as Σf <sub>1</sub> Kow <sub>1</sub> /Σf <sub>1</sub> each fraction, mixture.	م برا لبا	with octanol-water parti	of isomer with octanol-water partition coefficient $\operatorname{Kow}_{\mathbf{i}}$ in the

TABLE 4 Lent Partition Coeffici

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